NOBEL LAUREATE ROALD HOFFMANN

molecular
orbitals of
transition
metal complexes


YVES JEAN


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# Molecular Orbitals of Transition Metal Complexes 

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## Foreword

We are so divided. By the formal structure of university instructionorganic chemistry, inorganic chemistry, physical chemistry. By the incredible and unnecessary specialization of our journals. The molecular bounty we have ourselves created seems simply overwhelming-no wonder we seek compartmentalization in self-protection: It is easy to say, 'I'm an expert in Field x. And while I will listen to a seminar in y or $z$ (when I have time), please . . . let me be happy just in keeping up with my own field.'

The dangers of specialization are obvious-inbreeding, lack of scope, a kind of rococo elaboration of chemical complexity within a field. And we know that new ideas often come from an almost metaphorical importation of a way of thinking or a technique from another area.

Meanwhile, all along, nature persists in subverting the compartmentalizing simplicity of our minds. Through enzymes whose seeming magic is done by metal atoms and clusters at the active site, inorganic chemistry and biochemistry are rejoined. Transition metal carbides put organic carbon into some most unusual, patently inorganic environments. And, beginning in 1950, the explosion in organometallic chemistry has given us an incredible riches of structures and reactionsfrom ferrocene to olefin metathesis, metal-metal multiple bonds, to $\mathrm{C}-\mathrm{H}$ activation, and remarkable olefin polymerization catalysts. All from a combination of inorganic and organic chemistry.

Organometallic chemistry from its beginning also depended on, and also built, another bridge. This is to theoretical chemistry. The first, rationalizing accounts of the electronic structure of ferrocene and the Dewar-Chatt-Duncanson picture of metal-olefin bonding were followed by milestones such as the prediction of cyclobutadieneiron tricarbonyl and Cotton's beautiful elaboration of the idea of a metal-metal quadruple bond. The work of Leslie Orgel played a very important role in those early days. There were fecund interactions all along-compounds leading to calculations, and calculations pushing experimentalists to make new molecules. Often the theory was done by the experimentalists themselves, for the best kind of theory (the one that keeps the fertile dance of experiment moving) is a portable one. As easy to use molecular orbital theory, the theory of choice of the times, most certainly was and is.

It is hard to imagine a contemporary course in organometallic chemistry which does not contain a hefty, albeit qualitative component of molecular orbital theory. Yves Jean (with François Volatron) earlier wrote a classic teaching text on the orbitals of organic molecules. Here he has applied his great pedagogical skills to the construction of a beautifully thought through exposition of bonding in organometallic chemistry. Our undergraduate and graduate students will enjoy this book. And they, the chemists of the future, will use the knowledge gained here to enlarge our experience with new organometallic molecules, subverting once again the arbitrary division of organic and inorganic chemistry. Molecules whose beauty and utility we still cannot imagine.

Roald Hoffmann

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This book also owes a large debt to the students who have followed my lectures. In order to transmit knowledge and understanding, even in an area which is completely familiar, it is necessary to clarify one's ideas so as to make a consistent and, if possible, attractive presentation. Students' comments, remarks, and questions really help the teachers to achieve this goal. When it comes to writing a manuscript, one realizes just how useful the slow development and maturation of ideas has been.

I am very honoured that Roald Hoffmann accepted to write the preface, and the reader will discover during the chapters just how much the book owes to him.

Ecole Polytechnique (Palaiseau)
Université Paris Sud (Orsay)
April 2004

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## Introduction

This book starts from the most elementary ideas of molecular orbital theory, and it leads the reader progressively towards an understanding of the electronic structure, of the molecular geometry and, in some cases, the reactivity of transition metal complexes.

The use of simple notions, such as symmetry, overlap, and electronegativity, allows a qualitative method of analysis of the electronic structure of complexes, and of the properties which follow from it such as geometry or reactivity, to be developed. Qualitative in the sense that, for example, it enables us to understand why the structure of a particular complex is tetrahedral rather than planar, without being able to provide a reliable numerical value of the energy difference between these two structures. The quantitative level can be attained elsewhere-as is now standard practice in our laboratories-by more accurate methods such as $a b$ initio or density functional theories. But to interpret the results provided by more complex calculations, it is often necessary to return to the fundamental notions of symmetry, overlap, and electronegativity.

The qualitative approach used here is mainly based on the analysis of orbital interactions (atomic or molecular). Its application to transition metal complexes developed rapidly from about 1975, the leading exponent being Roald Hoffmann, winner of the Nobel prize for chemistry in 1981 with Kenichi Fukui. As a result, many experimental results can be rationalized, that is to say understood, on the basis of analyses and using a language that are accessible to every chemist. A colleague, Marc Bénard, spoke in the introduction to one of his lectures of the prodigious decade 1975-85 . . . Moreover, it has been possible to apply this approach to all of chemistry (organic, inorganic, organometallic, and the solid state), which is one of its strongest points. These are no doubt the main reasons for its success which has spread far beyond the realm of specialists: as Roald Hoffmann writes in the preface, it is a transferable theory which has marked our time.

It is certainly transferable to students, and the aim of this book is to encourage that process. By learning this method for the theoretical analysis of molecular electronic structure, a method which has so profoundly changed our approach to chemistry, the reader may be
encouraged to continue his exploration of the methods of quantum chemistry which nowadays are part of all chemical research.

In the first chapter, we present the rules for electron counting in transition metal complexes, the different coordination modes adopted by ligands and the essential properties of the orbitals that are involved on the metal and on the ligands. The main ligand fields are studied in the second chapter, where we limit ourselves to $\sigma$-type interactions between the metal and the ligands. The structure of the d block is established; knowledge of this structure, which is essential for transiton metal complexes, enables us to explore the relationships between the electronic configuration of complexes and their geometry. In the third chapter, we study the ways in which the analysis is changed when the ligands have $\pi$-type interactions with the metal (both $\pi$-donor and $\pi$-acceptor ligands). All these ideas are then used in the fourth chapter, which is a series of examples that illustrate how, starting from a knowledge of the orbital structure of complexes, we can understand their geometrical structure and, sometimes, their reactivity. The fifth chapter discusses the 'isolobal analogy' which shows how the electronic structures of transition metal complexes and of organic molecules can be related. A bridge is thus constructed between these two areas of chemistry that allows us to understand several resemblances (in particular, concerning structures) between species that appear to be very different. The last chapter contains a presentation of basic Group Theory, with applications to some of the complexes studied in the earlier chapters. This chapter is placed at the end of the book so as not to disrupt the flow of the more chemical aspects of the presentation, but the reader may consult it, if necessary, as and when reference is made to it in the book.

## Setting the scene

Transition metal complexes are molecules containing one or more metallic centres ( $\mathrm{Ti}, \mathrm{Fe}, \mathrm{Ni}$, etc.) bound to a certain number of 'ligands'. These latter may be atoms ( $\mathrm{H}, \mathrm{O}, \mathrm{Cl}$, etc.), molecular fragments $\left(\mathrm{CR}_{3}\right.$, $\mathrm{NR}_{2}, \mathrm{SH}$, etc.), or molecules that are themselves stable in the absence of any interaction with a metal ( $\left.\mathrm{NR}_{3}, \mathrm{PR}_{3}, \mathrm{R}_{2} \mathrm{C}=\mathrm{CR}_{2}\right)$, benzene, etc.). In this book, we shall study the electronic structure of these complexes by molecular orbital (MO) theory. We shall seek to establish the shape, the energetic ordering, and the electronic occupation of the MO; starting from this detailed description of the electronic structure, we shall consider problems of geometry and reactivity.

Certain important aspects of electronic structure can nevertheless be obtained from a far simpler description, which aims merely at providing a formal analysis of the electron distribution in the complex. Although much simpler and more limited in its applications, this approach to electronic structure turns out to be extremely useful, for at least two reasons:

1. It uses classical ideas and 'language' that are common to all chemists, such as electronegativity or Lewis structures for the ligands. It provides important information, such as the oxidation state (or number) of the metal in the complex, the number of electrons in the immediate environment of the metal, and what one normally calls the 'electronic configuration' of the complex.
2. In a way which can be a little surprising at first sight, it is very useful in the orbital approach when one wishes, for example, to know the number of electrons that must be placed in the complex's nonbonding MO.

There are two ways to obtain this formal distribution of the electrons (or electron count) in a complex. The first, based on a 'covalent' model of the metal-ligand bond, is mostly used in organometallic chemistry, that is, in complexes which possess one or more metal-carbon bonds. The second, based on an 'ionic' model of the metal-ligand bond in which the two electrons are automatically attributed to the ligand, is more frequently employed for inorganic complexes. In fact, the choice between the two methods is largely a matter of taste, as they lead, as we shall see, to identical conclusions.
${ }^{1}$ There is also a lone pair on the oxygen atom. We shall see later why CO binds preferentially through the carbon atom (§ 1.5.2.4 and Chapter 3, § 3.2.2).
${ }^{2}$ Complexes in which a dihydrogen molecule is bound to a transition metal were first characterized in the mid-1980s, and have since been extensively studied by both experimental and theoretical methods (Chapter 4, § 4.1.4).



### 1.1. Electron count in a complex: the covalent model

Consider a monometallic complex in which the transition metal M is bound to a certain number of ligands $(\mathrm{Lig})_{i}$, that may be either atoms or molecules. It is important to note that, in the covalent model, one always considers the ligands in their neutral form $\left(\mathrm{H}, \mathrm{Cl}, \mathrm{O}, \mathrm{CO}, \mathrm{CN}, \mathrm{PR}_{3}, \mathrm{CH}_{3}\right.$, etc.). Before making the formal electronic assignment for the complex, one must first categorize the ligands according to the nature of their electronic structure.

### 1.1.1. Ligand classification (L or X)

The main distinction is linked to the number of electrons that the ligand supplies to the metal's coordination sphere: if it supplies a pair of electrons, it is a ligand of type L, whereas a ligand that supplies just one electron is of type X. However, some ligands can supply more than two electrons to the metal. This notation, introduced by M. L. H. Green, is generalized to yield ligands of type $L_{\ell} X_{x}$.

### 1.1.1.1. L-type ligands

The simplest case concerns molecules which are coordinated to the metal through a lone pair located on one of their atoms (1-1). These molecules are L-type ligands, the metal-ligand bond being formed by the two electrons supplied by the ligand. Examples include amines $\mathrm{NR}_{3}$ and phosphines $\mathrm{PR}_{3}$ which contain a lone pair on the nitrogen or phosphorus atom, the water molecule or any ether $\left(\mathrm{OR}_{2}\right)$ which can bind to the metal through one of the lone pairs on the oxygen atom. Carbon monoxide is also an L-type ligand, due to the lone pair on the carbon atom. ${ }^{1}$





1-1

There are other cases in which the two electrons supplied by the ligand $L$ form a bond between two atoms of that ligand, rather than a lone pair. This can be a $\pi$-bond, as in the ethylene molecule, or, more surprisingly, a $\sigma$-bond, as in the dihydrogen molecule (1-2). ${ }^{2}$

In these examples, two atoms of the ligand are bound in an equivalent way to the metal centre. The hapticity of the ligand is said to be 2. This type of bond is indicated by the Greek letter $\eta$, the nomenclature used being $\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}$ or $\eta^{2}-\mathrm{H}_{2}$, respectively (1-2).

### 1.1.1.2. X-type ligands

These ligands supply only one electron to the metal's coordination sphere. As neutral entities, X-type ligands are radicals and the metalligand bond is formed by the unpaired electron of the ligand and a metal electron. Hydrogen $(\mathrm{H})$ is an X-type ligand, as are the halogens ( $\mathrm{F}, \mathrm{Cl}$, $\mathrm{Br}, \mathrm{I})$, alkyl radicals $\left(\mathrm{CR}_{3}\right)$, the amido $\left(\mathrm{NR}_{2}\right)$, alkoxy (OR), and cyano (CN) groups (1-3), etc.
-H






1-3

It should be noted that in some of the examples given above, the radical centre also possesses one or more lone pairs, so that one might have considered it to be an L-type ligand. However, the use of a lone pair for bond formation would lead to a complex with an unpaired electron on the metal (.L:-M). This electronic structure is less stable than that in which the unpaired electron and metal electron are paired to form the metal-ligand bond ( $: \mathrm{X}-:-\mathrm{M})$. It can be seen that in this case, all the electrons are paired, either as bonding pairs or as lone pairs.

### 1.1.1.3. Ligands of $\mathrm{L}_{\ell} \mathrm{X}_{x}$ type

In a more general notation, ligands can be represented as $\mathrm{L}_{\ell} \mathrm{X}_{x}$ when they use $\ell$ electron pairs and $x$ unpaired electrons to bind to the metal.

In the ground state, the oxygen atom possesses two unpaired electrons (1-4a). ${ }^{3}$ It is therefore a ligand of $\mathrm{X}_{2}$ type, which can bind to a transition metal to form an 'oxo' complex. The sulfido ( S ) and imido (N-R) (1-4a) ligands behave similarly. Atomic nitrogen, with three unpaired electrons, is an $X_{3}$ ligand ( $\mathbf{1 - 4 b}$ ), giving 'nitrido' complexes. In each case, one therefore considers all the unpaired electrons on the atom bound to the metal.


Conjugated polyenes constitute an important family of molecules which are ligands of $\mathrm{L}_{\ell} \mathrm{X}_{x}$ type; they form $\pi$-complexes with the metal, that is, complexes in which the $\pi$ system of the ligand interacts with the metal centre. Consider, for example, the cyclopentadienyl ligand

1-5

$1-6\left(\eta^{5}, \mathrm{~L}_{2} \mathrm{X}\right)$

Fe

1-7

$1-8\left(\eta^{3}, \mathrm{LX}\right)$

$\mathbf{1 - 1 0}\left(\eta^{4}, \mathrm{~L}_{2}\right)$

1-9 $\left(\eta^{1}, \mathrm{X}\right)$


1-11 $\left(\eta^{2}, \mathrm{~L}\right)$
$\mathrm{C}_{5} \mathrm{H}_{5}$ (also represented Cp ), whose Lewis structure (1-5) shows that the $\pi$ system contains five electrons (two $\pi$ bonds and one unpaired electron).

If this ligand is bound so that all five atoms are essentially at the same distance from the metal centre ( $\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}$ coordination), the five $\pi$ electrons are involved in the metal-ligand bonds, so that cyclopentadienyl is classified as an $\mathrm{L}_{2} \mathrm{X}$ ligand. Two graphical representations are therefore possible, depending on whether one gives a localized or delocalized description of the ligand's $\pi$ system (1-6).

Ferrocene is a particularly interesting $\pi$ complex $\left[\mathrm{Fe}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\right]$ (a 'sandwich' complex, in which an iron atom is placed between the planes of two cyclopentadienyl ligands, 1-7). At first sight, one could consider it either as a complex with two ligands, $\left[\mathrm{Fe}(\mathrm{Lig})_{2}\right]$ where $\mathrm{Lig}=\mathrm{C}_{5} \mathrm{H}_{5}$, or as one with 10 ligands, $\left[\mathrm{Fe}(\mathrm{Lig})_{10}\right]$, since the iron is bonded equivalently to all 10 carbon atoms. However, the L/X ligand classification shows us that each cyclopentadienyl ligand is of the $\mathrm{L}_{2} \mathrm{X}$ type, so ferrocene is therefore an $\left[\mathrm{FeL}_{4} \mathrm{X}_{2}\right]$ complex in which the iron must be considered as surrounded by six ligands, rather than two or ten. In fact, it is a pseudo-octahedral complex of $\left[\mathrm{Fe}(\mathrm{Lig})_{6}\right]$ type!

Two other coordination modes can be imagined for the cyclopentadienyl ligand, and they are indeed observed in some complexes. If only three $\pi$ electrons (a double bond and the unpaired electron) are supplied to the metal's coordination sphere, $\mathrm{C}_{5} \mathrm{H}_{5}$ acts as a ligand of LX type. In this case, only three carbon atoms are bound to the metal, and the coordination mode is $\eta^{3}-\mathrm{C}_{5} \mathrm{H}_{5}(\mathbf{1 - 8})$. Finally, the metal can bind just to the radical centre (X-type ligand), giving an $\eta^{1}-\mathrm{C}_{5} \mathrm{H}_{5}$ coordination (1-9). In this latter case, one can no longer describe it as a $\pi$ complex, since the metal centre interacts with only one of the ring carbon atoms, with which it forms a $\sigma$ bond.

This diversity of coordination behaviour is also found for other conjugated polyenes. Thus, butadiene can act as an $L_{2}$ ligand if the electrons of the two $\pi$ bonds are involved ( $\eta^{4}$-butadiene, 1-10) or as an L ligand involving a single $\pi$ bond ( $\eta^{2}$-butadiene, 1-11).

In the same way, benzene can bind in the $\eta^{6}\left(\mathrm{~L}_{3}\right.$ ligand, 1-12), $\eta^{4}\left(\mathrm{~L}_{2}\right.$ ligand, 1-13), or $\eta^{2}$ modes (Lligand, 1-14) (see Exercise 1.5). In the $\eta^{4}$ and $\eta^{2}$ coordination modes, the six carbon atoms become non-equivalent


1-12 $\left(\eta^{6}, \mathrm{~L}_{3}\right)$


1-13 $\left(\eta^{4}, \mathrm{~L}_{2}\right)$


1-14 ( $\left.\eta^{2}, \mathrm{~L}\right)$

This property is readily explained by molecular orbital theory: two electrons must be placed in two degenerate $\pi_{\mathrm{oo}}^{*}$ orbitals. The most favourable arrangement contains one electron in each orbital, with their spins parallel (a triplet state).


1-17a (L)


1-17b ( $\mathrm{X}_{2}$ )
(four or two, respectively, are bound to the metal), perturbing the $\pi$-electron conjugation. As a result, the ring becomes non-planar.

A slightly different case arises for ligands which can bind to a metal centre through several different sites without any conjugation of the electrons involved. These ligands are said to be polydentate (bidentate, tridentate, ...), in contrast to monodentate ligands such as $\mathrm{PR}_{3}, \mathrm{CR}_{3}$, etc. For example, 1,2-bis(dimethylephosphino)ethane is a bidentate ligand, since it can bind through its phosphino sites (1-15). As each of these has a lone pair, it behaves as an $\mathrm{L}_{2}$ ligand towards the metal. 1,2dioxyethane $\left(\mathrm{O}_{-} \mathrm{CH}_{2}-\mathrm{CH}_{2}-\mathrm{O}\right)$ is also a bidentate ligand (1-16), but each oxygen atom supplies only one electron to the metal (an $\mathrm{X}_{2}$ ligand).


1-15 ( $L_{2}$ )


1-16 ( $\mathrm{X}_{2}$ )

We end this section by discussing several ligands whose classification as L- or X-type can create difficulties.

The usual Lewis structure for the dioxygen molecule, $\mathrm{O}_{2}$, shows a double bond and two lone pairs on each oxygen atom. One might therefore conclude that it is an L-type ligand, and that the coordination would be either $\eta^{1}$ (through a lone pair) or $\eta^{2}$ (involving the $\pi$ bond). However, this Lewis structure, in which all the elctrons are paired, is not satisfactory since the magnetic moment measured experimentally shows that there are two unpaired electrons with parallel spin (the ground state is a triplet). ${ }^{4}$ This is why $\mathrm{O}_{2}$ behaves as an $\mathrm{X}_{2}$ ligand rather than an L ligand.

Carbene ligands, $\mathrm{CR}_{1} \mathrm{R}_{2}$, provide another example. These species contain two electrons on the carbon atom that do not participate in the formation of the C-R $\mathrm{R}_{1}$ and $\mathrm{C}-\mathrm{R}_{2}$ bonds. Depending on the nature of the $R_{1}$ and $R_{2}$ atoms or groups, the ground state is either diamagnetic, in which case the two nonbonding electrons are paired, forming a lone pair on the carbon atom (1-17a), or paramagnetic, in which case the two electrons are unpaired, giving a triplet state $\mathbf{1 - 1 7 b}$ ). In the first case, it is logical to consider the carbene as an L-type ligand, whereas it is an $\mathrm{X}_{2}$ ligand in the second case.

These two ways of describing a $\mathrm{CR}_{1} \mathrm{R}_{2}$ ligand are indeed used, and this distinction is at the origin of the two families of carbene complexes in organometallic chemistry: the Fischer-type ( L ) and the Schrock-type $\left(\mathrm{X}_{2}\right)$ carbenes. We shall return to this difference and offer an orbital interpretation (Chapter 4, § 4.3).


1-18


X ligand X ligand
1-19

5 There are two other transition series that correspond to the filling of the $4 f$ (lanthanides) and $5 f$ (actinides) sub-shells.

### 1.1.1.4. Bridging ligands

In bimetallic complexes, some ligands can be 'bridging', that is, bound simultaneously to the two metal centres. These cases are indicated by the nomenclature $\mu$. If one considers a bridging chlorine atom $\left(\mathrm{M}_{2}(\mu-\mathrm{Cl})\right.$, 1-18), it behaves as an X ligand towards the first metal centre, thanks to its unpaired electron, but as an L ligand towards the second, thanks to one of its lone pairs (the roles of the two metal centres can, of course, be interchanged). Overall, the chlorine atom is therefore an LX ligand; it supplies three electrons to the pair of metal centres. Other ligands in which an atom has an unpaired electron and at least one lone pair, such as $\mathrm{OR}, \mathrm{SR}, \mathrm{NR}_{2}, \mathrm{PR}_{2}$, etc. are analogous. A bridging oxygen atom is an X -type ligand towards each of the two metal centres, since it has two unpaired electrons (1-19), so it therefore acts as an $\mathrm{X}_{2}$ ligand overall.

### 1.1.2. Electron count and the 18 -electron rule

Once the nature of the ligands has been established, the second stage of our analysis of the electronic structure of transition metal complexes will require us to count the number of electrons around the metal and then to assign them, in a formal way, either to the metal or to the ligands. In what follows, we shall consider complexes written as $\left[\mathrm{ML}_{\ell} \mathrm{X}_{x}\right]^{q}$, in which the metal M is bound to $\ell$ ligands of L type and to $x$ ligands of X type, the overall charge being $q$.

### 1.1.2.1. Total number of electrons, the 18 -electron rule

Each ligand L supplies two electrons to the metal's environment, while each ligand X supplies a single electron. The total number of electrons supplied by the ligands is therefore equal to $2 \ell+x$. Only the valence electrons are considered for the transition metal, as, following the spirit of Lewis theory, we assume that core electrons play a negligible role. In what follows, we shall limit our analysis to transition elements corresponding to the progressive filling of the $3 d, 4 d$, and $5 d$ sub-shells (the transition metals of the $d$ block, see Table 1.1). The valenceelectron configuration of these elements is of the type $n d^{a}(n+1) s^{b}$, where $n$ equals 3,4 , or 5 , for the first, second, or third transition series, respectively. ${ }^{5}$ The metal therefore supplies $(a+b)$ electrons. We note that some authors do not consider zinc to be a transition element, as its $d$ sub-shell is full (valence-electron configuration $3 d^{10} 4 s^{2}$ ). This remark also applies to cadmium ( $\mathrm{Cd}, 4 \mathrm{~d}^{10} 5 s^{2}$ ) and to mercury ( $\mathrm{Hg}, 5 \mathrm{~d}^{10} 6 s^{2}$ ).

When we take account of the overall charge $q$ of the complex, the total number of valence electrons, $N_{\mathrm{t}}$, is:

$$
\begin{equation*}
N_{\mathrm{t}}=m+2 \ell+x-q \tag{1.1}
\end{equation*}
$$

Table 1.1. Electron configuration and number of valence electrons, $m$, for the d-block transition metals

| 1st series | Sc | Ti | V | Cr | Mn | Fe | Co | Ni | Cu | Zn |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
|  | $3 d^{1} 4 s^{2}$ | $3 d^{2} 4 s^{2}$ | $3 d^{3} 4 s^{2}$ | $3 d^{5} 4 s^{1}$ | $3 d^{5} 4 s^{2}$ | $3 d^{6} 4 s^{2}$ | $3 d^{7} 4 s^{2}$ | $3 d^{8} 4 s^{2}$ | $3 d^{10} 4 s^{1}$ | $3 d^{10} 4 s^{2}$ |
| 2nd series | Y | Zr | Nb | Mo | Tc | Ru | Rh | Pd | Ag | Cd |
|  | $4 d^{1} 5 s^{2}$ | $4 d^{2} 5 s^{2}$ | $4 d^{4} 5 s^{1}$ | $4 d^{5} 5 s^{1}$ | $4 d^{5} 5 s^{2}$ | $4 d^{7} 5 s^{1}$ | $4 d^{8} 5 s^{1}$ | $4 d^{10} 5 s^{0}$ | $4 d^{10} 5 s^{1}$ | $4 d^{10} 5 s^{2}$ |
| 3rd series | Lu | Hf | Ta | W | Re | Os | Ir | Pt | Au | Hg |
|  | $5 d^{1} 6 s^{2}$ | $5 d^{2} 6 s^{2}$ | $5 d^{3} 6 s^{2}$ | $5 d^{4} 6 s^{2}$ | $5 d^{5} 6 s^{2}$ | $5 d^{6} 6 s^{2}$ | $5 d^{7} 6 s^{2}$ | $5 d^{9} 6 s^{1}$ | $5 d^{10} 6 s^{1}$ | $5 d^{10} 6 s^{2}$ |
| $m$ | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 | 11 | 12 |

Several examples of the application of this rule are given below:

| Complex | $m$ | $2 \ell$ | $x$ | $q$ | $N_{t}$ |
| :--- | ---: | ---: | ---: | ---: | ---: |
| $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$ | 8 | 10 | 0 | 0 | 18 |
| $\left[\mathrm{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right]$ | 9 | 6 | 1 | 0 | 16 |
| $\left[\mathrm{Mn}(\mathrm{CO})_{6}\right]^{+}$ | 7 | 12 | 0 | +1 | 18 |
| $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$ | 10 | 0 | 5 | -3 | 18 |
| $\left[\mathrm{Zn}(\mathrm{Cl})_{4}\right]^{2-}$ | 12 | 0 | 4 | -2 | 18 |
| $\left[\mathrm{~V}(\mathrm{Cl})_{4}\right]$ | 5 | 0 | 4 | 0 | 9 |
| $\left[\mathrm{Cr}(\mathrm{CO})_{3}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\right]$ | 6 | 12 | 0 | 0 | 18 |
| $\left[\mathrm{Fe}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\right]$ | 8 | 8 | 2 | 0 | 18 |
| $\left[\mathrm{Cu}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\left(\mathrm{PMe}_{3}\right)\right]$ | 11 | 6 | 1 | 0 | 18 |
| $\left[\mathrm{Zr}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\left(\mathrm{CH}_{3}\right)\right]^{+}$ | 4 | 8 | 3 | +1 | 14 |
| $\left[\mathrm{Ti}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})_{3}\left(\mathrm{CH}_{3}\right)\right]$ | 4 | 4 | 4 | 0 | 12 |
| $\left[\mathrm{~W}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{CO})_{3}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$ | 6 | 12 | 0 | 0 | 18 |
| $\left[\mathrm{Ir}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})(\mathrm{H})_{2}\right]$ | 9 | 4 | 3 | 0 | 16 |
| $\left[\mathrm{Ni}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ | 10 | 12 | 0 | +2 | 20 |

By analogy with the octet rule, it has been proposed that a transition metal tends to be surrounded by the number of valence electrons equal to that of the following rare gas (electron configuration $n d^{10}(n+1) s^{2}$ $\left.(n+1) p^{6}\right)$. One thereby obtains the 18 -electron rule, for which we shall provide a first theoretical justification in this chapter (§ 1.6.3). However, in light of the examples given above, one must note that there are many exceptions to this rule; we shall analyse them in greater detail in the following chapters.

### 1.1.2.2. Oxidation state

In order to determine the oxidation state of the metal in the complex, one performs a fictitious dissociation of all the ligands, supposing that

Table 1.2. The Allred-Rochow electronegativity scale: (a) for the transition metals and (b) for the light elements
(a)

| Sc | Ti | V | Cr | Mn | Fe | Co | Ni | Cu | Zn |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1.20 | 1.32 | 1.45 | 1.56 | 1.60 | 1.64 | 1.70 | 1.75 | 1.75 | 1.66 |
| Y | Zr | Nb | Mo | Tc | Ru | Rh | Pd | Ag | Cd |
| 1.11 | 1.22 | 1.23 | 1.30 | 1.36 | 1.42 | 1.45 | 1.35 | 1.42 | 1.46 |
| Lu | Hf | Ta | W | Re | Os | Ir | Pt | Au | Hg |
| 1.14 | 1.23 | 1.33 | 1.40 | 1.46 | 1.52 | 1.55 | 1.44 | 1.41 | 1.44 |


each of them, either L or X, takes with it the electron pair that created the metal-ligand bond. The remaining charge on the metal after this decomposition is the oxidation state of the metal in the complex. This distribution of the electrons, which 'assigns' the bond pair to the ligand, can be partially justified when one notes that this latter is usually a more electronegative entity than is the transition metal (see Table 1.2, the Allred-Rochow electronegativity scale). The metal-ligand bonds are therefore polarized, and the electron pair is more strongly localized on the ligand than on the metal. To assign the two electrons of the bond just to the ligand is, however, a formal distribution, which exaggerates the tendency linked to the difference in electronegativity.

In the fictitious dissociation that we are considering, a ligand $L$ leaves with the two electrons that it had supplied, so the number of electrons on the metal is not changed in any way. However, an X-type ligand, which had supplied only a single electron to make the bond, leaves in its anionic form $\mathrm{X}^{-}$, carrying the two electrons from the bond with it. It therefore 'removes' an electron from the metal, that is, it oxidizes it by one unit. The result of this dissociation is therefore written:

$$
\begin{equation*}
\left[\mathrm{ML}_{\ell} \mathrm{X}_{x}\right]^{q} \rightarrow \ell \mathrm{~L}+x \mathrm{X}^{-}+\mathrm{M}^{(x+q)} \tag{1.2}
\end{equation*}
$$

The oxidation state (no) of the metal in the complex is therefore equal to the algebraic sum of the number of X-type ligands and the charge on the complex:

$$
\begin{equation*}
n o=x+q \tag{1.3}
\end{equation*}
$$

In a widely used notation to specify the oxidation state of the metal in a complex, the chemical symbol of the metal is followed by the oxidation state written in Roman letters ( $\mathrm{Mn}(\mathrm{I}), \mathrm{Fe}(\mathrm{II}), \mathrm{Cr}(\mathrm{III})$, etc.).

Examples

| Complex | $x$ | $q$ | no | Oxidation state |
| :--- | :---: | ---: | :---: | :---: |
| $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$ | 0 | 0 | 0 | $\mathrm{Fe}(0)$ |
| $\left[\mathrm{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right]$ | 1 | 0 | 1 | $\mathrm{Ir}(\mathrm{I})$ |
| $\left[\mathrm{Mn}(\mathrm{CO})_{6}\right]^{+}$ | 0 | +1 | 1 | $\mathrm{Mn}(\mathrm{I})$ |
| $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$ | 5 | -3 | 2 | $\mathrm{Ni}(\mathrm{II})$ |
| $\left[\mathrm{Zn}(\mathrm{Cl})_{4}\right]^{2-}$ | 4 | -2 | 2 | $\mathrm{Zn}(\mathrm{II})$ |
| $\left[\mathrm{V}(\mathrm{Cl})_{4}\right]$ | 4 | 0 | 4 | $\mathrm{~V}(\mathrm{IV})$ |
| $\left[\mathrm{Cr}(\mathrm{CO})_{3}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\right]$ | 0 | 0 | 0 | $\mathrm{Cr}(0)$ |
| $\left[\mathrm{Fe}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\right]$ | 2 | 0 | 2 | $\mathrm{Fe}(\mathrm{II})$ |
| $\left[\mathrm{Cu}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\left(\mathrm{PMe}_{3}\right)\right]$ | 1 | 0 | 1 | $\mathrm{Cu}(\mathrm{I})$ |
| $\left[\mathrm{Zr}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\left(\mathrm{CH}_{3}\right)\right]^{+}$ | 3 | +1 | 4 | $\mathrm{Zr}(\mathrm{IV})$ |
| $\left[\mathrm{Ti}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})_{3}\left(\mathrm{CH}_{3}\right)\right]$ | 4 | 0 | 4 | $\mathrm{Ti}(\mathrm{IV})$ |
| $\left[\mathrm{W}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{CO})_{3}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$ | 0 | 0 | 0 | $\mathrm{~W}(0)$ |
| $\left[\mathrm{Ir}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})(\mathrm{H})_{2}\right]$ | 3 | 0 | 3 | $\mathrm{Ir}(\mathrm{III})$ |
| $\left[\mathrm{Ni}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ | 0 | +2 | 2 | $\mathrm{Ni}(\mathrm{II})$ |

In bimetallic complexes, the oxidation state is calculated by supposing that the metal-metal bond(s), if any, is/ are broken homolytically. This procedure is justified by the fact that the electronegativities of the two metal centres are equal if they are identical, or similar in heteronuclear complexes (see Table 1.2(a)). The presence of one or more bonds between the metals therefore has no effect on their oxidation state. For example, the complex $\left[\mathrm{Mo}(\mathrm{Cl})_{2}\left(\mathrm{PR}_{3}\right)_{2}\right]_{2}$ can initially be considered to be decomposed into two monometallic neutral fragments $\left[\mathrm{Mo}\left(\mathrm{Cl}_{2}\right)_{2}\left(\mathrm{PR}_{3}\right)_{2}\right]$ in which the oxidation state of molybdenum is +2 .

In conclusion, we note that the oxidation state must not be equated to the real charge on the metal in the complex, as it is obtained from a formal distribution of the electrons between the metal and the ligands.

### 1.1.2.3. $d^{n}$ Configuration of a metal

The oxidation state of the metal, which supplies $m$ valence electrons, is equal to no after complex formation. The formal number of electrons remaining on the metal, $n$, is therefore given by the relationship:

$$
\begin{equation*}
n=m-n o \tag{1.4}
\end{equation*}
$$

We are considering here $n$ electrons which are not involved in the formation of metal-ligand bonds, in other words 'nonbonding' electrons. The electron configuration of the metal in the complex is represented as $d^{n}$.

## Examples

| Complex | no | $m$ | Configuration |
| :--- | ---: | :---: | :---: |
| $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$ | 0 | 8 | $d^{8}$ |
| $\left[\operatorname{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPH}_{3}\right)_{2}\right]$ | 1 | 9 | $d^{8}$ |
| $\left[\mathrm{Mn}(\mathrm{CO})_{6}\right]^{+}$ | 1 | 7 | $d^{6}$ |
| $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$ | 2 | 10 | $d^{8}$ |
| $\left[\mathrm{Zn}(\mathrm{Cl})_{4}\right]^{2-}$ | 2 | 12 | $d^{10}$ |
| $\left[\mathrm{~V}(\mathrm{Cl})_{4}\right]$ | 4 | 5 | $d^{1}$ |
| $\left[\mathrm{Cr}(\mathrm{CO})_{3}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\right]$ | 0 | 6 | $d^{6}$ |
| $\left[\mathrm{Fe}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\right]$ | 2 | 8 | $d^{6}$ |
| $\left[\mathrm{Cu}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\left(\mathrm{PMe}_{3}\right)\right]$ | 1 | 11 | $d^{10}$ |
| $\left[\mathrm{Zr}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\left(\mathrm{CH}_{3}\right)\right]^{+}$ | 4 | 4 | $d^{0}$ |
| $\left[\mathrm{Ti}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})_{3}\left(\mathrm{CH}_{3}\right)\right]$ | 4 | 4 | $d^{0}$ |
| $\left[\mathrm{~W}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{CO})_{3}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$ | 0 | 6 | $d^{6}$ |
| $\left[\mathrm{Ir}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})(\mathrm{H})_{2}\right]$ | 3 | 9 | $d^{6}$ |
| $\left[\mathrm{Ni}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ | +2 | 10 | $d^{8}$ |

This notation might seem surprising at first sight, as it implies that all the nonbonding electrons on the metal occupy $d$-type atomic orbitals (AO). Yet, for every metal except palladium, the $s$ orbital is at least partially occupied in the ground state of the isolated atom (see Table 1.1). A detailed study of the electronic structure of complexes, presented in Chapter 2, will show us that the nonbonding electrons on the metal do indeed occupy pure $d$-type orbitals, or molecular orbitals whose main component is a $d$-type atomic orbital.

### 1.2. An alternative model: the ionic model

There is a second method for counting the electrons in a complex and deducing the metal's oxidation state and electronic configuration. This is the ionic model, in which one supposes that a complex is formed by a metal centre and by ligands which always act as Lewis bases, supplying one (or several) pairs of electrons.

### 1.2.1. Lewis bases as ligands

In the covalent model, neutral ligands L (or $\mathrm{L}_{n}$ ) supply one (or $n$ ) electron pair(s) to the metal: for example, one for amines $\left(\mathrm{NR}_{3}\right)$, phosphines $\left(\mathrm{PR}_{3}\right)$, the carbonyl group (CO), and derivatives of ethylene ( $\mathrm{R}_{2} \mathrm{C}=\mathrm{CR}_{2}$ ), and three for benzene $\left(\mathrm{C}_{6} \mathrm{H}_{6}\right)$ in the $\eta^{6}$ coordination
mode. As these ligands already behave as Lewis bases in the covalent model, we shall continue to consider them in their neutral form L (or $\left.\mathrm{L}_{n}\right)$ in the ionic model.

However, an X-type ligand in the covalent model is a radical species which supplies only a single electron to the metal. To 'transform' it into a Lewis base, one must add an electron and therefore consider it in its anionic form $\mathrm{X}^{-}$. In this way, the radical ligands H (hydrogen), Cl (chlorine), and $\mathrm{CH}_{3}$ (methyl radical) of the covalent model become the $\mathrm{H}^{-}$(hydride), $\mathrm{Cl}^{-}$(chloride), and $\mathrm{CH}_{3}^{-}$(methyl anion) ligands in the ionic model. Analogously, $\mathrm{X}_{x}$ ligands in the covalent model, which have $x$ unpaired electrons, become $\mathrm{X}^{x-}$ ligands in the ionic model. For example, $\mathrm{O}\left(\mathrm{X}_{2}\right)$ and $\mathrm{N}\left(\mathrm{X}_{3}\right)$ are now described as $\mathrm{O}^{2-}$ and $\mathrm{N}^{3-}$. In general, one completes the ligand's valence-electron shell so that the octet rule is satisfied.

This is generalized for ligands of $\mathrm{L}_{\ell} \mathrm{X}_{x}$ type in the covalent model, which quite naturally become $\mathrm{L}_{\ell} \mathrm{X}_{x}^{x-}$ ligands in the ionic model. The cyclopentadienyl radical ( Cp ), a neutral species with five $\pi$ electrons (an $\mathrm{L}_{2} \mathrm{X}$ ligand, $\mathbf{1 - 5}$ ), is therefore considered in its monoanionic form ( $\mathrm{Cp}^{-}$with six $\pi$ electrons). Table 1.3 presents the numbers of electrons attributed to the principal ligands that have been considered so far in the covalent and ionic models.

The additional electron supplied to an X-type ligand to transform it into a Lewis base comes, of course, from the metal. The metalligand ensemble is therefore described as an $\mathrm{X}^{-}$ligand interacting with a metallic cation $\mathrm{M}^{+}$, thereby giving a purely ionic description of the metal-ligand bond. As a consequence, a complex which was written $\mathrm{ML}_{\ell} \mathrm{X}_{x}$ in the covalent model is represented, in the ionic model, as a metallic cation of charge $x$ bound to $(\ell+x)$ Lewis bases (1.5).

$$
\begin{equation*}
\left[\mathrm{ML}_{\ell} \mathrm{X}_{x}\right] \text { (covalent model) } \rightarrow\left[\mathrm{M}^{(x)+}(\mathrm{L})_{\ell}\left(\mathrm{X}^{-}\right)_{x}\right] \text { (ionic model). } \tag{1.5}
\end{equation*}
$$

If the complex has an overall charge $q$, the charge on the metallic centre in the ionic model becomes $(x+q)(1.6)$.

$$
\begin{equation*}
\left[\mathrm{ML}_{\ell} \mathrm{X}_{x}\right]^{q} \text { (covalent model) } \rightarrow\left[\mathrm{M}^{(x+q)+}(\mathrm{L})_{\ell}\left(\mathrm{X}^{-}\right)_{x}\right] \text { (ionic model). } \tag{1.6}
\end{equation*}
$$

This 'redistribution' of the electrons within the complex can be justified by the higher electronegativity of the ligands than of the metals (see Table 1.2): an X-type ligand 'attracts' the two electrons of the metalligand bond to itself, and becomes $\mathrm{X}^{-}$. In conclusion, we note that the name given to several complexes is directly linked to the ionic model. Thus, complexes with several H ligands $\left(\left[\mathrm{ReH}_{9}\right]^{2-}\right.$, for example) are called 'polyhydrides'.

Table 1.3. Number of electrons supplied by several common ligands according to the covalent and ionic models

| Covalent model |  | Ionic model |  |
| :---: | :---: | :---: | :---: |
| Ligand (type) | Number of electrons | Ligand | Number of electrons |
| $\mathrm{H}, \mathrm{Cl}, \mathrm{OR}, \mathrm{NR}_{2}, \mathrm{CR}_{3}, \mathrm{CN}$ ( X ligands) | 1 e | $\begin{aligned} & \mathrm{H}^{-}, \mathrm{Cl}^{-}, \mathrm{OR}^{-}, \mathrm{NR}_{2}^{-} \\ & \mathrm{CR}_{3}^{-}, \mathrm{CN}^{-} \end{aligned}$ | 2 e |
| $\begin{aligned} & \mathrm{CO}, \mathrm{NR}_{3}, \mathrm{PR}_{3}, \mathrm{H}_{2}, \\ & \mathrm{R}_{2} \mathrm{C}=\mathrm{CR}_{2} \text { (L ligands) } \end{aligned}$ | 2 e | $\begin{gathered} \mathrm{CO}, \mathrm{NR}_{3}, \mathrm{PR}_{3}, \mathrm{H}_{2}, \\ \mathrm{R}_{2} \mathrm{C}=\mathrm{CR}_{2} \end{gathered}$ | 2 e |
| O, S, NR ( $\mathrm{X}_{2}$ ligands) | 2 e | $\mathrm{O}^{2-}, \mathrm{S}^{2-}, \mathrm{NR}^{2-}$ | 4 e |
|  | 4 e |  | 4 e |
| $\eta^{4}$-diene ( $\mathrm{L}_{2}$ ligand) |  | $\eta{ }^{4}$-diene |  |
|  | 5 e |  | 6 e |
| $\eta^{5}-\mathrm{Cp}\left(\mathrm{L}_{2} \mathrm{X}\right.$ ligand) |  | $\eta^{5}-\mathrm{Cp}^{-}$ |  |
|  | 6 e |  | 6 e |
| $\eta^{6}$-arene ( $\mathrm{L}_{3}$ ligand) |  | $\eta^{6}$-arene |  |
| $\mu-\mathrm{Cl}$ (LX ligand) | 3 e | $\mu-\mathrm{Cl}^{-}$ | 4 e |
| $\mu-\mathrm{O}\left(\mathrm{X}_{2}\right.$ ligand $)$ | 2 e | $\mu-\mathrm{O}^{2-}$ | 4 e |

### 1.2.2. Equivalence of the covalent and ionic models: examples

### 1.2.2.1. Oxidation state and $d^{n}$ electronic configuration

In the covalent model, the oxidation state of the metal, no, is equal to the charge left on the metal after having carried out a fictitious dissociation of the complex in which all the ligands take the two bonding electrons with them ( $\left(\right.$ 1.1.2.2). For a complex whose general formula is $\left[\mathrm{ML}_{\ell} \mathrm{X}_{x}\right]^{q}$, one therefore obtains no $=x+q$ (see equations (1.2) and (1.3)). In the ionic formulation of this same complex, (see equation (1.6)), the charge on the metal is just equal to $x+q$, so the ionic and covalent models lead to the same oxidation state no for the metal. It follows that the same electronic configuration $d^{n}$ is obtained by the two models, since $n$ is equal to the number of valence electrons on the metal $(m)$, minus its oxidation state no (equation (1.4)).

## Examples

| Covalent model | Ionic model |
| :--- | :--- |
| $\left[\mathrm{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right]$ |  |
| $\left[\mathrm{Ir}(\mathrm{L})_{3}(\mathrm{X})\right]$ type | $\left.\left(\mathrm{Ir}^{+}\right)(\mathrm{CO})\left(\mathrm{Cl}^{-}\right)\left(\mathrm{PPh}_{3}\right)_{2}\right]$ |
| $x=1 ; q=0 \Rightarrow n o=+1$ | $n o=+1$ |
| $m=9 \Rightarrow n=9-1 \Rightarrow d^{8}$ | $m=9 \Rightarrow \mathrm{Ir}^{+}: d^{8}$ |
| $\left[\mathrm{Fe}\left(\eta^{5}-\mathrm{Cp}\right)_{2}\right]$ |  |
| $\left[\mathrm{Fe}\left(\mathrm{L}_{2} \mathrm{X}\right)_{2}\right]$ type | $\left[\left(\mathrm{Fe}^{2+}\right)\left(\mathrm{Cp}^{-}\right)_{2}\right]$ |
| $x=2 ; q=0 \Rightarrow n o=+2$ | $n o=+2$ |
| $m=8 \Rightarrow n=8-2 \Rightarrow d^{6}$ | $m=8 \Rightarrow \mathrm{Fe}^{2+}: d^{6}$ |
| $\left[\mathrm{Mn}(\mathrm{CO})_{6}\right]^{+}$ |  |
| $\left[\mathrm{Mn}(\mathrm{L})_{6}\right]^{+}$type | $\left[\left(\mathrm{Mn}^{+}\right)(\mathrm{CO})_{6}\right]$ |
| $x=0 ; q=+1 \Rightarrow n o=+1$ | $n o=+1$ |
| $m=7 \Rightarrow n=7-1 \Rightarrow d^{6}$ | $m=7 \Rightarrow \mathrm{Mn}^{+}: d^{6}$ |
| $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$ |  |
| $\left[\mathrm{Ni}(\mathrm{X})_{5}\right]^{3-}$ type | $\left[\left(\mathrm{Ni}^{2+}\right)\left(\mathrm{CN}^{-}\right)_{5}\right]$ |
| $x=5 ; q=-3 \Rightarrow n o=+2$ | $n o=+2$ |
| $m=10 \Rightarrow n=10-2 \Rightarrow d^{8}$ | $m=10 \Rightarrow \mathrm{Ni}^{2+}: d^{8}$ |

### 1.2.2.2. Total number of electrons

The equivalence of the two models for the calculation of the total number of electrons in a complex $\left(N_{t}\right)$ is shown by a further look at the four examples above.

| Covalent model |  | Ionic model |  |
| :---: | :---: | :---: | :---: |
| $\left[\operatorname{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right]$ |  |  |  |
| $\left[\operatorname{Ir}(\mathrm{L})_{3}(\mathrm{X})\right]$ type |  | $\left[\left(\mathrm{Ir}^{+}\right)(\mathrm{C}\right.$ | $\left.)\left(\mathrm{PPh}_{3}\right)_{2}\right]$ |
| Ir | 9 e | $\mathrm{Ir}^{+}$ | 8 e |
| 1 CO | 2 e | 1 CO | 2e |
| $2 \mathrm{PPh}_{3}$ | 4 e | $2 \mathrm{PPh}_{3}$ | 4 e |
| 1 Cl | 1 e | $1 \mathrm{Cl}^{-}$ | 2e |
| $N_{\text {t }}$ | 16e | $N_{\text {t }}$ | 16 e |
| [Fe( $\left.\left.\eta^{5}-\mathrm{CP}\right)_{2}\right]$ |  |  |  |
| $\left[\mathrm{Fe}\left(\mathrm{L}_{2} \mathrm{X}\right)_{2}\right]$ type |  | $\left[\left(\mathrm{Fe}^{2+}\right)\right.$ |  |
| Fe | 8 e | $\mathrm{Fe}^{2+}$ | 6 e |
| 2 Cp | 10 e | $2 \mathrm{Cp}^{-}$ | 12 e |
| $N_{\text {t }}$ | $18 e$ | $N_{\text {t }}$ | $18 e$ |
| $\left[\mathrm{Mn}(\mathrm{CO})_{6}\right]^{+}$ |  |  |  |
| $\left[\mathrm{Mn}(\mathrm{L})_{6}\right]^{+}$type |  | $\left[\left(\mathrm{Mn}^{+}\right)\right.$ |  |
| Mn | 7 e | $\mathrm{Mn}^{+}$ | 6 e |
| 6 CO | 12e | 6 CO | 12 e |
| Charge | $-1 \mathrm{e}$ | $N_{\text {t }}$ | $18 e$ |
| $N_{\text {t }}$ | $18 e$ |  |  |
| $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$ |  |  |  |
| $\left[\mathrm{Ni}(\mathrm{X})_{5}\right]^{3-}$ type |  | $\left[\left(\mathrm{Ni}^{2+}\right)\right.$ |  |
| Ni | 10e | $\mathrm{Ni}^{2+}$ | 8 e |
| 5 CN | 5 e | $5 \mathrm{CN}^{-}$ | 10e |
| Charge | 3 e | $N_{\text {t }}$ | $18 e$ |
| $N_{\text {t }}$ | $18 e$ |  |  |

Figure 1.1. Interaction diagram for two orbitals with the same energy.

### 1.3. Principles of orbital interactions

When we use molecular orbital (MO) theory, the term 'orbital structure' of a complex, or of any molecule, means the shape and the energetic order of the MO. Usually, these orbitals are expressed as Linear Combinations of Atomic Orbitals (LCAO) of the different atoms that make up the system being studied. The shape of an MO is determined by the relative magnitudes and the signs of the different coefficients. The electronic structure is then obtained by placing electrons in these orbitals, filling first those which are lowest in energy.

To construct the MO, it is often advantageous to decompose the molecular system being studied into two simpler sub-systems whose orbitals, either atomic or molecular, are already known. The MO of the complete system are then obtained by allowing the orbitals of the two fragments to interact. In this paragraph, we shall remind the reader of the principal rules which control the interaction between two orbitals on two fragments. For simplicity, we shall treat atomic orbitals, but this limitation will not affect the general nature of our conclusions in any way.

### 1.3.1. Interaction between two orbitals with the same energy

Consider, for example, the interaction between two identical orbitals of $s$ type, $\chi_{1}$ and $\chi_{2}$ (Figure 1.1).

The interaction produces a bonding ( $\phi_{+}$) and an antibonding MO ( $\phi_{-}$). The first is the in-phase combination of the two orbitals $\chi_{1}$ and $\chi_{2}$ (coefficients with the same sign), while the second is the out-of-phase combination (coefficients with opposite signs) of these same orbitals. In each MO, the coefficients of $\chi_{1}$ and $\chi_{2}$ have the same magnitude, since the interacting orbitals are identical.

In energy terms, the bonding MO is lower in energy than the initial AO, but the antibonding MO is higher. It is important to notice that the destabilization of the antibonding level $\left(\Delta E^{-}\right)$is larger than the stabilization of the bonding level $\left(\Delta E^{+}\right)$. It can be shown that these

${ }^{6}$ The overlap $S_{i j}$ between two orbitals $\phi_{i}$ and $\phi_{j}$ is equal to the integral evaluated over all space of the product of the functions $\phi_{i}^{*}$ (the complex conjugate function of $\phi_{i}$ ) and $\phi_{j}: S_{i j}=\left\langle\phi_{i}^{*} \mid \phi_{i}\right\rangle$. For real functions, the integral of the product of the two functions is evaluated over all space.

7 We assume here that the total electronic energy is equal to the sum of the individual electronic energies. This relationship, which has the advantage of being simple, is obtained when the electronic Hamiltonian is written as a sum of monoelectronic Hamiltonians, as in the Hückel and extended Hückel methods. This approximate formula has, of course, limited application, but it is acceptable for a qualitative analysis of orbital interactions. Further details may be found in Structure électronique des molécules, by Y. Jean and F. Volatron, Volume 2, Chapter 13, Dunod, Paris (2003).
two quantities are proportional to the overlap $S$ between the interacting orbitals. ${ }^{6}$ Therefore, as this overlap increases, the stabilization of the bonding MO and the destabilization of the antibonding MO both become larger.

We shall often be interested later in this book by interactions involving either two or four electrons. In the first case (1-20), after the interaction the two electrons both occupy the bonding MO, producing a stabilization of the electronic energy equal to $2 \Delta E^{+}$. ${ }^{7}$ We deduce that the stabilization associated with a two-electron interaction between orbitals of the same energy is proportional to the overlap $S$.


In the case of a four-electron interaction (1-21), both the bonding and antibonding orbitals are doubly occupied. Since $\Delta E^{-}$is larger than $\Delta E^{+}$, the four-electron interaction is destabilizing, and it can be shown that the destabilization is proportional to the square of the overlap, $S^{2}$.

### 1.3.2. Interaction between two orbitals with different energies

We now consider the more general case, where the two orbitals $\chi_{1}$ and $\chi_{2}$, have different energies $\left(\varepsilon_{1}<\varepsilon_{2}\right.$, Figure 1.2). Their interaction leads to the formation of a bonding orbital $\left(\phi_{+}\right)$, lower in energy than the lowest orbital $\left(\chi_{1}\right)$, and an antibonding orbital $\left(\phi_{-}\right)$, higher in energy than the highest orbital $\left(\chi_{2}\right)$. As in the preceding case, the stabilization $\left(\Delta E^{+}\right)$of the bonding orbital, compared to the energy of $\chi_{1}$, is smaller than the destabilization $\left(\Delta E^{-}\right)$of the antibonding orbital compared to the energy of $\chi_{2}$ (Figure 1.2). It can be shown that these two quantities are both proportional to the square of the overlap between the orbitals and inversely proportional to their energy difference $(\Delta \varepsilon)$, that is, proportional to $S^{2} / \Delta \varepsilon$. A strong interaction therefore requires both a good overlap between the orbitals and a small energy difference between them.

## Comment

This formula is approximate and cannot be used when the two orbitals are too close in energy. It is clear that the expression tends to infinity as $\Delta \varepsilon$ tends


Figure 1.2. Interaction diagram for two orbitals with different energies.


1-22


1-23
towards zero (orbitals of the same energy). For orbitals whose energies are only slightly different, it is safer to use the result from the preceding paragraph (proportional to $S$ ). An example will be discussed in Chapter 4, § 4.1.3.

As far as the coefficients are concerned, the bonding orbital $\left(\phi_{+}\right)$ is concentrated on the centre (or the fragment) that has the lowestenergy orbital ( $\chi_{1}$ here), whereas the opposite polarization is found for the antibonding orbital $\left(\phi_{-}\right)$, where the coefficient is larger for $\chi_{2}$ (Figure 1.2). From a chemical viewpoint, this means that the bonding MO is mainly based on the more electronegative centre (or fragment), but the antibonding MO on the less electronegative centre (or fragment).

If we consider a two-electron interaction between doubly occupied $\chi_{1}$ and empty $\chi_{2}(\mathbf{1 - 2 2})$, the two electrons are stabilized by $2 \Delta E^{+}$. The stabilization associated with a two-electron interaction between two orbitals of different energy is therefore proportional to the square of the overlap and inversely proportional to the energy difference between the two orbitals, that is, proportional to $S^{2} / \Delta \varepsilon$. However, a four-electron interaction is destabilizing, since $\Delta E^{-}$is larger than $\Delta E^{+}$(1-23). It can be shown that this four-electron destabilization is proportional to the square of the overlap, $S^{2}$.

The two-orbital interaction diagrams (1-20 to $\mathbf{1 - 2 3})$ enable us to establish a link between the idea of a bonding pair in the Lewis sense and the MO description. The bonding pair corresponds to double occupation of the bonding MO with the antibonding MO empty. There is thus a single bond in $\mathrm{H}_{2}$ (identical orbitals, 1-20) and in $\mathrm{HeH}^{+}$(different orbitals, 1-22). However, if four electrons are involved, the antibonding orbital is doubly occupied and no chemical bond exists between the two atoms. This is the situation in $\mathrm{He}_{2}$, for example (identical orbitals, 1-21), and $\mathrm{HeH}^{-}$(different orbitals, 1-23), species where the two atoms remain separate.

### 1.3.3. The role of symmetry

The interaction between two orbitals $\chi_{1}$ and $\chi_{2}$ leads to a stabilization (destabilization) of the bonding (antibonding) MO, proportional to the overlap if the orbitals have the same energy but to $S^{2} / \Delta \varepsilon$ if their energies are different. In both cases, there is clearly no interaction if the overlap is zero. Now $S$ is equal to the integral over all space of the product of the functions $\chi_{1}^{*}$ and $\chi_{2}$. In order for this integral to be non-zero, these two functions must be bases for the same irreducible representation of the molecular symmetry group, or, in simpler terms, they must have the same symmetry (Chapter $6, \S 6.5 .1$ ). If they have different symmetries, the integral is exactly equal to zero, and one says that the overlap is zero by symmetry.


Figure 1.3. Construction of the $\sigma_{\mathrm{CC}}$ and $\pi_{\mathrm{CC}}$ MO in ethylene from $n_{\sigma}$ and $n_{p}$ orbitals on each $\mathrm{CH}_{2}$ fragment.

In the general case, where two fragments each with several orbitals interact, this comment allows us to simplify the interaction diagrams very considerably: only orbitals of the same symmetry interact.

### 1.3.4. $\sigma$ and $\pi$ interactions

Two types of interactions are often distinguished: $\sigma$ interactions, which concern an axial orbital overlap, and $\pi$ interactions, where the orbital overlap occurs laterally, or 'sideways'. These two types of overlap are illustrated in 1-24 and 1-25, respectively, for two $p$ orbitals whose axes of revolution are either co-linear (axial overlap) or parallel (sideways overlap). Notice that another way to characterize $\pi$ interactions is to observe that the orbitals involved share a common nodal plane ( $P$, 1-25).

In general, $\sigma$ interactions are stronger than $\pi$ interactions, since axial overlap is more efficient than sideways overlap. The energy separation between the resulting orbitals is therefore larger for $\sigma$ (bonding) and $\sigma^{*}$ (antibonding) MO than for the $\pi$ and $\pi^{*}$ MO.

The ethylene molecule provides a typical example. The construction of the $\sigma_{\mathrm{CC}}$ and $\pi_{\mathrm{CC}} \mathrm{MO}$ from nonbonding orbitals (represented by $n_{\sigma}$ and $n_{p}$ ) of the $\mathrm{CH}_{2}$ fragments is presented in Figure 1.3: the order of the four resulting MO, in terms of increasing energy, is $\sigma_{\mathrm{CC}}<\pi_{\mathrm{CC}}<$ $\pi_{\mathrm{CC}}^{*}<\sigma_{\mathrm{CC}}^{*}$.

### 1.4. Metal orbitals

In the case of monometallic transition metal complexes, it seems quite natural to construct the MO by allowing the orbitals on the metal centre to interact with those on the ligands. We are now going to examine just which orbitals one must consider on the metal (§ 1.4) and on the ligands (§ 1.5) so as to obtain, after interaction, a satisfactory description of the orbital structure of the complex.

For the metal centre, the atomic orbitals ( AO ) describing the core electrons will not be considered for the construction of the complex's MO. This approximation can be justified by noting that the amplitude of these orbitals is significant only close to the nucleus, so they can therefore play only a negligible role in bond formation. One must, however, consider the valence AO that are occupied in the ground state of the isolated atom ( $n d$ and $(n+1) s)$, see Table 1.1), together with the $(n+1) p$ orbitals, which, even though they are empty in the isolated atom, do contribute to bond formation in the complexes of transition metals. There are, therefore, nine atomic orbitals in all which participate on the metal, five $d$-type orbitals, one $s$-type, and three $p$-type orbitals.
${ }^{8}$ Y. Jean and F. Volatron in An Introduction to Molecular Orbitals, Oxford University Press, NY. (1993). Chapter 2.
${ }^{9}$ In these expressions, the angular part is expressed using the ratios $x / r, y / r$, and $z / r$ rather than the spherical coordinates $(r, \theta$, and $\phi)$. This transformation enables us to make the link between the analytical expression of the orbital and the name which is attributed to it.

### 1.4.1. Description of the valence orbitals

For the $s$ and $p$ orbitals, we shall use the usual conventional representation ${ }^{8}(\mathbf{1 - 2 6})$ which takes their essential features into account:

1. The spherical symmetry of the $s$ orbital.
2. The existence of an axis of revolution for the $p_{x}, p_{y}$, and $p_{z}$ orbitals (the $O x, O y$, or $O z$ directions, respectively) and of a nodal plane perpendicular to this axis (the $y O z, x O z$, or $x O y$ planes, respectively), that is, a plane in which the orbital amplitude is zero. The $p$ orbitals change sign on crossing the nodal plane, which is why they are represented by two 'lobes', one grey (positive amplitude), the other white (negative amplitude).


It should be noted that the representation of the orbital whose axis of revolution is perpendicular to the plane of the page ( $p_{x}, \mathbf{1 - 2 6}$ ) poses a special problem. This function is exactly zero in this plane (the nodal plane $y O z$ ), the positive lobe being directed towards the reader and the negative lobe away from him or her. When one takes account of the symmetry of revolution around the $O x$ axis, the intersections of these lobes with planes parallel to the nodal plane, either 'above' or 'below' it, are circles. The conventional representation of this orbital shows two offset circles, which represent each lobe seen in perspective.

The presence of five valence $d$-type orbitals is, of course, the chief characteristic of the metals of the first three transition series. For hydrogenoïd atoms (those with only one electron but a nuclear charge equal to $+Z$ ), exact analytical solutions to the Schrödinger equation can be obtained (which is not the case for polyelectronic atoms in general). The expressions for the $3 d$ orbitals are given below (formulae (1.7)-(1.12)), where both the radial $\left(R_{3,2}(r)\right)$ and angular parts are normalized. ${ }^{9}$ Analogous expressions are obtained for the $4 d$ and $5 d$ orbitals of the hydrogenoïd atoms, only the radial part of the functions $\left(R_{4,2}\right.$
and $R_{5,2}$ ) being modified.

$$
\begin{align*}
3 d_{x y} & =R_{3,2}(r) \sqrt{\frac{60}{16 \pi}} \frac{x y}{r^{2}}  \tag{1.7}\\
3 d_{x z} & =R_{3,2}(r) \sqrt{\frac{60}{16 \pi}} \frac{x z}{r^{2}}  \tag{1.8}\\
3 d_{y z} & =R_{3,2}(r) \sqrt{\frac{60}{16 \pi}} \frac{y z}{r^{2}}  \tag{1.9}\\
3 d_{x^{2}-y^{2}} & =R_{3,2}(r) \sqrt{\frac{15}{16 \pi}} \frac{x^{2}-y^{2}}{r^{2}}  \tag{1.10}\\
3 d_{z^{2}} & =R_{3,2}(r) \sqrt{\frac{5}{16 \pi}} \frac{2 z^{2}-x^{2}-y^{2}}{r^{2}}, \tag{1.11}
\end{align*}
$$

where

$$
\begin{equation*}
R_{3,2}(r)=\frac{4}{81 \sqrt{30}} \sqrt{\frac{Z^{3}}{a_{0}^{3}}}\left(\frac{Z r}{a_{0}}\right) \exp \left(-\frac{Z r}{3 a_{0}}\right)^{2} . \tag{1.12}
\end{equation*}
$$

In this last expression (1.12), $a_{0}$ is the Bohr radius, equal to $0.529 \AA$, and $Z$ is the nuclear charge. To what extent are these hydrogenoïd orbitals suitable to describe the $d$ orbitals of transition metals? In polyelectronic atoms, it is only the radial part of the orbitals that is different from hydrogenoïd orbitals; it is modified to take account of the charge on the nucleus and the screening effect created by the other electrons. Since the angular part of the orbitals is conserved, the expressions that are obtained for the $3 d$ orbitals of hydrogenoìd atoms enable us to analyse the symmetry properties of the $d$ orbitals of all the transition metals.

We note first that the names given to these orbitals $\left(d_{x y}, d_{x z}, d_{y z}, d_{x^{2}-y^{2}}, d_{z^{2}}\right)$ are directly related to the formulae of their angular parts. At a given distance $r$ from the nucleus, the amplitude of the $d_{x z}$ orbital is directly proportional to the product of the $x$ and $z$ coordinates for that point (formula (1.8)). The same applies for the $d_{y z}, d_{x y}$, and $d_{x^{2}-y^{2}}$ orbitals. But the $d_{z^{2}}$ orbital is a special case. Its name suggests that it is concentrated wholly along the $z$-axis. But in fact, this orbital also has a small amplitude, of opposite sign, in the $x y$ plane, so according to formula (1.11), it would be more logical to call it $d_{2 z^{2}-\left(x^{2}+y^{2}\right)}$.

It is important to define carefully the graphical representations that we shall use for these orbitals throughout this book. They show the orbitals' symmetry properties, the regions of space where their amplitudes are largest and where they are zero (nodal surfaces), all important aspects for our subsequent analysis of interactions between the $d$ orbitals


1-27


1-28


1-29


1-31
and orbitals on the ligands. Consider the $d_{y z}$ orbital as an example. Its analytical expression (formula (1.9) shows that its amplitude is zero if $y=0$ (i.e. at all points in the $x z$ plane) and if $z=0$ ( $x y$ plane): $x z$ and $x y$ are therefore two nodal planes for the $d_{y z}$ orbital. In contrast, the amplitude is greatest along the bisectors of the $y$-and $z$-axes. Finally, it is positive where $y$ and $z$ have the same sign, but negative otherwise. All of these properties are clearly shown by the graphical representation 1-27.

The $d_{x y}$ and $d_{x z}$ orbitals (formulae (1.7) and (1.8)) may be obtained from the $d_{y z}$ orbital by a rotation of $90^{\circ}$ around the $y$-and $z$-axes, respectively. They have analogous symmetry properties, with two nodal planes ( $x z$ and $y z$ for $d_{x y}, x y$ and $y z$ for $d_{x z}$ ), a maximum amplitude along the bisectors of the $(x, y)$ or $(x, z)$ axes and alternating signs for the lobes. Their graphical representation poses the same problem as that already met for the $p_{x}$ orbital (1-26), since the plane of the page is one of the nodal planes. In the same way as before, we represent the intersection of the lobes with planes parallel to the plane of the page $(y z)$, placed either in front or behind, with the back part of the orbital being partly hidden by the front part (1-28 and 1-29).

The $d_{x^{2}-y^{2}}$ orbital (formula ((1.10)) has its maximum amplitude along the $x$ - and $y$-axes, and it also possesses two nodal planes which are the planes bisecting the $x$ - and $y$-axes (1-30a). An alternative representation of this orbital is given in $\mathbf{1 - 3 0 b}$, where the $x$-axis is perpendicular to the plane of the page. The lobes directed along this axis are now represented by two offset circles.


1-30a


1-30b

The shape of the $d_{z^{2}}$ orbital is very different from those we have already seen. Its analytical expression (formula (1.11)) shows that its maximal amplitude lies along the $z$-axis, and that it is positive for both positive and negative $z$. But it is negative in the $x y$ plane $(z=0)$, and this change in sign implies the existence of a nodal surface. The equation of this surface, $z^{2}=\left(x^{2}+y^{2}\right) / 2$ from formula (1.11), defines a cone whose apex angle, $\theta$, is equal to $109.5^{\circ}$ (the tetrahedral angle). All of these properties are reproduced by the conventional representation given in 1-31.

We close this section by noting that the sign in the analytical expressions of the orbitals is arbitrary. The same remark applies for

Table 1.4. Energies (in eV ) of the $s$ and $d$ orbitals for d-block transition elements obtained from spectroscopic data.

| 1st series | Sc | Ti | V | Cr | Mn | Fe | Co | Ni | Cu | Zn |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\varepsilon_{3 d}$ | -7.92 | -9.22 | -10.11 | -10.74 | -11.14 | -11.65 | -12.12 | -12.92 | -13.46 | -17.29 |
| $\varepsilon_{4 s}$ | -6.60 | -7.11 | -7.32 | -7.45 | -7.83 | -7.90 | -8.09 | -8.22 | -8.42 | -9.39 |
| 2nd series | Y | Zr | Nb | Mo | Tc | Ru | Rh | Pd | Ag | Cd |
| $\varepsilon_{4 d}$ | -6.48 | -8.30 | -8.85 | -9.14 | -9.25 | -9.31 | -9.45 | -9.58 | -12.77 | -17.85 |
| $\varepsilon_{5 s}$ | -6.70 | -7.31 | -7.22 | -7.24 | -7.21 | -7.12 | -7.28 | -7.43 | -7.57 | -8.99 |
| 3rd series | Lu | Hf | Ta | W | Re | Os | Ir | Pt | Au | Hg |
| $\varepsilon_{5 d}$ | -5.28 | -6.13 | -7.58 | -8.76 | -9.70 | -10.00 | -10.21 | -10.37 | -11.85 | -15.58 |
| $\varepsilon_{6 s}$ | -7.04 | -7.52 | -8.45 | -8.51 | -8.76 | -8.81 | -8.83 | -8.75 | -9.22 | -10.43 |

${ }^{10}$ J. B. Mann, T. L. Meek, E. T. Knight, J. F. Capitani, L. C. Allen J. Amer. Chem. Soc. 122, 5132 (2000). The calculations use the ionization potential of the atom and the energy of the atom in its ground state, as well as the energy of the cation formed by removal of an electron from the $s$ orbital or from a $d$ orbital.
their graphical representations. This means that one can change all the signs in the representation of an orbital, for example, representing the $d_{x^{2}-y^{2}}$ orbital by negative (white) lobes along the $x$-axis and positive (grey) lobes along the $y$-axis. All the orbital's properties, such as the regions of maximum amplitude, the changes in sign between neighbouring lobes or the nodal surfaces, are retained in this new representation.

## Notation

In the interest of simplification, the five $d$ orbitals are often written $x y, x z$, $y z, x^{2}-y^{2}$, and $z^{2}$. This is the notation that we shall use henceforth.

### 1.4.2. Orbital energies

It is possible to determine the energy of the $n d$ and $(n+1) s$ orbitals for transition metals in their ground-state electronic configuration $n d^{a}$ $(n+1) s^{b}$ from spectroscopic data. ${ }^{10}$

The values that are obtained are presented in Table 1.4. They invite several comments which will be helpful when we come to the construction of diagrams for the interaction between metal and ligand orbitals. On moving from left to right in a given series, the energy of both the $s$ and $d$ orbitals decreases (becoming more negative). This decrease in orbital energy arises from the increase in nuclear charge which strengthens the interaction between the nucleus and the electrons. The variation is nonetheless less pronounced for the $s$ orbitals than for the $d$ orbitals, because the $(n+1) s$ electrons are strongly shielded by the $n d$ electrons. As a consequence, the effective charge $Z_{(n+1) s}^{*}$
(the nuclear charge $Z$ reduced by the screening $\sigma_{(n+1) s}$ ) experienced by the $s$ electrons varies little from one element to the next: the increase of the nuclear charge by one unit is largely cancelled by the presence of an additional $d$ electron with its screening effect. However, the $d$ electrons are only weakly screened by the $s$ electrons, so the effective charge $Z_{n d}^{*}$ increases by close to one unit from one element to the next, leading to a substantial stabilization of the energy of the $d$ orbitals. When we consider the variation within a group, the energetic ordering of the $d$ orbitals in the first four columns is $\varepsilon_{3 d}<\varepsilon_{4 d}<\varepsilon_{5 d}$, but there is an inversion of the $4 d$ and $5 d$ levels for the four following groups $\left(\varepsilon_{3 d}<\varepsilon_{4 d}>\varepsilon_{5 d}\right)$. For all the elements except four ( $\mathrm{Y}, \mathrm{Lu}, \mathrm{Hf}$, and Ta ), the nd orbital is lower in energy than the $(n+1)$ s orbital. The $(n+1) p$ orbital is always higher in energy than the $(n+1) s$, as it is everywhere in the periodic table. For the great majority of the d-block transition metals, the orbital energy ordering is therefore, $\varepsilon_{n d}<\varepsilon_{(n+1) s}<\varepsilon_{(n+1) p}$.

### 1.5. Ligand orbitals

It is not possible to define a single set of orbitals that can be used to describe the interactions with the metal for any type of ligand. Two conditions must be met: the ligand orbitals must be close in energy to those on the metal, and their overlap must also be substantial (§ 1.3.2). Depending on the nature of the ligand, one or several orbitals may satisfy these criteria.

### 1.5.1. A single ligand orbital: $\sigma$ interactions

The case where it is clearest that only one orbital need be considered involves the ligand H , since it possesses only one valence orbital, $1_{s_{\mathrm{H}}}$. This orbital, which contains one electron (an X-type ligand), can be used to form a $\sigma_{\mathrm{M}-\mathrm{H}}$ bond by combination with a metal orbital such as the $z^{2}$ orbital (1-32). ${ }^{11}$

For certain more complicated ligands, it is also possible, as a first approximation, to consider only a single orbital to describe the metalligand interaction. This is the case for ligands of the type $\mathrm{AH}_{3}$ (or more generally $\mathrm{AR}_{3}$ ) whose orbital structure is summarized in Figure 1.4. Therefore, for an amine or phosphine (L-type ligands), it is in general sufficient to consider the nonbonding orbital $2 \mathrm{a}_{1}$ (Figure 1.4) that characterizes the lone pair on the nitrogen or phosphorus atom (1-33a). Analogous remarks may be made for the methyl ligand, $\mathrm{CH}_{3}$, or more generally for an alkyl radical $\mathrm{CR}_{3}$, the nonbonding orbital being only singly occupied in this case (an X-type ligand) (1-33b). This is the highest occupied orbital on the ligand, and its energy is not very different from that of the $d$ orbitals for most of the transition metals. Moreover, its
overlap with a metal orbital (e.g. $z^{2}, \mathbf{1 - 3 4}$ ) is substantial since it is a hybrid orbital polarized towards the metal centre. The resulting interaction produces a bonding and an antibonding MO. If the former is doubly occupied and the latter is empty, there is a $\sigma$ bond between the metal and the ligand ( $\sigma_{\mathrm{M}-\mathrm{PR} 3}$, or $\sigma_{\mathrm{M}-\mathrm{CR} 3}, \mathbf{1 - 3 4}$ ).


1-33a


1-33b


1-34

If we consider only the nonbonding orbital on ligands such as $\mathrm{NR}_{3}$, $\mathrm{PR}_{3}$, or $\mathrm{CR}_{3}$, we are effectively supposing that the interactions of the other MO with the metal orbitals are much weaker than the $\sigma$ interaction already described. This hypothesis can be justified by analysing the orbital structure of pyramidal $\mathrm{AH}_{3}$ molecules (Figure 1.4), where we see three bonding molecular orbitals $\left(\sigma_{\mathrm{A}}-\mathrm{H}\right)$ that account for the A-H bonds, the nonbonding orbital that is concentrated on the central atom, and three antibonding orbitals $\left(\sigma_{\mathrm{A}-\mathrm{H}}^{*}\right)$. The bonding orbitals of the ligand can interact with the metal orbitals. But they are far too low in energy, since they are orbitals describing the $\sigma_{\mathrm{A}-\mathrm{H}}$ bonds. Moreover, they are partially distributed over the hydrogen atoms, that is, in the direction away from the metal.

For both these reasons (substantial energy gap and poor overlap), the interactions involving the bonding orbitals are weaker than those concerning the nonbonding orbital. In a similar way, interactions between the antibonding $\sigma_{\mathrm{A}-\mathrm{H}}^{*}$ orbitals and the metal centre are usually negligible for the description of the metal-ligand bond; these MO are at high energy and partially oriented away from the metal.


Figure 1.4. Molecular orbitals for $\mathrm{AH}_{3}$ pyramidal molecules (with the electronic occupation appropriate for molecules with eight valence electrons, such as $\mathrm{NH}_{3}$ or $\mathrm{PH}_{3}$ ).

Figure 1.5. Molecular orbitals for $\mathrm{AH}_{2}$ molecules (with the electronic occupation appropriate for molecules with six valence electrons, such as $\mathrm{CH}_{2}$ or $\mathrm{SiH}_{2}$ in their lowest singlet state).

## Comment

It must, however, be noted that the antibonding orbitals can be strongly stabilized by very electronegative substituents, as, for example, in trifluorophosphine, $\mathrm{PF}_{3}$. In such a case, these orbitals can interact to a non-negligible extent with the $d$ orbitals of the metal centre. This is particularly important for the 2e orbitals (Figure 1.4) which can be involved in $\pi$-type interactions (see Chapter 3).

### 1.5.2. Several orbitals: $\sigma$ and $\pi$ interactions

For other ligands the situation is more complicated, since it is necessary to take several orbitals into account to obtain a satisfactory description of the bond with the metal. As a first example, we shall consider bent $\mathrm{AH}_{2}$ molecules.

### 1.5.2.1. Ligands of the $\mathrm{AH}_{2}$ type

The orbital structure of bent $\mathrm{AH}_{2}$ molecules (or more generally $\mathrm{AR}_{2}$ ), see Figure 1.5, shows that there are two bonding orbitals, which describe the $\sigma_{\mathrm{A}-\mathrm{H}}$ bonds, the two corresponding antibonding orbitals $\sigma_{\mathrm{A}-\mathrm{H}}^{*}$, and at an intermediate energy level, two nonbonding molecular orbitals: the $2 \mathrm{a}_{1}$ orbital, which is a hybrid pointing in the direction away from the hydrogen atoms, and the $1 \mathrm{~b}_{1}$ orbital, which is a pure $p$ atomic orbital perpendicular to the molecular plane, (see Chapter 6, $\S$ 6.5.2). Depending on the nature of the atom A , these two orbitals can contain one electron $\left(\mathrm{BH}_{2}, \mathrm{AlH} 2\right)$, two $\left(\mathrm{CH}_{2}, \mathrm{SiH}_{2}\right)$, three $\left(\mathrm{NH}_{2}, \mathrm{PH}_{2}\right)$, or four $\left(\mathrm{OH}_{2}, \mathrm{SH}_{2}\right)$.

As in the previous example, the bonding and antibonding orbitals can, in a first approximation, be neglected for the description of the metal-ligand interactions. However, it is necessary to take both nonbonding orbitals into account, as they are close in energy and both can lead to interactions with the metal centre. The $2 \mathrm{a}_{1}$ orbital plays the


$\sigma_{\mathrm{M}-\mathrm{L}}$
1-35


1-36
same role as the $2 \mathrm{a}_{1}$ orbital in $\mathrm{AH}_{3}$ molecules, and its interaction with a metal orbital ( $z^{2}$, for example) leads to the formation of an MO that characterizes a $\sigma$ bond (1-35).

The $1 \mathrm{~b}_{1}$ orbital $\left(p_{x}\right)$ has the correct symmetry to interact with the $x z$ orbital: the overlaps above and below the $y z$ plane have the same sign, so that the total overlap (the sum of the partial overlaps) is non-zero (1-36). This new metal-ligand interaction is said to be a $\pi$ interaction, since the orbitals concerned share a common nodal plane $(y z)$. It leads to the formation of two MO , one bonding, $\pi_{\mathrm{M}-\mathrm{L}}$, represented schematically in 1-36, and one antibonding, $\pi_{\mathrm{M}-\mathrm{L}}^{*}$. This interaction is particularly important for the description of the metal-ligand bond when two electrons are concerned: the bonding molecular orbital (1-36) is then doubly occupied but the antibonding orbital is empty. A $\pi$-type interaction therefore adds to and reinforces the $\sigma$ interaction (1-35), giving a double-bond character to the metal-ligand bond. It should be noted that there is an important difference between the $\sigma(\mathbf{1 - 3 5})$ and $\pi(\mathbf{1 - 3 6})$ interactions concerning the change in the overlap when there is a rotation about the M-L bond: the $\sigma$ overlap does not change, as it possesses cylindrical symmetry with respect to the bond direction, but the $\pi$ overlap is eliminated by a rotation of $90^{\circ}$.

## Comment

This is exactly analogous to what happens in an organic molecule such as ethylene (Figure 1.3), in which there is a $\sigma_{\mathrm{CC}}$ bond for any relative orientation of the two methylene groups, whereas the existence of a $\pi_{\mathrm{CC}}$ bond depends on the two groups being coplanar.

### 1.5.2.2. AH ligands

Following the reasoning developed in the previous paragraph, three molecular orbitals need to be considered for an A-H (or A-R) ligand (see Figure 1.6): the nonbonding orbital $2 \sigma$, analogous to the $2 \mathrm{a}_{1}$ orbital in $\mathrm{AH}_{3}$ and $\mathrm{AH}_{2}$ ligands, which allows a metal-ligand $\sigma$ bond to be formed, and the two degenerate $\pi$ orbitals ( $p_{x}$ and $p_{y}$ ) which may be involved in $\pi$ interactions with the metal centre. Depending on the nature of $A$, these three orbitals may contain two electrons $(\mathrm{BH}, \mathrm{AlH})$, three $(\mathrm{CH}, \mathrm{SiH})$, four $(\mathrm{NH}, \mathrm{SH})$, five $(\mathrm{OH}, \mathrm{SH})$, or six $(\mathrm{FH}, \mathrm{ClH})$.

### 1.5.2.3. Monoatomic ligands $A$

With the exception of the ligand H , for which there is only a single valence orbital $1 s_{\mathrm{H}}$ to consider (a $\sigma$ interaction), one must treat all the valence $s$ and $p$ orbitals on a monoatomic ligand A. Now the $s$ orbital is usually much lower in energy than the $d$ orbitals on the metal, especially

Figure 1.6. Molecular orbitals for AH molecules (with the electronic occupation appropriate for molecules with four valence electrons, such as BH or AIH in their lowest singlet state).

if A is a fairly electronegative element. In this case, one can therefore neglect the interaction of the $s$ orbital (which, if it is doubly occupied, therefore describes a lone pair localized on A), and only consider the three $p$ orbitals on the ligand, which are higher in energy and therefore closer to the metal $d$ orbitals. The one which points towards the metal centre ( $p_{z}, \mathbf{1 - 3 7}$ ) is used for the $\sigma$ interaction, and the two orbitals whose axis of revolution is perpendicular to the bond ( $p_{x}$ and $p_{y}$ ) can lead to $\pi$ interactions (1-38).


## Comment

In a more sophisticated model of the $\sigma$ orbitals, one can suppose that $p_{z}$ and $s$ are mixed, to form two $s-p$ hybrid orbitals, one pointing towards the metal to form the $\sigma$ bond, and the other in the opposite direction, so as to describe a lone pair of $\sigma$ type on A. This new approach does not fundamentally change anything in the simplified description given above.

The preceding examples show that one must always consider the atomic or molecular orbital on the ligand that allows a $\sigma$ bond to be formed with the metal. This orbital may be a nonbonding $s$ orbital (H), a hybrid $s-p$ orbital $\left(\mathrm{AH}_{3}, \mathrm{AH}_{2}\right.$, and AH molecules), or a $p$ orbital which points towards the metal centre (A atoms). When the atom bound to the metal also possesses nonbonding $p$ orbitals perpendicular to the metal-ligand bond $\left(\mathrm{AH}_{2}, \mathrm{AH}, \mathrm{A}\right)$, it is also necessary to consider them, since they lead to $\pi$-type interactions with the metal orbitals.

Even though it is an approximation to neglect the other MO on the ligand, this usually leads to an acceptable description of the bond with
the metal centre. It is particularly appropriate when the neglected ligand orbitals are very low in energy and the antibonding orbitals very high. These conditions are usually satisfied when the orbitals are involved in the $\sigma$ bonds of the ligand. However, when the ligand possesses one or more $\pi$ bonds, its $\pi$ bonding and $\pi^{*}$ antibonding molecular orbitals must usually be taken into account.

### 1.5.2.4. Ligands with a $\pi$ system: the example of CO

When the ligand possesses a $\pi$ bond involving the atom bound to the metal ( $\eta^{1}$ coordination), this leads to the presence of a $\pi$ bonding and $\pi^{*}$ antibonding orbital on the ligand. In general, the $\pi$ orbital is higher in energy than the MO that describe the $\sigma$ bonds, while the $\pi^{*}$ orbital is lower than the $\sigma^{*} \mathrm{MO}$ (§1.3.4). Although none of these $\pi$ or $\pi^{*}$ orbitals is nonbonding, in contrast to the $p$ orbitals of $\mathrm{AH}_{2}, \mathrm{AH}$, and A ligands, their energy level is neither sufficiently low ( $\pi$ ) nor sufficiently high $\left(\pi^{*}\right)$ for one to be able, a priori, to neglect their role in the metal-ligand interaction.

Carbon monoxide, CO, also called the carbonyl ligand, is an example which illustrates these points nicely. The essential features of its electronic structure are shown in Figure 1.7. The highest occupied orbital is a nonbonding $\sigma$ orbital, mainly concentrated on the carbon atom and polarized in the direction away from the oxygen atom. This orbital, which describes the lone pair on the carbon atom, is the one which allows a $\sigma_{\mathrm{M}}$ Co bond to be formed (L-type ligand). The two $\pi_{\mathrm{CO}}$ bonding orbitals associated with the $\pi$ bonds in $\mathrm{C} \equiv \mathrm{O}$ are lower in energy. They are mainly concentrated on oxygen, as that atom is more electronegative than carbon. The lowest empty orbitals are the antibonding $\pi_{\mathrm{CO}}^{*}$ orbitals, which have a larger contribution from carbon than oxygen. These four orbitals can lead to $\pi$-type interactions with orbitals of suitable symmetry on the metal, similar to those we have already seen with the nonbonding $p$ orbitals of the $\mathrm{AH}_{2}$ and AH molecules.

We shall therefore have to study a set of five orbitals (one $\sigma$, two $\pi$, and two $\pi^{*}$ ) when we wish to analyse the metal-carbonyl bond. ${ }^{12}$

### 1.5.2.5. $\pi$ complexes

In the examples studied so far, the ligand is bound to the metal centre by only one of its atoms. The situation is different when several atoms of the ligand are bound in an equivalent manner to the metal centre $\left(\eta^{x}\right.$ coordination). This is the case for $\pi$ complexes, in which the $\pi$ system of the ligand is oriented towards the metal. All the $\pi$ orbitals of the ligand, both occupied and empty, must now be considered to describe the metal-ligand bonds. As an example, we shall treat an $\eta^{2}$-ethylene complex in detail in Chapter 3 (§ 3.4).

Figure 1.8. Simplified diagram for the interaction of the atomic orbitals on a metal centre and the $\ell$ ligands which surround it ( $\sigma$ interactions only).

### 1.6. Initial orbital approach to $M_{\ell}$ complexes

The shape and the energy of the molecular orbitals of a complex depend on the number of ligands and their geometrical arrangement around the metal. It is possible to obtain some important information on the MO without defining the particular complex studied. The purpose of this paragraph is thus to derive the general characteristics of the orbital structure which do not depend (or depend only slightly) on the complex.

### 1.6.1. Simplified interaction diagram

We shall consider, for simplicity, a complex in which the metal is surrounded by $\ell$ identical ligands, each with just a single orbital that can take part in the metal-ligand interaction (a $\sigma$ interaction, $\S$ 1.5.1).

A simplified diagram for the interaction between the $\ell$ ligand orbitals and the nine atomic orbitals on the metal (five $d$ orbitals, one $s$ orbital, and three $p$ orbitals, without distinction) is given in Figure 1.8. In this diagram, the metal orbitals are placed higher in energy than those on the ligands, since the latter are more electronegative. The $\ell$ ligand orbitals interact with $\ell$ metal orbitals, to form $\ell$ bonding MO and the associated $\ell$ antibonding MO. There are therefore $(9-\ell)$ nonbonding orbitals that remain on the metal.



Given the relative energies for the participating orbitals (§ 1.3.2), we can make the following points:

1. The bonding MO , which describe the $\sigma_{\mathrm{M}}$-Lig bonds, are mainly concentrated on the ligand orbitals. An example is given for an $\mathrm{M}-\mathrm{H}$ bond involving the metal $z^{2}$ orbital (1-39).
2. The corresponding antibonding MO are mainly concentrated on the metal orbitals (1-40).
3. The nonbonding MO are orbitals that are localized on the metal centre. A more detailed analysis of the orbital structure of complexes (Chapter 2) will show that usually, but not always, these nonbonding MO are pure $d$ orbitals, or orbitals in which the principal component is of $d$ type.

### 1.6.2. Strong-field and weak-field complexes

The separation between the energy levels ( $\Delta E^{-}$, Figure 1.8) of the nonbonding and antibonding $\mathrm{MO}\left(\sigma_{\mathrm{M}-\mathrm{Lig}}^{*}\right)$ is directly linked to the strength of the interaction between the ligand orbitals and those on the metal. The stronger this interaction, the more the antibonding orbitals are destabilized, so the larger the energy gap $\Delta E^{-}$. When the metal-ligand interaction is strong, $\Delta E^{-}$is large and one refers to strongfield complexes; in contrast, when $\Delta E^{-}$is small, one refers to weak-field complexes.

### 1.6.3. Electronic configuration and the 18 -electron rule

In so far as the electronic occupation of the molecular orbitals is concerned, the stability of an $\mathrm{ML}_{\ell}$ complex is generally maximized when the bonding MO , of which there are $\ell$, together with the nonbonding MO , of which there are $(9-\ell)$, are doubly occupied, but when the $\ell$ antibonding MO remain empty (Figure 1.8). The bonding MO describe the M-Lig bonds and the nonbonding MO represent lone pairs on the metal. In this situation, the total number of electrons is:

$$
\begin{equation*}
N_{\mathrm{t}}=(2 \times \ell)+2 \times(9-\ell)=18 \tag{1.13}
\end{equation*}
$$

In this way, we rationalize the 18 -electron rule that was previously discussed with reference to the valence electronic structure of the nearest noble gas (§ 1.1.2.1).

The electrons that occupy the nonbonding MO are not used to form metal-ligand bonds. They therefore correspond to the $n$ electrons that 'remain' on the metal in the classical counting scheme (§ 1.1.2.3). The $d^{n}$ notation for the electronic configuration of the metal assumes, however,
that all the occupied nonbonding orbitals on the metal are of $d$ type (see Chapter 2).

Even if the 18 -electron rule is often obeyed, we must not forget that there are many exceptions. There are some complexes that have fewer than 18 electrons. For example, the $\left[\mathrm{M}(\mathrm{Lig})_{4}\right]$ complexes that adopt a 'square-planar' geometry (four ligands at the vertices of a square whose centre is occupied by the metal) with a $d^{8}$ electronic configuration (e.g.: $\left[\operatorname{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right], \S$ 1.1.2.1) are 16 -electron complexes. Since the bonding MO that describe the bonds are doubly occupied, this analysis shows that one of the nonbonding orbitals in Figure 1.8 is empty, and a more detailed study of the electronic structure is necessary to understand this result (Chapter 2, §2.2). These 16 -electron complexes are stable, but often reactive towards other molecules since they have a tendency to form 18 -electron complexes, by binding other ligands. For example, Wilkinson's catalyst $\left[\mathrm{Rh}\left(\mathrm{PPh}_{3}\right)_{3} \mathrm{Cl}\right]$ is used industrially for the catalytic hydrogenation of olefins (see Exercise 1.3).

There are also some complexes with more than 18 electrons, such as $\left[\mathrm{Ni}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ which possesses 20 electrons (§ 1.1.2.1). Some antibonding MO must therefore be occupied, which can happen only if they are sufficiently low in energy, as is the case in weak-field complexes. Organometallic complexes, characterized by the presence of one or several metal-carbon bonds, are strong-field complexes. It is therefore rare for them to possess more than 18 electrons.

### 1.6.4. Analogy with the octet rule

In the same way, one can construct a simplified interaction diagram for $\mathrm{AH}_{n}\left(\right.$ or $\left.\mathrm{AR}_{n}\right)$ molecules in which A is an element from the second or third row of the periodic table ( $\mathrm{C}, \mathrm{Si}, \mathrm{N}, \mathrm{P}, \mathrm{O}, \mathrm{S}$, etc.). There are thus four valence orbitals on the central atom: one $s$ AO, and three $p$ AO. The $\sigma$ interactions with the orbitals on the atoms bound to A lead to the formation of $n$ bonding MO, $n$ antibonding MO and $(4-n)$ nonbonding MO. If the bonding and nonbonding MO are doubly occupied, the number of electrons $N_{\mathrm{t}}$ is equal to

$$
\begin{equation*}
N_{\mathrm{t}}=(2 \times n)+2 \times(4-n)=8 \tag{1.14}
\end{equation*}
$$

We have therefore derived the octet rule. As examples, we can quote:
$\mathrm{CH}_{4} \quad n=4 \quad 4$ bonding MO, 4 antibonding MO, 0 nonbonding MO
$\mathrm{NH}_{3} n=33$ bonding MO, 3 antibonding MO, 1 nonbonding MO
$\mathrm{OH}_{2} n=22$ bonding MO, 2 antibonding MO, 2 nonbonding MO
FH $n=1 \quad 1$ bonding MO, 1 antibonding MO, 3 nonbonding MO

It is straightforward to verify that the number of bonding, nonbonding, and antibonding MO predicted by this simple model agree consistently with the detailed orbital structures of the $\mathrm{NH}_{3}, \mathrm{OH}_{2}$, and HF molecules (§ 1.1.4-1.1.6).

## Exercises

## 1.1

What are the two coordination modes of the allyl ligand, $\mathrm{H}_{2} \mathrm{C}-\mathrm{CH}-\mathrm{CH}_{2}$, for which a Lewis structure is given below? In each case, categorize the ligand as $\mathrm{L}_{\ell} \mathrm{X}_{x}$.

1.2

Write each of the following complexes as $\left[\mathrm{ML}_{\ell} \mathrm{X}_{x}\right]^{q}$, give the oxidation state of the metal no, the electronic configuration $d^{n}$, and the total number of electrons $N_{t}$. (1) $\left[\operatorname{Cr}(\mathrm{CO})_{6}\right]$; (2) $\left[\mathrm{W}(\mathrm{CO})_{5}\right] ;$ (3) $\left[\mathrm{Mn}(\mathrm{CO})_{5} \mathrm{Cl}\right] ;$ (4) $\left[\mathrm{TiCl}_{4}\right] ;$ (5) $\left[\mathrm{Co}(\mathrm{CO})_{3}(\mathrm{Et})\right]$; (6) $\left[\mathrm{Re}\left(\mathrm{PR}_{3}\right)(\mathrm{CO})_{4} \mathrm{Cl}\right]$; (7) $\left[\mathrm{Fe}(\mathrm{CO})_{4}(\mathrm{H})_{2}\right]$; (8) $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$; (9) $\left[\mathrm{ReH}_{9}\right]^{2-} ; \quad$ (10) $\quad\left[\operatorname{ReH}_{5}\left(\mathrm{PR}_{3}\right)_{2}\left(\mathrm{SiR}_{3}\right)_{2}\right] ;$ (11) $\quad\left[\mathrm{Ni}(\mathrm{CO})_{4}\right]$; (12) $\left[\mathrm{Cu}(\mathrm{SR})_{3}\right]^{2-}$; (13) $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$; (14) $\left[\mathrm{RhI}_{3}(\mathrm{CO})_{2}(\mathrm{Me})\right]^{-}$; (15) $\left[\operatorname{RhI}_{3}(\mathrm{CO})(\mathrm{COMe})\right]^{-} ;(16)\left[\mathrm{MoF}_{4}(\mathrm{O})_{2}\right]^{2-} ;(17)\left[\operatorname{Re}\left(\mathrm{NR}_{3}\right)_{4}(\mathrm{O})_{2}\right]^{+}$; (18) $\left[\mathrm{Mn}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)(\mathrm{CO})_{3}\right]^{+} ; \quad$ (19) $\quad\left[\mathrm{Zr}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}(\mathrm{H})(\mathrm{Cl})\right]$; (20) $\left[\mathrm{Nb}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}(\mathrm{Me})_{3}\right] ;(21)\left[\mathrm{Os}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)(\mathrm{CO})_{2} \mathrm{Cl}\right] ;(22)\left[\mathrm{WCl}_{6}\right]$;
(23) $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right] ;(24)\left[\mathrm{Zn}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)(\mathrm{Me})\right]$.

## 1.3

Wilkinson's catalyst $\left[\mathrm{RhL}_{3} \mathrm{Cl}\right]$, where $\mathrm{L}=\mathrm{PPh}_{3}$, is used for the hydrogenation of alkenes. The reaction proceeds as follows:

$$
\begin{gathered}
{\left[\mathrm{RhL}_{3} \mathrm{Cl}\right] \xrightarrow{\mathrm{H}_{2}}\left[\mathrm{RhL}_{3} \mathrm{Cl}(\mathrm{H})_{2}\right] \longrightarrow\left[\mathrm{RhL}_{2} \mathrm{Cl}\left(\mathrm{H}_{2}\right)\right]+\mathrm{L} \stackrel{\mathrm{C}_{2} \mathrm{H}_{4}}{\longrightarrow}\left[\mathrm{RhL}_{2} \mathrm{Cl}(\mathrm{H})_{2}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]+\mathrm{L}} \\
\downarrow \\
{\left[\mathrm{RhL}_{3} \mathrm{Cl}\right]+\mathrm{CH}_{3}-\mathrm{CH}_{3} \longleftarrow\left[\mathrm{RhL}_{3} \mathrm{Cl}(\mathrm{H})\left(\mathrm{CH}_{2} \mathrm{CH}_{3}\right)\right] \longleftarrow\left[\mathrm{RhL}_{2} \mathrm{Cl}(\mathrm{H})\left(\mathrm{CH}_{2} \mathrm{CH}_{3}\right)\right]+\mathrm{L}}
\end{gathered}
$$

For each step, give the oxidation state of the metal (no), the electronic configuration of the complex $\left(d^{n}\right)$, and the total number of electrons $\left(N_{\mathrm{t}}\right)$.

## 1.4

1. Give the charge on the ligands $\mathrm{CO}, \mathrm{Cl}, \mathrm{Et}, \mathrm{PR}_{3}, \mathrm{H}, \mathrm{H}_{2}, \mathrm{SiR}_{3}, \mathrm{SR}$, $\mathrm{CN}, \mathrm{I}, \mathrm{Me}, \mathrm{COMe}, \mathrm{F}, \mathrm{O}, \mathrm{NR}_{3}, \mathrm{C}_{2} \mathrm{H}_{4}, \mathrm{C}_{6} \mathrm{H}_{6}, \mathrm{C}_{5} \mathrm{H}_{5}$ if the ionic model is adopted (§ 1.2).
2. Formulate the complexes given in $\S$ 1.1.2.2 if this model is adopted, and hence deduce the oxidation state of the metal.

## 1.5

1. Give the oxidation state of the metal no, the electronic configuration $d^{n}$, and the total number of electrons $N_{t}$ for the complexes $\left[\mathrm{Cr}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)_{2}\right]$ and $\left[\mathrm{Ru}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\left(\eta^{4}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\right]$.
2. How can one explain the change in the coordination mode $\eta^{6} \rightarrow$ $\eta^{4}$ of one of the ligands when chromium in the complex is replaced by ruthenium?
3. Using the same argument, predict the hapticity $x$ for a cyclopentadienyl ligand in the following complexes: (i) $\left[\mathrm{Mn}(\mathrm{CO})_{3}\left(\eta^{x}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\right]$; (ii) $\left[\mathrm{W}(\mathrm{CO})_{2}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\left(\eta^{x}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\right]$; (iii) $\left[\mathrm{Fe}(\mathrm{CO})_{2}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\left(\eta^{x}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\right]$.

## 1.6

The borohydride ligand $\left(\mathrm{BH}_{4}^{-}\right.$in the ionic model) can bind in the $\eta^{1}, \eta^{2}$, or $\eta^{3}$ modes, depending on whether one, two, or three B-H bonds interact with the metal centre.


$\eta^{2}$


Rationalize the coordination mode in the following complexes: $\left[\mathrm{Cu}\left(\mathrm{PR}_{3}\right)_{3}\left(\eta^{1}-\mathrm{BH}_{4}\right)\right],\left[\mathrm{Cu}\left(\mathrm{PR}_{3}\right)_{2}\left(\eta^{2}-\mathrm{BH}_{4}\right)\right]$, and $\left[\mathrm{Ti}(\mathrm{CO})_{4}\left(\eta^{3}-\mathrm{BH}_{4}\right)\right]^{-}$.

## 1.7

What is the oxidation state of the metal centres in the following binuclear complexes: (1) $\left[\operatorname{Re}(\mathrm{CO})_{5}\right]_{2}$; (2) $\left[\mathrm{ReCl}_{4}\left(\mathrm{H}_{2} \mathrm{O}\right)\right]_{2}^{2-}$; (3) $\left[\mathrm{MoCl}_{2}\left(\mathrm{PR}_{3}\right)_{2}\right]_{2}$; (4) $\left[\mathrm{Pd}\left(\eta^{3}-\mathrm{C}_{3} \mathrm{H}_{5}\right)(\mu-\mathrm{Cl})\right]_{2}$; (5) $\left[\mathrm{Mo}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)(\mathrm{CO})_{2}(\mu-\mathrm{SR})\right]_{2}$.

## 1.8

Consider the interaction between a metal atom and a ligand with an $s$ orbital.

1. Show that there can always be an interaction between the ligand orbital and the $s$ orbital of the metal. Sketch the shapes of the resulting bonding and antibonding MO.
2. What position must the ligand adopt to give (i) maximal and (ii) minimal overlap with the $p_{x}$ orbital of the metal? What is its value in the latter case?
3. Repeat question 2 for the $x y$ and $z^{2}$ metal orbitals.

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## Principal ligand fields: $\sigma$ interactions

${ }^{1}$ Those ligands for which it is necessary to consider $\pi$-type interactions with the metal are studied in Chapter 3.

In this chapter, we shall construct the molecular orbitals (MO) of monometallic complexes $\mathrm{ML} \ell$ and thereby deduce their electronic structure by distributing the electrons in these orbitals. We shall study different types of 'ligand field', each one being characterized by the number of ligands and by their geometrical arrangement around the metal centre (octahedral complexes $\mathrm{ML}_{6}$, tetrahedral, or square-planar $\mathrm{ML}_{4}$, etc.).

We shall always begin the analysis by establishing the molecular orbitals of the associated model complex, in which all the ligands (i) are identical and (ii) only have $\sigma$-type interactions with the metal (Chapter 1, $\S$ 1.5.1). The resulting orbital scheme is characteristic of the ligand field being studied (octahedral, tetrahedral, square-planar, etc.). In a first approximation, it is applicable to all complexes of this type, even if the two conditions specified above are not met exactly. The main reason for this is that $\sigma$ interactions exist in all complexes, and they are stronger than $\pi$ interactions when the latter are present. Therefore, if $\pi$-type interactions are added in a 'real' complex to $\sigma$ interactions, while the results are different from those obtained for the model complex, they are not completely transformed. The $\pi$ effects will therefore be treated subsequently as perturbations to be added to the orbital scheme established for the model complex. ${ }^{1}$ In the same way, different $\sigma$ interactions associated with non-equivalent ligands produce changes compared to the model complex that can initially be neglected.

There are several ways of representing the ligand orbital that is involved in the $\sigma$ interaction with the metal: as an $s$ orbital ( $\mathbf{2 - 1 \mathbf { 1 a } \text { ), as a } p}$ orbital (2-1b), or as a hybrid $(s-p)$ orbital directed towards the metallic centre (2-1c). We shall use this last representation, since, except for hydride ligands which possess only a single valence orbital $1 s$ and very electronegative monoatomic ligands such as F in which the contribution of the $p$ orbital dominates, it is the most appropriate for all other ligands. The orbital on ligand $\mathrm{L}_{i}$ will be written $\sigma_{i}$, due to the nature of the bond which it can form with the metal.


To construct the molecular orbitals of a complex, we have to determine their shape, by which we mean the contributions of the metal and ligand orbitals to each one, as well as their relative energies. The 'fragment method' is widely used for the construction of the MO of a complex $\mathrm{ML}_{\ell}$; we allow the atomic orbitals ( AO ) of the central metal to interact with those on the ligands. Knowledge of the symmetry properties of these orbitals allows us to simplify the construction of these diagrams very considerably. If the two orbitals have the same symmetry properties, their overlap is non-zero and an interaction can occur. On the other hand, two orbitals of different symmetry have zero overlap ('by symmetry') and do not interact at all (Chapter 1, § 1.3.3). In this context, we shall sometimes use basic ideas from group theory, as well as several results established in Chapter 6 which the reader may consult when necessary.

### 2.1. Octahedral $\mathrm{ML}_{6}$ complexes

Consider a complex in which the metal centre is surrounded by six identical ligands ( $\mathrm{L}_{1}-\mathrm{L}_{6}$ ) placed at the vertices of an octahedron (2-2). The metal is placed at the origin, and the ligands on the axes $x\left(\mathrm{~L}_{2}, \mathrm{~L}_{4}\right)$, $y\left(\mathrm{~L}_{1}, \mathrm{~L}_{3}\right)$, and $z\left(\mathrm{~L}_{5}, \mathrm{~L}_{6}\right)$. Each ligand $\mathrm{L}_{i}$ possesses an orbital $\sigma_{i}$ directed towards the metal (2-3).


2-2


2-3


2-4 $(S \neq 0)$

### 2.1.1. Initial analysis of the metal-ligand orbital interactions

Our first task is to examine how each of the orbitals on the ligands can interact with the $s, p$, and $d$ orbitals on the metal centre.

As the $s$ orbital has spherical symmetry, its overlap with any one of the $\sigma_{i}$ orbitals is non-zero ( $\sigma_{1}$ in 2-4). And as the six ligands are equivalent, this overlap is the same no matter which orbital $\sigma_{i}$ is considered, so the interactions between $s$ orbital and each of the ligand orbitals are identical.

We therefore find a bonding MO between the metal and the ligands (2-5) and a corresponding antibonding MO (2-6), where each


2-5
(bonding)


2-7 $(S \neq 0)$


2-6
(antibonding)

$2-8(S=0)$

(bonding)


2-11
(bonding)


2-10
(antibonding)


2-12 (antibonding)


2-13
(bonding)
2-14
(antibonding)
$\sigma_{i}$ contribution is equal in both orbitals. Since the ligands are more electronegative than the metal, the bonding orbital is mainly concentrated on the ligands but the antibonding orbital on the metal (Chapter 1, § 1.6.1).

We now turn to the $p$-type orbitals on the central atom. The $p_{y}$ orbital overlaps with the orbitals on the ligands $\mathrm{L}_{1}$ and $\mathrm{L}_{3}$ located on the $y$-axis ( $\sigma_{3}$ in 2-7). But the four other ligands $\mathrm{L}_{2}, \mathrm{~L}_{4}-\mathrm{L}_{6}$, are placed in the nodal plane of $p_{y}(x z)$, which is a symmetry element of the complex. The $\sigma_{i}$ orbitals associated with these ligands are symmetric with respect to this plane, whereas the $p_{y}$ orbital is antisymmetric: the overlaps between $p_{y}$ and each of $\sigma_{2}, \sigma_{4}-\sigma_{6}$ are therefore zero 'by symmetry'. In the example shown in 2-8, it is clear that the $\sigma_{6}$ orbital has a positive overlap with the grey lobe of the $p_{y}$ orbital but an equivalent negative overlap with the white lobe, so the total overlap is equal to zero. Therefore, there cannot be any interaction between the $p_{y}$ orbital and any of these four ligand orbitals.

The $p_{y}$ orbital on the metal therefore combines with the $\sigma_{1}$ and $\sigma_{3}$ orbitals to form a bonding MO, mainly based on the ligands (2-9), and an antibonding MO mainly based on the metal (2-10).

This analysis is readily extended to the $p_{x}$ and $p_{z}$ orbitals, which can combine only with $\sigma_{2}$ and $\sigma_{4}$ for $p_{x}$ (2-11 and 2-12) or $\sigma_{5}$ and $\sigma_{6}$ for $p_{z}$ (2-13 and 2-14). The interactions involving the $p$ orbitals of the metal therefore lead to the formation of three bonding and three antibonding MO. The three bonding orbitals differ from each other only by their orientations: they have the same energy, and are therefore degenerate. The same is clearly true for the antibonding orbitals.

We now turn to the interactions that involve the $d$ orbitals of the metal centre. We consider first the three orbitals $x y, x z$, and $y z$ that each contain two nodal planes, $(x z, y z),(x y, y z)$, and $(x y, x z)$, respectively. The six ligands, placed on the $x$-, $y$-, or $z$-axes (2-2), are all located in one of the two nodal planes of the $d$ orbitals, and sometimes even at the intersection of these two planes. The overlap between any one of the $\sigma_{i}$ orbitals and the $y z(\mathbf{2}-15), x y(\mathbf{2 - 1 6})$ or $x z$ orbitals (2-17) is therefore equal to zero.


Therefore, there cannot be any interaction between the three $x y$, $x z$, and $y z$ orbitals and the $\sigma_{i}$ orbitals on the ligands. In an octahedral complex, these three $d$ orbitals therefore form a degenerate nonbonding set, located only on the metal (2-18 to 2-20).


The situation is different for the two remaining $d$ orbitals. The $x^{2}-y^{2}$ orbital points towards the ligands located on the $x$ - and $y$-axes ( $\mathrm{L}_{1}-\mathrm{L}_{4}$ ). The overlaps with the orbitals on these four ligands ( $\sigma_{1}$ in 2-21) are therefore non-zero and all equal (in absolute value), since the $x^{2}-y^{2}$ orbital has the same amplitude (in absolute value), at a given distance from the metal, in the direction of each of the ligands. But the ligands $\mathrm{L}_{5}$ and $\mathrm{L}_{6}$ are located on the $z$-axis, that is, the intersection of the two nodal planes of the $x^{2}-y^{2}$ orbital (the planes which bisect the $x$ - and $y$-axes). No overlap is therefore possible between $x^{2}-y^{2}$ and the orbitals on these last two ligands ( $\sigma_{6}$ in 2-22).

The interaction of the $x^{2}-y^{2}$ orbital with the ligand orbitals therefore leads to a bonding MO, mainly concentrated on the ligands (2-23), and an antibonding MO, essentially the metal orbital (2-24). In each of these MO , the coefficients on the ligands $\mathrm{L}_{1}-\mathrm{L}_{4}$ are equal in absolute value, since the overlaps concerned are identical, but the coefficients on $\mathrm{L}_{5}$ and $\mathrm{L}_{6}$ are zero.

The last orbital to be considered is the $z^{2}$ orbital. It is mainly oriented along the $z$ axis, but also, to a lesser degree, in the $x y$ plane; none of the ligands is located in its nodal cone (Chapter 1, Scheme 1-31). It can therefore interact with the orbitals of all six ligands, though the overlap with a ligand placed on the $z$-axis ( $\sigma_{6}$ in 2-25) is larger than that with a ligand in the $x y$ plane ( $\sigma_{1}$ in 2-26).

The bonding (2-27) and antibonding (2-28) MO formed from the $z^{2}$ orbital therefore contain contributions from the six ligands, with the coefficients for $\mathrm{L}_{5}$ and $\mathrm{L}_{6}$, placed on the $z$-axis, being larger than those for $\mathrm{L}_{1}-\mathrm{L}_{4}$, in the $x y$ plane (Chapter 6, § 6.6.5).

This initial analysis of the orbital structure of octahedral complexes has relied simply on the existence, or the absence, of an overlap between the orbitals of the metal and those on the ligands, and on the difference in electronegativity between the metal and the ligands. It enables us to obtain several important results: (i) there are six bonding MO, mainly concentrated on the ligands, and six antibonding MO, mostly on the metal; (ii) there are three remaining orbitals located on the metal which take no part in the metal-ligand bonds (nonbonding orbitals). This description is consistent with the simplified scheme given in Chapter 1 (§ 1.6.1, Figure 1.8) which predicted the formation of $\ell$ bonding MO, $\ell$ antibonding MO, and $(9-\ell)$ nonbonding MO in ML $\ell$ complexes. It also provides a qualitative description of the shape of
${ }^{2}$ In other words, the linear combinations of orbitals which are bases for an irreducible representation of the symmetry point group (see Chapter 6, § 6.4).

[^0]these MO, and shows that the three nonbonding orbitals are, in the particular case of an octahedral complex, pure $d$ orbitals on the metallic centre.

### 2.1.2. Complete interaction diagram

The analysis presented above allows us to establish the main features of the orbital structure of octahedral complexes. It also shows that certain molecular orbitals are degenerate. This is the case, for example, for the three bonding MO constructed from the metal $p$ orbitals (2-9, 2-11, and $\mathbf{2 - 1 3}$ ), or for the three nonbonding orbitals (2-18 to 2-20). However, other degeneracies exist which cannot readily be shown by the analysis above. To obtain a more complete description of the orbital structure, it is necessary to exploit all the symmetry properties of the octahedron and not simply this or that symmetry element as we have done so far.

There are four stages in the general procedure for constructing the MO of an ML $\ell$ complex:

1. Find the appropriate point-group symmetry. In this case, it is the octahedral group, $O_{h}$.
2. Determine the symmetry properties of the orbitals on the central metal atom. They are given directly by the character table for the point group.
3. Do not consider the ligand orbitals individually, but use linear combinations of these orbitals which are 'adapted' to the symmetry of the complex. ${ }^{2}$
4. Allow the metal and ligand orbitals to interact (fragment method). Only orbitals of the same symmetry can interact, since their overlap is non-zero.

### 2.1.2.1. The symmetry of metal orbitals

Examination of the character table of the octahedral point group $\left(O_{h}\right)$ gives us the symmetry of the metal orbitals directly, as this atom is located at the origin of the $x, y, z$ frame (Table 2.1). The last two columns indicate the symmetry of several functions of $x, y$, and $z$ which have a direct link with the analytical expressions, in Cartesian coordinates, of these orbitals (Chapter 1, § 1.4.1).

The $s$ orbital, which has spherical symmetry, has the same symmetry properties as the function $x^{2}+y^{2}+z^{2}$. It is therefore a basis of the irreducible representation $\mathrm{A}_{1 g}$ (the totally symmetric representation of the $O_{h}$ group), or more simply 'it has $\mathrm{A}_{1 g}$ symmetry'. The $p_{x}, p_{y}$, and $p_{z}$ orbitals transform as $x, y$, and $z$, respectively. Their symmetry is therefore $\mathrm{T}_{1 u}$ (a three-dimensional, or triply degenerate, representation). ${ }^{3}$ Finally, the

Table 2.1. Character table of the point group $O_{h}$

| $\mathrm{O}_{h}$ | E | $8 C_{3}$ | $6 C_{2}^{\prime}$ | $6 C_{4}$ | $3 C_{2}$ | $i$ | $8 S_{6}$ | $6 \sigma_{d}$ | $6 S_{4}$ | $3 \sigma_{h}$ |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | :--- |
| $\mathrm{~A}_{1 g}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | $x^{2}+y^{2}+z^{2}$ |
| $\mathrm{~A}_{2 g}$ | 1 | 1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 | -1 |  |
| $\mathrm{E}_{g}$ | 2 | -1 | 0 | 0 | 2 | 2 | -1 | 0 | 0 | 2 | $\left(z^{2}, x^{2}-y^{2}\right)$ |
| $\mathrm{T}_{1 g}$ | 3 | 0 | -1 | 1 | -1 | 3 | 0 | -1 | 1 | -1 |  |
| $\mathrm{~T}_{2 g}$ | 3 | 0 | 1 | -1 | -1 | 3 | 0 | 1 | -1 | -1 |  |
| $\mathrm{~A}_{1 u}$ | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 |  |
| $\mathrm{~A}_{2 u}$ | 1 | 1 | -1 | -1 | 1 | -1 | -1 | 1 | 1 | -1 |  |
| $\mathrm{E}_{u}$ | 2 | -1 | 0 | 0 | 2 | -2 | 1 | 0 | 0 | -2 |  |
| $\mathrm{~T}_{1 u}$ | 3 | 0 | -1 | 1 | -1 | -3 | 0 | 1 | -1 | 1 | $(x, y, x z, y z)$ |
| $\mathrm{T}_{2 u}$ | 3 | 0 | 1 | -1 | -1 | -3 | 0 | -1 | 1 | 1 |  |



Figure 2.1. Symmetry-adapted $\sigma$ orbitals for an octahedral complex $\mathrm{ML}_{6}$, and orbitals of the same symmetry on the metal centre (consult Chapter 6, $\S 6.6 .5$ for the analytical expressions of these symmetry-adapted orbitals and the method for finding them).
$d$ orbitals have $\mathrm{E}_{g}$ symmetry for the $z^{2}$ and $x^{2}-y^{2}$ pair, but $\mathrm{T}_{2 g}$ for the $x y, x z$, and $y z$ set. It is thus important to realize that the five $d$ orbitals, which clearly have the same energy for an isolated atom, are separated into two groups according to their symmetry properties: a doubly degenerate representation $\mathrm{E}_{g}$ and a triply degenerate set $\mathrm{T}_{2 g}$.

### 2.1.2.2. Symmetry-adapted orbitals on the ligands

None of the ligand orbitals $\sigma_{i}$, considered individually, is a basis for the irreducible representations of the $O_{h}$ group. For the construction of the orbital interaction diagram, it is therefore necessary to use linear combinations of these orbitals which are adapted to the symmetry of the octahedron. These are presented in Figure 2.1, next to the metal orbitals with the same symmetry.

In the orbital of $\mathrm{A}_{1 g}$ symmetry, all the coefficients have the same value and sign. It is clearly a totally symmetric orbital, just like the $s$ orbital on the central atom. Out-of-phase combinations on two ligands in trans positions appear in the three $\mathrm{T}_{1 u}$ orbitals, along the $x, y$, or $z$ axes. It is clear by inspection of the shape of these orbitals that they differ only in their orientations and that they have the same symmetry properties as the $p_{x}, p_{y}$, and $p_{z}$ orbitals of the metal centre (Figure 2.1). These combinations of ligand orbitals are therefore degenerate by symmetry, just like the metal $p$ orbitals ( $\mathrm{T}_{1 u}$ representation).

In the case of the first orbital of $\mathrm{E}_{g}$ symmetry, the coefficients for the four ligands in the $x y$ plane are all of equal value, but of opposite sign, depending whether they are on the $x$ - or the $y$-axis. The symmetry
properties of this orbital are therefore just the same as those of the metal $x^{2}-y^{2}$ orbital, which also has $\mathrm{E}_{g}$ symmetry. The second $\mathrm{E}_{g}$ orbital contains contributions from all six ligands. The coefficients for the ligands located on the $z$-axis are twice as large as those for the ligands in the $x y$ plane, and of opposite sign. The $z^{2}$ orbital, which also has $\mathrm{E}_{g}$ symmetry, shares these properties; its amplitude is proportional to $\left(2 z^{2}-x^{2}-y^{2}\right)$ (Chapter 1, formula (1.11)). It is important to realize that the two $\mathrm{E}_{g}$ orbitals that are degenerate by symmetry necessarily have the same energy (Chapter 6). In contrast to the case of the three orbitals of $\mathrm{T}_{1 u}$ symmetry, this result is not at all obvious just from the coefficients for the ligands in the two orbitals. This additional information, obtained after proper consideration of the octahedral symmetry of the complex, enables us to produce a complete description of their orbital structure.

## Comment

In what follows, the symmetry of atomic or molecular orbitals will be represented with lower-case letters ( $\mathrm{a}_{1 g}, \mathrm{t}_{1 u}$, $\mathrm{e}_{g}$, $\mathrm{t}_{2 g}$, etc.), while uppercase letters will be reserved for the different irreducible representations in character tables, for the symmetry of electronic states and for spectroscopic terms.

### 2.1.2.3. Interaction diagram

We can now construct all the MO of an octahedral complex, by allowing the metal orbitals (Figure 2.2, left-hand side) to interact with the symmetry-adapted combinations of orbitals on the six ligands (Figure 2.2, right-hand side).

The energetic ordering adopted for the metal orbitals follows the usual rule for the $d$-block transition metals: $\varepsilon(d)<\varepsilon(s)<\varepsilon(p)$ (Chapter 1, § 1.4.2). The six ligand-orbital combinations are not all at the same energy. To understand the energetic ordering in Figure 2.2, one must analyse the bonding or antibonding character of each for the ligand-ligand interactions (Figure 2.1). The $\mathrm{a}_{1 g}$ orbital is the most stable, since all the interactions are bonding. In the $\mathrm{t}_{1 u}$ orbitals, there is an antibonding interaction between two ligands that are trans to each other. Due to the large distance between these ligands, the orbitals may be considered to be practically nonbonding. For the two $\mathrm{e}_{g}$ orbitals (which necessarily have the same energy), it is easier to consider the first: the four interactions between the closer ligands (cis) are antibonding, but two weakly bonding interactions occur between the trans ligands. Overall, the $\mathrm{e}_{g}$ orbitals are therefore antibonding, leading to the energetic ordering: $\varepsilon\left(\mathrm{a}_{1 g}\right)$ (bonding) $<\varepsilon\left(\mathrm{t}_{1 u}\right)$ (nonbonding $)<\varepsilon\left(\mathrm{e}_{g}\right)$ (antibonding). One must remember that the bonding or
antibonding interactions just described never involve atoms that are directly bonded to each other. The overlaps involved are therefore weak, due to the distance between the centres concerned (e.g. if the metal-ligand bond distances are $2.2 \AA$, the corresponding cis-ligand separation is $3.11 \AA$, and as much as $4.40 \AA$ for the trans ligands). This explains why the energy separations between the different ligand orbitals are small (Figure 2.2, right-hand side). We conclude by noting that the ligand orbitals are placed lower in energy than those on the metal, since the ligands are in general more electronegative than the metal.

The molecular orbitals are obtained by allowing the orbitals of the M and $\mathrm{L}_{6}$ fragments to interact (Figure 2.2). We note features that were established in § 2.1.1: the formation of six bonding MO $\left(1 \mathrm{a}_{1 g}, 1 \mathrm{t}_{1 u}\right.$, and $1 \mathrm{e}_{\mathrm{g}}$ ), by bonding interactions involving the fragment orbitals whose symmetries are $\mathrm{a}_{1 g}, \mathrm{t}_{1 u}$, and $\mathrm{e}_{g}$, and the formation of six corresponding antibonding MO $\left(2 \mathrm{e}_{g}, 2 \mathrm{a}_{1 g}\right.$, and $\left.2 \mathrm{t}_{11}\right)$. The new feature concerns the degeneracy of the $\mathrm{e}_{g}$ orbitals (bonding or antibonding): it was not possible to establish this just from the shapes of the orbitals that were already determined (2-23 and $\mathbf{2 - 2 7}$ for the bonding MO, 2-24 and 2-28 for the antibonding MO). Three nonbonding MO of $\mathrm{t}_{2 g}$ symmetry are found at an intermediate energy level. They are pure metal $d$ orbitals ( $x y, x z$, and $y z$ ), which remain nonbonding as there are no orbitals of the same symmetry on the $\mathrm{L}_{6}$ fragment.

This general scheme, with six bonding MO, three nonbonding MO, and six antibonding MO is appropriate for all octahedral complexes, since it is based on the symmetry properties of the fragment orbitals. Moreover, the lowest antibonding MO (situated above the nonbonding $\mathrm{t}_{2 g}$ orbitals) are always the two degenerate $2 \mathrm{e}_{g}$ orbitals, since they involve the metal $d$ orbitals, which are lower in energy than the $s$ and $p$ orbitals which contribute to the antibonding MO $2 \mathrm{a}_{1 g}$ and $2 \mathrm{t}_{1 u}$, respectively. However, the energetic ordering given in Figure 2.2 for both the bonding and antibonding MO can depend on the nature of the metal and the ligands. For example, the bonding level $1 \mathrm{e}_{g}$ can be found between the bonding $1 \mathrm{a}_{1 g}$ and $1 \mathrm{t}_{1 u}$ levels, or the antibonding $2 \mathrm{a}_{1 g}$ may be placed lower than the antibonding $2 \mathrm{t}_{1 u}$ level.

## Comment

There are six valence atomic orbitals on the metal which contribute to the formation of the six metal-ligand bonds: the $s$ orbital, the three $p$ orbitals, and two of the $d$ orbitals. This result may be compared with that given by hybridization theory, where one forms six equivalent hybrids, directed towards the six vertices of an octahedron. One does indeed talk of $d^{2} s p^{3}$ hybridisation, each hybrid orbital being a suitable linear combination of two $d$ orbitals, the $s$ orbital, and the three $p$ orbitals.

octahedral $\mathrm{ML}_{6}$ complex in the orientation given in Scheme 2-2.

### 2.1.2.4. Molecular orbitals and the definition of the d block

The shapes of the molecular orbitals may be obtained without difficulty from the shapes of the fragment orbitals (Figure 2.3). Given the greater electronegativity of the ligands, the bonding MO are more concentrated on them, whereas the antibonding MO are more concentrated on the metal. Among these MO, seven involve contributions from the metal $d$ orbitals: the two bonding MO $1 \mathrm{e}_{g}$, the three nonbonding $\mathrm{t}_{2 g} \mathrm{MO}$ and the two antibonding MO $2 \mathrm{e}_{g}$. The term ' d block' of an octahedral complex (and also for other types of complex) is used to refer to the five molecular orbitals that are mainly concentrated on the d orbitals of the metal. The three nonbonding MO of $\mathrm{t}_{2 g}$ symmetry clearly belong to the d block, since they are formed from pure $d$ orbitals $(x y, x z$, and $y z)$. For the MO of $\mathrm{e}_{g}$ symmetry, the bonding combinations ( $1 \mathrm{e}_{g}$ ) are mainly concentrated on the ligands while the largest coefficient for the antibonding combinations $\left(2 \mathrm{e}_{g}\right)$ is on the metal. It is these last two orbitals which are considered to belong to the d block. The d block of an octahedral complex is therefore made up of three nonbonding degenerate ( $\mathrm{t}_{2 g}$ ) MO and two antibonding MO which are also degenerate ( $2 \mathrm{e}_{g}$ ). These latter orbitals are often represented as $x^{2}-y^{2}$ and $z^{2}$, with the antibonding contributions on the ligands being implicitly taken into account in this simplified notation.

### 2.1.2.5. Two other representations of the d block

The orbitals of the d block, presented in the way we have just established, correspond to the arrangement of the ligands that is indicated in 2-2: two ligands placed on each of the $x$-, $y$-, and $z$-axes. The three nonbonding orbitals are then $x y, x z$, and $y z$, and the two antibonding orbitals are $x^{2}-y^{2}$ and $z^{2}$. But other arrangements of the ligands with respect to the axes are of course possible. Two in particular will be useful to us in what follows: (i) where four metal-ligand bonds bisect the $x$ - and $y$-axes (2-29); (ii) where the $z$-axis is a threefold symmetry axis $\left(C_{3}\right)$ of the octahedron (2-30). It is clear that a change in the axis system cannot have any consequence for the structure of the $d$ block, which consists of three degenerate nonbonding orbitals and two antibonding orbitals that are also degenerate. However, the expressions of some of these orbitals in the new axes systems do change.



2-30

Figure 2.4. Representation of the orbitals of the d block for an octahedral complex $\mathrm{ML}_{6}$ in the orientation shown in Scheme 2-29.


$x y$




We examine first the situation represented in 2-29. The six ligands are still placed in one or other of the two nodal planes of the $x z$ and $y z$ orbitals. These two orbitals are therefore still nonbonding (Figure 2.4), as they were in the preceding case (Figure 2.3). Nor is there any change for the antibonding orbital $z^{2}$, since there are still two ligands on the $z$-axis and four in the $x y$ plane. However, the $x y$ orbital which was nonbonding now points towards all four ligands placed on the bisectors of the $x$ - and $y$-axes: this orbital becomes one of the antibonding orbitals of the d block. On the other hand, these four ligands are now situated in the nodal planes of the $x^{2}-y^{2}$ orbital, which becomes, with $x z$ and $y z$, the third nonbonding orbital of the d block. This new representation of the $d$ block (Figure 2.4) may be distinguished from the preceding one (Figure 2.3) by 'exchange' of the $x y$ and $x^{2}-y^{2}$ orbitals.

Analysis of the situation is more complex for the geometrical arrangement shown in $\mathbf{2 - 3 0}$. The only orbital for which the result is obtained easily is $z^{2}$. The angle between the M-L bonds and the $z$-axis as it is now defined is exactly equal to $109.5^{\circ}$, which is the value of the angle of the nodal plane of the $z^{2}$ orbital (Chapter 1, Scheme 1-31). As the six ligands are located on the nodal surface of this orbital, no interaction of $\sigma$ type is possible, so $z^{2}$ is one of the three nonbonding orbitals. The difficulty in finding the two other nonbonding orbitals arises because none of the four ligands $\mathrm{L}_{1}-\mathrm{L}_{4}$ is now placed in one of the nodal planes of the $x y, x z, y z$, or $x^{2}-y^{2}$ orbitals. None of these orbitals


2-31 is therefore nonbonding. Nor is any as antibonding as the $2 \mathrm{e}_{g}$ orbitals were in the preceding representations, since they do not point directly towards the ligands. An example is given in 2-31 for the $x z$ orbital. The ligands $\mathrm{L}_{5}$ and $\mathrm{L}_{6}$ are in the nodal plane $y z$ (coefficients zero). The orbitals on the ligands $\mathrm{L}_{1}-\mathrm{L}_{4}$ interact with $x z$, though they are not placed on the bisectors of the $x$ - and $z$-axes where the amplitude of this orbital is greatest.

Figure 2.5. Representation of the orbitals of the $d$ block of the octahedron when the $z$-axis coincides with a $C_{3}$ axis (Scheme 2-30).







In fact, rather than considering the metal $d$ orbitals individually, we must form new combinations that either 'avoid' the six ligands (nonbonding orbitals) or point towards them (antibonding orbitals). These linear combinations, which allow the original orbitals to be reoriented, are indicated below and shown graphically in Figure 2.5.

$$
\begin{align*}
\phi_{1} & =\sqrt{\frac{2}{3}}\left(x^{2}-y^{2}\right)-\sqrt{\frac{1}{3}} y z  \tag{2.1}\\
\phi_{2} & =\sqrt{\frac{2}{3}} x y-\sqrt{\frac{1}{3}} x z  \tag{2.2}\\
\phi_{3} & =\sqrt{\frac{1}{3}}\left(x^{2}-y^{2}\right)+\sqrt{\frac{2}{3}} y z  \tag{2.3}\\
\phi_{4} & =\sqrt{\frac{1}{3}} x y+\sqrt{\frac{2}{3}} x z \tag{2.4}
\end{align*}
$$

The first two combinations ( $\phi_{1}$ and $\phi_{2}$ ) correspond to two orbitals which, with $z^{2}$, form the $t_{2 g}$ block. The last two ( $\phi_{3}$ and $\phi_{4}$ ), on the other hand, have maximum overlap with the ligands and are the two antibonding $\mathrm{e}_{\mathrm{g}}$ orbitals (Figure 2.5).

### 2.1.2.6. Weak fields and strong fields

In the absence of any ligand (isolated metal atom), the five $d$ orbitals of course all have the same energy. One can therefore represent the effect of complexation on the $d$ orbitals by Scheme 2-32, where the $d$ block of the complex is considered as defined earlier. The fivefold degeneracy of the $d$ orbitals is lifted, to give three nonbonding orbitals ( $\mathrm{t}_{2 g}$ ) and two antibonding orbitals ( $\mathrm{e}_{\mathrm{g}}$ ).


2-32

The energy separation $\left(\Delta_{0}\right)$ between the $t_{2 g}$ orbitals (nonbonding) and the $\mathrm{e}_{g}$ orbitals (antibonding) depends on the strength of the $\sigma$ interactions between the metal and the ligands. The value of this energy gap allows us to distinguish, in the family of octahedral complexes, strongfield (large $\Delta_{0}$ ) from weak-field complexes (small $\Delta_{0}$ ) (see Chapter 1, $\S$ 1.6.2). Measurement of the energies of $d-d$ transitions allows us to estimate the value of $\Delta_{0}$ in a large number of complexes, and to establish a spectrochemical series, in which the ligands are ranked according to the strength of the field (value of $\Delta_{0}$ ) that they create:

$$
\begin{aligned}
\mathrm{I}^{-} & <\mathrm{Br}^{-}<\mathrm{Cl}^{-}<\mathrm{F}^{-}<\mathrm{OH}^{-}<\mathrm{O}^{2-}<\mathrm{H}_{2} \mathrm{O}<\mathrm{NH}_{3} \\
& <\mathrm{NO}_{2}^{-}<\mathrm{CH}_{3}^{-}<\mathrm{C}_{6} \mathrm{H}_{5}^{-}<\mathrm{CN}^{-}<\mathrm{CO}
\end{aligned}
$$

Organometallic complexes, which contain one or several metal-carbon bonds, are thus strong-field complexes.

## Comment

The value of $\Delta_{0}$, and therefore the order of the ligands in the spectrochemical series, does not depend only on $\sigma$ interactions. We shall see in Chapter 3 ( $\S 3.3 .4 .2$ ) that the presence of $\pi$ interactions can either decrease the value of $\Delta_{0}$ (' $\pi$-donor' ligands such as $\mathrm{I}^{-}$or $\mathrm{Br}^{-}$) or increase it (' $\pi$-acceptor' ligands such as $\mathrm{CN}^{-}$or CO ).

### 2.1.3. Electronic structure

### 2.1.3.1. $d^{6}$ diamagnetic complexes

The notion that a complex is stable if all its bonding and nonbonding MO are doubly occupied is verified for many octahedral complexes. The six bonding MO, which form the six metal-ligand bonds, and the
 three nonbonding $\mathrm{MO}\left(\mathrm{t}_{2 g}\right)$ of the d block are thus doubly occupied, giving diamagnetic complexes (all the electrons are paired) with 18 electrons. As six electrons occupy the $d$ block, the electronic configuration of these complexes is written as $d^{6}$ or $\left(\mathrm{t}_{2 g}\right)^{6}(\mathbf{2 - 3 3})$.

If we limit ourselves to octahedral complexes with six identical ligands that only have $\sigma$ interactions with the metal (the model
studied in this chapter), examples include the complexes $\left[\mathrm{Fe}\left(\mathrm{H}_{6}\right)\right]^{4-}$, $\left[\mathrm{Ru}\left(\mathrm{NH}_{3}\right)_{6}\right]^{2+},\left[\mathrm{Co}\left(\mathrm{NH}_{3}\right)_{6}\right]^{3+},\left[\mathrm{Rh}\left(\mathrm{NH}_{3}\right)_{6}\right]^{3+}$, and $\left[\operatorname{Ir}\left(\mathrm{NH}_{3}\right)_{6}\right]^{3+}$ which are all 18 -electron complexes, whose electronic configuration is $d^{6}$ (complexes of $\mathrm{Fe}(\mathrm{II}), \mathrm{Ru}(\mathrm{II}), \mathrm{Co}(\mathrm{III}), \mathrm{Rh}(\mathrm{III})$, and $\mathrm{Ir}(\mathrm{III})$, respectively). The presence of ligands that can have $\pi$-type interactions with the metal will of course modify the interaction scheme shown in Figure 2.2, but will not change it drastically. Thus, the complexes $\left[\mathrm{M}(\mathrm{CO})_{6}\right](\mathrm{M}=\mathrm{Cr}$, Mo, W), $\left[\operatorname{Re}(\mathrm{CO})_{6}\right]^{+},\left[\mathrm{M}(\mathrm{CN})_{6}\right]^{5-}(\mathrm{M}=\mathrm{Mn}, \mathrm{Re}),\left[\mathrm{M}(\mathrm{CN})_{6}\right]^{4-}(\mathrm{M}=\mathrm{Fe}$, $\mathrm{Os}),\left[\mathrm{Co}(\mathrm{CN})_{6}\right]^{3-},\left[\mathrm{Mn}(\mathrm{CNR})_{6}\right]^{+},\left[\mathrm{Fe}(\mathrm{CNR})_{6}\right]^{2+},\left[\mathrm{M}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{3+}(\mathrm{M}=\mathrm{Co}$, $\mathrm{Rh}, \mathrm{Ir}),\left[\mathrm{MCl}_{6}\right]^{3-}(\mathrm{M}=\mathrm{Rh}, \mathrm{Ir})$ are also diamagnetic $d^{6}$ complexes. In other complexes, the presence of different ligands can, from the rigorous perspective of group theory, lower the symmetry of the complex considerably: the bond lengths are no longer all equal and the bond angles can deviate from the ideal value of $90^{\circ}$. However, the orbital scheme of the octahedron is essentially preserved (one may speak of 'pseudo-octahedral' symmetry), in particular as far as the main features of the d block are concerned: there are three energy levels below two, even if, within each of these groups, the strict degeneracy due to symmetry has disappeared. Complexes such as $\left[\mathrm{M}(\mathrm{CO})_{5}\right]^{-}$ $(\mathrm{M}=\mathrm{Cr}, \quad \mathrm{Mo}), \quad\left[\mathrm{M}(\mathrm{CO})_{5} \mathrm{Cl}\right] \quad(\mathrm{M}=\mathrm{Mn}, \quad \mathrm{Re}), \quad\left[\mathrm{Ru}(\mathrm{CO})_{4}(\mathrm{Cl})\left(\mathrm{CH}_{3}\right)\right]$, $\left[\operatorname{Ir}(\mathrm{CO})_{2}\left(\mathrm{PR}_{3}\right)(\mathrm{I})(\mathrm{Cl})(\mathrm{Me})\right], \quad\left[\mathrm{Mo}\left(\mathrm{PR}_{3}\right)_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)_{2}\right]$, the molecular hydrogen complex $\left[\mathrm{W}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{CO})_{3}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$ and even ferrocene $\left[\mathrm{Fe}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\right]$ (Chapter 1, Scheme 1-7) are all pseudo-octahedral diamagnetic complexes, whose electronic configuration is $d^{6}$.

### 2.1.3.2. Other cases

There are octahedral complexes that have fewer than 18 electrons. $d^{0}$ complexes in which the d block is completely empty are an extreme case: examples include $\left[\mathrm{MF}_{6}\right]$ complexes ( $\mathrm{M}=\mathrm{Cr}, \mathrm{Mo}$ ), $\left[\mathrm{WCl}_{6}\right],\left[\mathrm{MF}_{6}\right]^{-}(\mathrm{M}=\mathrm{V}, \mathrm{Ta}),\left[\mathrm{MCl}_{6}\right]^{-}(\mathrm{M}=\mathrm{Ta}, \mathrm{Nb}),\left[\mathrm{MF}_{6}\right]^{2-}(\mathrm{M}=\mathrm{Ti}$, $\mathrm{Zr}),\left[\mathrm{ZrCl}_{6}\right]^{2-},\left[\mathrm{Mo}(\mathrm{OMe})_{6}\right]$, and the organometallic $\mathrm{Ti}(\mathrm{IV})$ complex $\left[\mathrm{Ti}\left(\mathrm{PR}_{3}\right)_{2}\left(\mathrm{Cl}_{3}\left(\mathrm{CH}_{3}\right)\right]\right.$. These are very electron-deficient complexes, since formally there are only the 12 electrons that form the 6 metal-ligand bonds. However, we note that the ligands usually have one or two lone pairs that are not involved in the metal-ligand bond (halide ligands, for example). These lone pairs play an important role in stabilizing complexes that apparently have too few electrons (Chapter 3, § 3.5). Another interesting case concerns complexes in which only three electrons occupy the d block ( 15 -electron complexes whose electronic configuration is $d^{3}$ ). In the ground state, one electron occupies each of the orbitals in the $t_{2 g}$ group, and the three electrons have parallel spin (Hund's rule). As typical examples, we mention $\left[\mathrm{Cr}\left(\mathrm{NH}_{3}\right)_{6}\right]^{3+},\left[\mathrm{Mo}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{3+}$, $\left[\mathrm{V}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+},\left[\mathrm{MF}_{6}\right](\mathrm{M}=\mathrm{Rh}, \mathrm{Ir}),\left[\mathrm{RuF}_{6}\right]^{-},\left[\mathrm{ReF}_{6}\right]^{2-},\left[\mathrm{MCl}_{6}\right]^{2-}$
$(\mathrm{M}=\mathrm{Mn}, \mathrm{Re}),\left[\mathrm{MoCl}_{6}\right]^{3-}$ and the organometallic complexes $\left[\mathrm{V}(\mathrm{CN})_{6}\right]^{4-}$, $\left[\mathrm{TaCl}_{2}(\mathrm{dmpe})_{2}\right],\left[\mathrm{MnMe}_{4}(\mathrm{dmpe})\right]$ (dmpe= dimethylphosphino-ethane).

Two other striking aspects of the electronic structure of octahedral complexes may be noted for weak-field complexes ( $\S$ 2.1.2.6). In this case, the energy gap between the nonbonding $\left(\mathrm{t}_{2 g}\right)$ and the antibonding ( $e_{g}$ ) orbitals of the $d$ block is small. When there are four, five, or six electrons to place in these orbitals we do not find them paired in the three orbitals of lowest energy. The ground state corresponds to the configuration in which as many orbitals of the d block as possible are occupied by a single electron, all with parallel spin. The exchange energy arising from this arrangement, together with the reduction of the interelectron repulsion, more than compensates the energy that is necessary to promote one or two electrons from the $t_{2 g}$ level to the $\mathrm{e}_{g}$ ( $\Delta_{0}$ is small). We thus obtain a 'high-spin' complex, which may be contrasted with a 'low-spin' complex in which as many electrons as possible are paired. Thus, the ground-state electronic configuration of high-spin $d^{5}$ complexes is that shown in 2-34, with one electron in each of the five orbitals of the d block $\left(\mathrm{t}_{2 g}^{3} \mathrm{e}_{g}^{2}\right)$, rather than that for the low-spin configuration shown in $2-35\left(\mathrm{t}_{2 g}^{5}\right)$. This latter arrangement is found for strong-field $d^{5}$ complexes such as $\left[\mathrm{Mn}(\mathrm{CN})_{6}\right]^{4-}$, $\left[\mathrm{Fe}(\mathrm{CN})_{6}\right]^{3-},\left[\mathrm{Mo}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\right]^{+}$, or $\left[\operatorname{Ir}(\mathrm{Cl})_{4}\left(\mathrm{PR}_{3}\right)_{2}\right]$. Many $d^{6}$ complexes are diamagnetic (electronic configuration $\left(\mathrm{t}_{2 g}\right)^{6}$, § 2.1.3.1). But high-spin $d^{6}$ complexes also exist, such as $\left[\mathrm{Fe}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ and $\left[\mathrm{CoF}_{6}\right]^{3-}$, whose electronic configuration is $\left(\mathrm{t}_{2}^{4} \mathrm{e}_{g}^{2}\right)$, with four unpaired electrons, two in the $t_{2 g}$ block and two in the $e_{g}$ block. The fact that the destabilization of the antibonding levels in weak-field complexes is so small allows us to understand the existence of complexes with more than 6 electrons in the d block, that is, complexes with more than 18 electrons. As examples, we may mention the octahedral complexes $\left[\mathrm{Ni}\left(\mathrm{NH}_{3}\right)_{6}\right]^{2+}$ and $\left[\mathrm{Ni}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}, d^{8}$ complexes that have 20 electrons. It is very rare to find more than 18 electrons in organometallic complexes, as these are strong-field complexes in which the antibonding orbitals are at high energy.


We shall now study other geometrical arrangements (ligand fields) that are frequently met in transition metal complexes. We shall generally limit ourselves to the characterization of the 'structure' of the d block, that is, the shape and the relative energy of the five orbitals
that make it up. Knowledge of this structure will enable us to anticipate the (or those) electronic configuration(s) that will be particularly favoured for a given family of complexes (e.g. $d^{6}$ for octahedral diamagnetic complexes). Moreover, the occupied $d$ orbitals are in general the highest-energy occupied orbitals of the complex (nonbonding or weakly antibonding orbitals), and they therefore play an important role in problems linked to the geometrical structure and the reactivity of complexes (see Chapter 4).

### 2.2. Square-planar $\mathrm{ML}_{4}$ complexes

In a square-planar $\mathrm{ML}_{4}$ complex, the metal is placed in the centre of a square whose corners are occupied by the four ligands. One can therefore consider, at least formally, that a square-planar complex is formed by removing two ligands from an octahedral complex, for example, those situated on the $z$-axis (2-36). To establish the structure of the $d$ block, it is convenient to start from the results already obtained for octahedral complexes.


### 2.2.1. Characterization of the $d$ block

We consider first the three nonbonding orbitals of the octahedron ( $x y$, $x z$, and $y z$ ), where the representation chosen is given in Figure 2.3. These orbitals are nonbonding in the octahedron, as the coefficients for the six ligands are zero by symmetry. Removal of two of these ligands therefore has no effect on either the shape or the energy of these three orbitals, which stay pure nonbonding $d$ orbitals in the square-planar complex $(\mathbf{2 - 3 7 a}-\mathbf{c})$. The $x^{2}-y^{2}$ orbital is antibonding in the octahedron, because of its interactions with the four ligands in the $x y$ plane. In contrast, the coefficients on the ligands located on the $z$-axis are zero by symmetry. Removal of these two ligands therefore has no influence on the shape and energy of this orbital, which remains antibonding in the square-planar complex (2-37d). The only $d$ orbital which is modified on passing from the octahedron to the square plane is the antibonding $z^{2}$ orbital. Now the two main antibonding interactions in the octahedron concern the ligands placed on the $z$-axis. Removal of these ligands therefore leads

Figure 2.6. Derivation of the orbitals of the $d$ block for a square-planar $\mathrm{ML}_{4}$ complex from those of an octahedral $\mathrm{ML}_{6}$ complex.
to a substantial stabilization of this orbital, though it remains weakly antibonding due to the small interactions with the four ligands placed in the $x y$ plane (2-37e).


These changes in energy for the orbitals of the $d$ block are illustrated in Figure 2.6, where the symmetry of the orbitals in the square-planar complex is also indicated (point-group symmetry $D_{4 h}$ ). Notice that while the three nonbonding orbitals $x y, x z$, and $y z$ are degenerate from the energetic point of view, only two of them ( $x z$ and $y z$ ) are degenerate by symmetry ( $\mathrm{e}_{g}$ representation whose dimension is 2 ).

The main difference between the $d$ blocks of octahedral and squareplanar complexes concerns the number of nonbonding or weakly antibonding orbitals: there are three in the former but four in the latter.


### 2.2.2. Electronic structure for 16 -electron $d^{8}$ complexes

The favoured electronic configuration for a square-planar diamagnetic complex $\mathrm{ML}^{4}$ involves the double occupation of the four low-energy orbitals of the d block (2-38). There are indeed many diamagnetic complexes with a $d^{8}$ electronic configuration, such as $\left[\mathrm{Pd}\left(\mathrm{NH}_{3}\right)_{4}\right]^{2+}$, $\left[\mathrm{M}(\mathrm{CN})_{4}\right]^{2-}(\mathrm{M}=\mathrm{Ni}, \mathrm{Pd}),\left[\mathrm{PtCl}_{4}\right]^{2-},\left[\mathrm{PtHBr}\left(\mathrm{PR}_{3}\right)_{2}\right],\left[\mathrm{AuBr}_{4}\right]^{-}$, Vaska's complex $\left[\operatorname{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right]$, and Wilkinson's hydrogenation catalyst $\left[\mathrm{Rh}(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{3}\right]$.

If we include the four doubly occupied MO that form the four metalligand bonds (these are not represented in Figure 2.6), we find that the total number of electrons in a $d^{8}$ complex is 16 . But if we consult the qualitative analysis given in Chapter 1 ( $\S$ 1.6.1), it is at first sight surprising that the 18 -electron rule is not respected in this family of complexes, since the bonding orbitals, the nonbonding, and the weakly antibonding orbital $z^{2}$ are all doubly occupied. This 'lack' of two electrons arises because, despite appearances, there is a nonbonding orbital on the metal that is empty. It is not an orbital from the d block, but the metal's $p_{z}$ orbital whose nodal plane contains the four ligands (2-39). Although nonbonding, this orbital is never occupied: it is an $(n+1) p$ orbital whose energy is substantially higher (several eV ) than that of an $n d$ orbital.

This aspect of their electronic structure means that these complexes are likely to undergo addition reactions, in which their 18 -electron shell is completed by the arrival of new ligands. For example, the addition of a molecule of dihydrogen $\mathrm{H}_{2}$ to Vaska's complex leads to the formation of an octahedral complex (2-40). In this addition product, the oxidation state of the metal is +3 (a $d^{6}$ complex), which, when the 12 electrons associated with the bonds are included, does indeed correspond to an 18 -electron complex. We also note that this addition reaction is accompanied by a change in the metal's oxidation state, from +1 in the reactant $\left(d^{8}\right)$ to +3 in the product $\left(d^{6}\right)$. This is therefore called an oxidative addition reaction, since the oxidation state of the metal has increased. The reverse reaction is called reductive elimination ( $\operatorname{Ir}(\mathrm{III}) \rightarrow \operatorname{Ir}(\mathrm{I})$ ).


2-40

### 2.3. Square-based pyramidal $\mathrm{ML}_{5}$ complexes

In an ML ${ }_{5}$ complex which adopts a square-based pyramidal (SBP) geometry, four ligands ( $L_{1}-L_{4}$ ) are located at the corners of a square which
is the base of the pyramid, while the fifth, or apical, ligand $\left(\mathrm{L}_{5}\right)$, is placed on the summit (or apex) of the pyramid (2-41). The metal centre may, depending on the complex, be either in the basal plane $\left(\theta=90^{\circ}\right)(\mathbf{2 - 4 1 a}$, $\S$ 2.3.1) or above this plane $\left(\theta>90^{\circ}\right)(\mathbf{2 - 4 1 b}, \S 2.3 .2)$.


### 2.3.1. Characterization of the $d$ block (metal in the basal plane)

In the complex shown in 2-41a, all the angles between the apical bond $\mathrm{M}-\mathrm{L}_{5}$ and the basal bonds $\mathrm{M}-\mathrm{L}_{1-4}$ are equal to $90^{\circ}$. This structure may formally be obtained by removing one of the ligands from an octahedral complex ( $\mathrm{L}_{6}, \mathbf{2 - 4 2}$ ). We may therefore establish the MO of the d block for the SBP by starting from those that we already know for an octahedron, following the method previously used for square-planar complexes (§ 2.2.1).



2-42

### 2.3.1.1. Derivation of the $d$ orbitals from those of the octahedron

In the octahedral complex, the coefficients of the $x y, x z, y z$, and $x^{2}-y^{2}$ orbitals are zero for the ligands located on the $z$-axis (left-hand side of 2-37). The removal of one of these ligands therefore makes no change to the shape or energy of these four orbitals: in the $d$ block of the SBP ML5 complex, we therefore find that the three orbitals $x y, x z$, and $y z$ are nonbonding, while $x^{2}-y^{2}$ is a strongly antibonding orbital. This result is just the same as that found when studying the square-planar geometry as a derivative of the octahedron (Figure 2.6). The $z^{2}$ orbital is stabilized by the elimination of one of the two antibonding interactions with the ligands placed on the $z$-axis (2-43). But this stabilization is not as large
as that observed when passing from an octahedral to a square-planar complex, since in this latter case, both the antibonding interactions along the $z$-axis are removed (Figure 2.6).


2-43

These results are shown in Figure 2.7, with the orbitals' symmetry ( $\mathrm{C}_{4 v}$ point group).


Figure 2.7. Derivation of the d-block orbitals for an $\mathrm{ML}_{5}$ complex with an SBP geometry (the metal is in the basal plane) from those of an octahedral complex ML ${ }_{6}$.


polarized $z^{2}$

The orbital shown in 2-43 is in fact stabilized by a bonding interaction with the $p_{z}$ orbital (see Appendix A). In Scheme 2-44 (left-hand side), the two orbitals are shown separated for greater clarity, but one must of course imagine that they are superposed. The way in which they must be combined to obtain a bonding mixture ( $S>0$ ) is not obvious, since one of the orbitals is delocalized over all the centres. We shall examine their overlap if the $p_{z}$ orbital is oriented along negative $z(\mathbf{2 - 4 4 )}$. It can be decomposed into three terms: (i) the overlap between the atomic orbitals $z^{2}$ and $p_{z}$ : this term is zero since the two atomic orbitals involved are located on the same atom; (ii) the overlap between $p_{z}$ and the orbitals on the ligands that define the base of the SBP complex: this term is also zero, since these ligands are placed in the nodal plane ( $x y$ ) of $p_{z}$; (iii) the overlap between $p_{z}$ and the orbital of the axial ligand (placed on the $z$-axis), which is the only non-zero term. It is clear that the choice adopted in 2-44 for the orientation of $p_{z}$ leads to the overlap being positive. But how should this orbital be represented graphically? Where $z$ is negative, the amplitudes of $z^{2}$ and $p_{z}$ have the same sign, so they add. But where $z$ is positive, they have the opposite sign, so they tend to cancel each other. The two lobes along the $z$-axis are therefore not equivalent, as the one along negative $z$ is 'larger' (2-44, right-hand side). The participation of the $p_{z}$ orbital therefore leads to a polarization of the $z^{2}$ orbital towards the vacant site of the octahedron. From the energetic point of view, this mixing stabilizes the orbital, since the polarization reduces the antibonding interaction with the axial ligand, by reducing the size of the lobe of the $z^{2}$ orbital that points towards that ligand.

### 2.3.2. Characterization of the $d$ block (metal out of the basal plane)

In most SBP complexes, the metal is located above the base of the pyramid $\left(\theta>90^{\circ}, \mathbf{2 - 4 1 b}\right)$. This geometrical deformation of the preceding structure leads to a displacement, of the same amplitude, of the ligands $\mathrm{L}_{1}$ and $\mathrm{L}_{3}$ in the $y z$ plane, and of the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{4}$ in the $x z$ plane $(\mathbf{2}-45)$. These movements change the shape and energy of some of the
orbitals of the d block represented in Figure 2.7. We shall consider each of these orbitals in turn.


The $x y$ orbital has its nodes in the $x z$ and $y z$ planes, the two planes in which the movements of the ligands occur. As the ligands remain located in the nodal planes, no interaction is possible with the $x y$ orbital ( $S=0$ ), whose shape and energy are unchanged (a pure $d$ orbital that is strictly nonbonding, 2-46).


No interaction can take place between the $y z$ orbital and the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{4}$ that stay in the nodal plane $x z$ for all values of the angle $\theta$. However, the ligands $\mathrm{L}_{1}$ and $\mathrm{L}_{3}$ leave the nodal plane $x y$ and move into the $y z$ plane in which that $d$ orbital is mainly concentrated. As a result, the orbitals on the ligands $\mathrm{L}_{1}$ and $\mathrm{L}_{3}$ can interact with the $y z$ orbital when $\theta$ is greater than $90^{\circ}$. The bonding combination is a lowenergy orbital (an MO that represents a bond), but the antibonding combination is the new d-block orbital (2-47). The $y z$ orbital, which is nonbonding if $\theta=90^{\circ}$, is therefore destabilized if $\theta$ is larger than $90^{\circ}$, its energy rising as the overlap with the ligand orbitals increases. The greatest destabilization occurs when $\theta=135^{\circ}$, the value for which the ligands $L_{1}$ and $L_{3}$ are in the region where the amplitude of the $y z$ orbital is largest.

${ }^{5}$ Since the angle of the nodal cone for $z^{2}$ is $109.5^{\circ}$, the corresponding value of $\theta$ is $125.25^{\circ}$. However, this is only an approximate value in the present situation, as the $z^{2}$ orbital is polarized by the $p_{z}$ orbital.

But this destabilization of the $y z$ orbital is reduced by a bonding mixture with the $p_{y}$ metal orbital, in a way that is exactly analogous to that described in $\S$ 2.3.1.2 for the $z^{2}$ orbital (see Appendix A). As the overlap between the $y z$ and $p_{y}$ orbitals on the metal is zero, a bonding interaction is obtained if the overlap between $p_{y}$ and the orbitals on ligands $L_{1}$ and $L_{3}$ is positive. The appropriate combination is shown in $\mathbf{2 - 4 8}$. If we imagine the superposition of the two orbitals represented in the left-hand part of this sketch, the grey lobe of $p_{y}$ points towards the orbital on $\mathrm{L}_{3}$ (grey lobe directed towards the metal) and the white lobe towards the orbital on $\mathrm{L}_{1}$ (white lobe pointing towards the metal). This mixture with $p_{y}$ leads to a change in the amplitude of the lobes of the $d$ orbital (2-48). If we consider first the right-hand part of the $y z$ orbital (grey lobe towards the top, white lobe towards the bottom), we must add the grey lobe $p_{y}$ to it (2-48). As a result, the grey lobe of $y z$ is enlarged (the amplitudes add) but the white lobe decreases in size (the amplitudes tend to cancel). On the left-hand side of the orbital, the opposite effect is observed: the amplitude of the white lobe increases in size, as it is added to that of the $p_{y}$ orbital (white lobe pointing towards the left), while the amplitude of the grey lobe is diminished.


The polarization of the $y z$ orbital by $p_{y}$ therefore leads to a reduction in the size of the lobes pointing towards ligands $L_{1}$ and $L_{3}$, but an increase in the opposite direction (2-48). The antibonding character of the $y z$ orbital is thereby reduced, so that in fact it becomes a weakly antibonding orbital which can be occupied in stable complexes.

The $x z$ orbital behaves in just the same way as $y z$, but the roles of the ligands $\left(L_{1}, L_{3}\right)$ and $\left(L_{2}, L_{4}\right)$ are interchanged (2-49). It is now the mixture with the $p_{x}$ orbital which polarizes $x z$ in the direction opposed to the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{4}$. For any value of the angle $\theta$, the $x z$ and $y z$ orbitals are related to each other by a rotation of $90^{\circ}$ around the $z$-axis: they are degenerate by symmetry (e symmetry in the $\mathrm{C}_{4 v}$ point group).

The $z^{2}$ orbital, which had antibonding interactions with all the ligands, is affected by the movement of the four basal ligands towards its nodal cone. If the angle $\theta$ is close to $125^{\circ}$, these ligands are even in this cone, ${ }^{5}$ leading to zero overlap between the ligand orbitals and the $z^{2}$ orbital of the d block. For this latter orbital, the deformation that we are


2-49
studying therefore produces a decrease, even perhaps an elimination, of the antibonding contributions on the basal ligands (2-50). The energy of this orbital is therefore lowered, though it remains higher than that of the nonbonding orbital $x y$, since there is still an antibonding interaction with the apical ligand $\mathrm{L}_{5}$.


2-50
We turn last to the case of the antibonding orbital $x^{2}-y^{2}(\mathbf{2 - 5 1})$. An increase in the angle $\theta$ decreases the antibonding interactions with the orbitals on the four basal ligands, since these are no longer located in the regions where the amplitude of this orbital is greatest. As a result, the $x^{2}-y^{2}$ orbital is stabilized, by an amount which increases with the angle $\theta$. However, this orbital remains the least stable of all the five in the d block.


2-51

In Figure 2.8 we present a schematic correlation diagram, which shows the changes to the shape and the energy of the d-block orbitals during this geometrical deformation, as the angle $\theta$ varies from $90^{\circ}$ to about $110^{\circ}$.

## Comment

It must be noted that a crossing between the energy levels of the orbitals that are destabilized ( $x z$ and $y z$ ), and the $z^{2}$ orbital that is stabilized, can occur if $\theta$ becomes sufficiently large. Since these sets of orbitals have different symmetries (e and $a_{1}$, respectively), they cannot interact even if their energies are very similar or even equal (an 'allowed' energy-level crossing). The

Figure 2.8. Correlation diagram for the d-block orbitals of an $\mathrm{ML}_{5}$ complex with an SBP geometry where the angle between the apical and the basal bonds varies from $90^{\circ}$ to about $110^{\circ}$.

value of the angle $\theta$ for which this crossing occurs depends on the exact nature of the metal and the ligands.

### 2.3.3. Electronic structure and geometry

### 2.3.3.1. $d^{8}$ or $d^{6}$ diamagnetic complexes

In the d block of $\mathrm{ML}_{5}$ complexes which adopt an SBP geometry, there is only one strongly antibonding orbital, $x^{2}-y^{2}$ (Figure 2.8). The four nonbonding or weakly antibonding orbitals ( $x y, x z, y z$, and $z^{2}$ ) are therefore likely to be doubly occupied, leading to diamagnetic complexes with a $d^{8}$ electronic configuration (e.g. $\left[\mathrm{Co}(\mathrm{H})_{5}\right]^{4-},\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-},\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]^{-}$, $\left[\mathrm{Co}(\mathrm{NCPh})_{5}\right]^{+}$, or $\left.\left[\mathrm{PtI}\left(\mathrm{PMe}_{3}\right)_{4}\right]^{+}\right)$. Including the ten electrons associated with the five bonds, these are 18 -electron complexes. There are also diamagnetic complexes whose electronic configuration is $d^{6}$ ( 16 -electron complexes), such as $\left[\mathrm{M}(\mathrm{CO})_{5}\right](\mathrm{M}=\mathrm{Cr}, \mathrm{Mo}, \mathrm{W})$ or $\left[\mathrm{W}(\mathrm{CO})_{4}(\mathrm{CS})\right]$. The presence of an empty low-energy orbital in the d block (the $z^{2}$ orbital, Figure 2.8) confers special properties on these complexes. They have only been isolated in rare-gas matrices, at very low temperatures, and they can easily bind a sixth ligand of the L type to give an octahedral 18-electron complex.

### 2.3.3.2. Other cases

Complexes with an intermediate electronic configuration, $d^{7}$, are also known, such as $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right],\left[\operatorname{Re}(\mathrm{CO})_{5}\right],\left[\operatorname{Cr}(\mathrm{CO})_{5}\right]^{-}$, and $\left[\mathrm{Co}(\mathrm{CN})_{5}\right]^{3-}$. These are radical complexes with 17 electrons; the unpaired electron occupies the $z^{2}$ orbital. They can therefore dimerize, just like organic radicals, to form bimetallic complexes such as $\left[\mathrm{Mn}_{2}(\mathrm{CO})_{10}\right],\left[\mathrm{Re}_{2}(\mathrm{CO})_{10}\right]$, or $\left[\mathrm{Re}_{2}(\mathrm{CN})_{10}\right]^{6-}$, etc. For some electronic
counts, the presence of four orbitals with similar energies can favour the existence of high-spin complexes. For example, the ground-state electronic configuration of the complex $\left[\mathrm{MnCl}_{5}\right]^{2-}\left(d^{4}\right)$ corresponds to the occupation of each of the orbitals $(x y),(x z),(y z)$, and $\left(z^{2}\right)$ by a single electron, with their spins parallel. High-spin complexes with a $d^{6}$ or $d^{8}$ electronic configuration are also known, which implies that the antibonding $x^{2}-y^{2}$ orbital is singly occupied. This can only happen for weak-field complexes. Deoxyhaemoglobin is a well-known example: it is a $d^{6}$ iron complex, whose electronic configuration is $(x y)^{2}(x z)^{1}(y z)^{1}\left(z^{2}\right)^{1}\left(x^{2}-y^{2}\right)^{1}$. We remark in closing this section that there are a few complexes, such as $\left[\mathrm{VOF}_{4}\right]^{-}$or $\left[\mathrm{Nb}\left(\mathrm{NMe}_{2}\right)_{5}\right]$, in which the $d$ block is empty ( $d^{0}$ configuration). As for the octahedral complexes, these very electron-deficient systems are only observed when the ligands possess lone pairs.

### 2.3.3.3. Electronic count and geometry

Correlation diagrams for the d-block orbitals, of the type shown in Figure 2.8, are often used to interpret the changes in the structures of complexes as a function of their electronic configuration $d^{n}$. In order to use them, one adopts the hypothesis, which in most cases is verified a posteriori, that the geometry of the complex is controlled by the energy changes of the highest-occupied molecular orbital (the HOMO rule). ML; complexes with an SBP geometry provide an interesting illustration of this rule.

## Comment

The structures observed for complexes whose electronic configuration is $d^{0}$ (d-block empty) are characterized by values of the angle $\theta$ greater than $90^{\circ}$ (metal above the basal plane). This preference, which is caused by bonding orbitals in the complex, can, as we shall see, be either modified or maintained when there are electrons in the d block.

We consider first a diamagnetic $d^{6}$ complex. In the structure with $\theta=90^{\circ}$, three nonbonding orbitals ( $x y, x z$, and $y z$ ) are doubly occupied (Figure 2.8). An increase in the value of $\theta$ above $90^{\circ}$ leads to a destabilization of two of these orbitals ( $x z$ and $y z$ ), which is energetically unfavourable (Figure 2.8). We may therefore predict that these complexes will adopt a structure in which the metal stays in the basal plane $\left(\theta=90^{\circ}\right)$. Experimental values for these complexes are indeed close to $90^{\circ}\left(91-94^{\circ}\right.$ for the $\left[\mathrm{M}(\mathrm{CO})_{5}\right]$ complexes if $\left.\mathrm{M}=\mathrm{Cr}, \mathrm{W}\right)$. But in a $d^{8}$ diamagnetic complex, the two additional electrons occupy the $z^{2}$ orbital, which is stabilized when $\theta$ increases above $90^{\circ}$ (Figure 2.8). Application of the HOMO rule leads to the prediction that in these complexes, the
metal will be located above the basal plane. This change in geometry due to the number of $d$ electrons is indeed observed: the angle $\theta$ is $101.0^{\circ}$ in $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$ and $102.6^{\circ}$ in $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]^{-}$.

In the light of this analysis, we present the energy levels and orbital occupations for the d block of $d^{6}$ and $d^{8}$ diamagnetic complexes in Scheme 2-52.


We can extend this analysis to other electron counts. In $d^{7}$ complexes, the $z^{2}$ orbital is singly occupied. We may therefore expect to find that $\theta$ is larger than $90^{\circ}$, but not as large as in the low-spin $d^{8}$ complexes in which this orbital is doubly occupied. The values of $\theta$ for $d^{7}$ complexes are indeed often intermediate between those found for low-spin $d^{6}$ and $d^{8}$ complexes: about $95^{\circ}$ for $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right],\left[\operatorname{Re}(\mathrm{CO})_{5}\right]$, or $\left[\mathrm{Cr}(\mathrm{CO})_{5}\right]^{-}$, $97.6^{\circ}$ for $\left[\mathrm{Co}(\mathrm{CN})_{5}\right]^{3-}$.

For a given number of electrons, the change from a low-spin to a high-spin complex can also lead to a change in geometry, since the HOMO is different. In a diamagnetic $d^{6}$ complex, whose electronic configuration is $(x y)^{2}(x z)^{2}(y z)^{2}$, we have seen that the angle $\theta$ is close to $90^{\circ}$. However, in a high-spin $d^{6}$ complex, whose electronic configuration is $(x y)^{2}(x z)^{1}(y z)^{1}\left(z^{2}\right)^{1}\left(x^{2}-y^{2}\right)^{1}$, the HOMO, $x^{2}-y^{2}$, is strongly stabilized by an increase in $\theta$ (Figure 2.8), thereby favouring a displacement of the metal out of the basal plane. This is exactly what is observed in deoxyhaemoglobin (a high-spin complex of Fe (II)), where the iron atom is placed well above the plane $\left(\theta=110^{\circ}\right)$ defined by the porphyrin ring (a ligand of $\mathrm{L}_{2} \mathrm{X}_{2}$ type).

### 2.4. Tetrahedral $\mathrm{ML}_{4}$ complexes

In a complex of this type, the metal is placed at the centre of a tetrahedron whose vertices are occupied by the four ligands. There are at least two ways to derive the d-block orbitals for a tetrahedral complex. In the 'direct' method, we allow the $d$ orbitals of the metal to interact with the symmetry-adapted combination of ligand orbitals (Chapter 6, § 6.6.2.2
and Exercise 2.11), taking account of their symmetry properties in the $T_{d}$ group. We may also start from the $d$ orbitals of a square-planar complex and study their changes as the ligands are moved towards the tetrahedral geometry. We shall use the second approach here, as it not only gives us the d-block structure for a tetrahedral complex, but it also yields the orbital correlation diagram linking the two geometries most frequently found for $\mathrm{ML}_{4}$ complexes.



2-53

The change from one structure to the other is studied following the mechanism described in Scheme 2-53, where the ligands $L_{1}$ and $L_{3}$ move upwards in the $y z$ plane, and the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{4}$ downwards in the $x z$ plane. In this way, we change from a square-planar ( $\mathrm{D}_{4 h}$ symmetry) to a tetrahedral complex ( $\mathrm{T}_{d}$ symmetry), maintaining $D_{2 d}$ symmetry at intermediate points.

## Comment

This is not the only way to move from one structure to the other. One can, for example, rotate one $\mathrm{ML}_{2}$ unit by $90^{\circ}$ with respect to the other, progressively adjusting the values of the bond angles ( $D_{2}$ symmetry is conserved).

### 2.4.1. Characterization of the d block

We shall consider each of the five $d$ orbitals of a square-planar complex as they are presented in Figure 2.6 (right-hand side). On moving to the tetrahedral structure, the $x y$ orbital stays unchanged in both shape and energy ( $\mathbf{2 - 5 4 a}$ ), as the four ligands move in one or other of the two nodal planes of this orbital ( $\mathrm{L}_{2}$ and $\mathrm{L}_{4}$ in $x z, \mathrm{~L}_{1}$ and $\mathrm{L}_{3}$ in $y z$ ). The $y z$ and $x z$ orbitals behave in just the same way as they do when the metal moves out of the basal plane in SBP ML ${ }_{5}$ complexes (§ 2.3.2). On passing from the square plane to the tetrahedron, two ligands move in the plane where the amplitude of the orbital is greatest, but the other two ligands stay in one of the nodal planes. As a result, the $y z$ orbital is destabilized by antibonding interactions with the ligands $\mathrm{L}_{1}$ and $\mathrm{L}_{3}$, but it has no interaction with $L_{2}$ and $L_{4}$ which are still located in the nodal plane $x z(\mathbf{2}-\mathbf{5 4 b})$. In the same way, the $x z$ orbital is destabilized by an interaction with the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{4}(\mathbf{2 - 5 4 c})$. These two orbitals, $y z$
and $x z$, which were nonbonding in a square-planar complex, therefore become antibonding in a tetrahedral complex, while staying degenerate. However, just as in ML 5 complexes ( $\S 2.3 .2$ ), the antibonding interactions are reduced by mixing with a $p$ orbital on the metal ( $p_{y}$ in the case of $y z, p_{x}$ for $x z$ ) which polarizes the $d$ orbital in the direction opposite to that of the two ligands with which it is interacting $(\mathbf{2 - 5 4 b})$ and see also Appendix A). The $z^{2}$ orbital becomes rigorously nonbonding (zero coefficients on the ligands, $\mathbf{2 - 5 4 d}$ ), and thus at the same energy as the $x y$ orbital, since the four ligands in the tetrahedral structure are moved precisely on to its nodal cone (whose angle is exactly equal to $109.5^{\circ}$ ). The last orbital, $x^{2}-y^{2}$, is strongly stabilized since the four ligands leave the $x$ - and $y$-axes on which they were located. The antibonding interactions are decreased but not eliminated, since in the tetrahedral structure, the ligands are not in the nodal planes of the $x^{2}-y^{2}$ orbital (2-54e).


Table 2.2. Character table for the $T_{d}$ group (tetrahedral $\mathrm{ML}_{4}$ complex)

| $T_{d}$ | $E$ | $8 C_{3}$ | $3 C_{2}$ | $6 S_{4}$ | $6 \sigma_{d}$ |  |  |
| :---: | :---: | ---: | ---: | ---: | ---: | :--- | :---: |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 | 1 | 1 | $x^{2}+y^{2}+z^{2}$ |  |
| $\mathrm{~A}_{2}$ | 1 | 1 | 1 | -1 | -1 |  |  |
| E | 2 | -1 | 2 | 0 | 0 |  | $\left(2 z^{2}-x^{2}-y^{2}, x^{2}-y^{2}\right)$ |
| $\mathrm{T}_{1}$ | 3 | 0 | -1 | 1 | -1 |  |  |
| $\mathrm{~T}_{2}$ | 3 | 0 | -1 | -1 | 1 | $(x, y, z)$ | $(x y, x z, y z)$ |

Inspection of the last column shows that the five $d$ orbitals are found in two groups of degenerate orbitals, of e (doubly degenerate) and $t_{2}$ (triply degenerate) symmetry, respectively. As we have already shown that there are just two nonbonding orbitals, it is clear that they must have e symmetry. The antibonding orbitals therefore have $t_{2}$ symmetry, and all three are degenerate by symmetry. The $d$ block of the tetrahedron therefore contains two nonbonding degenerate orbitals (e) and three antibonding orbitals that constitute another degenerate set $\left(\mathrm{t}_{2}\right)$.

## Comment

We note that the roles of $x y$ and $x^{2}-y^{2}$ orbitals are interchanged in a tetrahedral complex, depending on whether we adopt the analysis developed above ( $x y$ nonbonding (2-54a), $x^{2}-y^{2}$ antibonding (2-54e)), or the character table ( $x^{2}-y^{2}$ degenerate with $z^{2}$, therefore nonbonding, $x y$ degenerate with $x z$ and $y z$, therefore antibonding). This arises simply from a different definition of the axes (a rotation of $45^{\circ}$ around the $z$-axis) which interchanges the role of the $x y$ and $x^{2}-y^{2}$ orbitals, as we have already seen in the case of the octahedron (Figures 2.3 and 2.4) or the square-planar complex plan (see Exercise 2.4).

We can now sketch the correlation diagram linking the d-block orbitals of a square-planar complex to those of a tetrahedron (Figure 2.9), following the deformation shown in 2-53.


Figure 2.9. Correlation diagram linking the d-block orbitals of a square-planar $\mathrm{ML}_{4}$ complex and those of a tetrahedral $\mathrm{ML}_{4}$ complex, following the deformation shown in 2-53.

### 2.4.2. Electronic structure

### 2.4.2.1. $d^{10}$ diamagnetic complexes

The separation between the $e$ and $t_{2}$ levels of the $d$ block is sufficiently

e
 small (the $t_{2}$ orbitals are only weakly antibonding) that all five orbitals can be occupied. Diamagnetic complexes with a $d^{10}$ electronic configuration are thus obtained (2-55), such as $\left[\mathrm{Ni}(\mathrm{H})_{4}\right]^{4-},\left[\mathrm{Ni}(\mathrm{CO})_{4}\right]$, $\left[\mathrm{M}\left(\mathrm{PF}_{3}\right)_{4}\right] \quad(\mathrm{M}=\mathrm{Ni}, \quad \mathrm{Pd}), \quad[\mathrm{Pt}(\mathrm{dppe}) 2] \quad$ (dppe= diphenylphosphinoethane), $\left[\mathrm{Ni}(\mathrm{CN})_{4}\right]^{4-},\left[\mathrm{Co}(\mathrm{CO})_{4}\right]^{-},\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]^{2-},\left[\mathrm{Cu}(\mathrm{CN})_{4}\right]^{3-}$, $\left[\mathrm{Cu}\left(\mathrm{PMe}_{3}\right)_{4}\right]^{+},\left[\mathrm{Ag}\left(\mathrm{PPh}_{3}\right)_{4}\right]^{+}$, or $\left[\mathrm{Zn}(\mathrm{Cl})_{4}\right]^{2-}$. These are 18 -electron complexes, since eight additional electrons are associated with the four metal-ligand bonds.

### 2.4.2.2. Other cases

$\mathrm{ML}_{4}$ complexes whose d block is empty ( $d^{0}$ electronic configuration), such as $\left[\mathrm{TiCl}_{4}\right]$ and $\left[\mathrm{MnO}_{4}\right]^{-}$, also adopt a tetrahedral geometry. As in the examples of $d^{0}$ complexes already mentioned ( $\S$ 2.1.3.2 and 2.3.3.2), the ligands have lone pairs which play an important role in the stabilization of these species that apparently are strongly electron-deficient (formally, only eight electrons around the metal!). We may also mention tetrahedral $d^{8}$ complexes, such as $\left[\mathrm{NiCl}_{4}\right]^{2-},\left[\mathrm{Ni}\left(\mathrm{PPh}_{2} \mathrm{Et}_{2}\right)_{2} \mathrm{Br}_{2}\right]$, or $\left[\mathrm{CoBr}\left(\mathrm{PR}_{3}\right)_{3}\right]$. The structure of the d block, with the three $\mathrm{t}_{2}$ degenerate orbitals, leads to a paramagnetic ground state $\left(e^{4} t_{2}^{4}\right)$, with two unpaired electrons in the $t_{2}$ orbitals (a high-spin tetrahedral complex, right-hand side of 2-57).

### 2.4.3. $\mathrm{ML}_{4}$ complexes: square-planar or tetrahedral?

Not all $\mathrm{ML}_{4}$ complexes have a square-planar or tetrahedral structure. 'Intermediate' geometries, of lower symmetry, can be observed (e.g. $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$ is a high-spin $d^{8}$ complex with $C_{2 v}$ symmetry, see $\S$ 2.8.3). However, a large number of $\mathrm{ML}_{4}$ complexes adopt, either exactly or with only small deviations, one or other of these high-symmetry geometries. In this paragraph, we shall show how a knowledge of the d block for each type of structure ( $\S 2.2 .1$ and 2.4.1) enables us to establish a link between the $d^{n}$ electronic configuration of the complex and the geometry that is observed experimentally. We shall concentrate our attention on the $d^{10}$ and $d^{8}$ electronic configurations, which are the most common for $\mathrm{ML}_{4}$ complexes.

It is straightforward to understand the structural preference for $d^{10}$ complexes: in the tetrahedral geometry, all five $d$ orbitals are low in energy and can thus receive ten electrons, whereas there are only four low-energy $d$ orbitals in the square-planar geometry

[^1](Figure 2.9). Complexes with a $d^{10}$ electronic configuration therefore adopt a tetrahedral structure ( $\$$ 2.4.2.1). ${ }^{6}$

In diamagnetic $d^{8}$ complexes (low-spin), four $d$ orbitals must be doubly occupied. This count is ideal for the square-planar geometry, characterized by four low-energy $d$ orbitals (three strictly nonbonding, one very weakly antibonding ( $\mathbf{2 - 5 6}$, left-hand side)), well separated from the fifth which is strongly antibonding. The situation is less favourable for the tetrahedral geometry: first, there are only four nonbonding electrons instead of six, and second, the distribution of four electrons in three orbitals ( $\mathbf{2 - 5 6}$, right-hand side) is energetically unfavourable since it does not obey Hund's rule (two electrons should occupy two different orbitals, with parallel spins). In summary, the d-block of a tetrahedral complex is not well suited to accomodate four pairs of electrons, and in fact diamagnetic $d^{8}$ complexes adopt a square-planar geometry $\left(\left[\mathrm{Ni}(\mathrm{CN})_{4}\right]^{2-},\left[\mathrm{Pd}\left(\mathrm{NH}_{3}\right)_{4}\right]^{2+},\left[\mathrm{RhCl}\left(\mathrm{PPh}_{3}\right)_{3}\right]\right.$, or $\left[\operatorname{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right]$, for example).


Square-planar ( $d^{8}$ low-spin)


Tetrahedral ( $d^{8}$ low-spin)

This leads us naturally to consider the case of paramagnetic $d^{8}$ complexes (high-spin), that have two unpaired electrons with parallel spin. The situation is now favourable for the tetrahedral geometry: Hund's rule is obeyed, and the occupied $d$ orbitals are either nonbonding or only weakly antibonding ( $\mathbf{2}-57$, right-hand side). In contrast, in the square-planar geometry, an electron must be placed in the strongly antibonding orbital of the d block ( $\mathbf{2}-57$, left-hand side). The tetrahedral structure is therefore favoured for high-spin $d^{8}$ complexes ( $\S$ 2.4.2.2). As a consequence, the low-spin $\rightarrow$ high-spin change in $d^{8} \mathrm{ML}_{4}$ complexes is accompanied with a change in geometry, from square-planar to tetrahedral.

It is more difficult to predict whether a given complex with a $d^{8}$ electronic configuration will be low-spin, with a square-planar structure, or high-spin, with a tetrahedral structure. From a purely electronic point of view, the square-planar geometry seems to be favoured, with six nonbonding and two weakly antibonding electrons, instead of four nonbonding and four more strongly antibonding electrons in the tetrahedron (2-58). However, in this latter structure, the exchange energy due to the two unpaired electrons is a favourable factor. Steric factors


2-57 Square planar ( $d^{8}$ high-spin) Tetrahedral ( $d^{8}$ high-spin)
also play an important role, since they favour the tetrahedral structure, with bond angles of $109.5^{\circ}$ rather than the square-planar arrangement with $90^{\circ}$ bond angles. So complexes with very bulky ligands (triphenylphosphine, $\mathrm{PPh}_{3}$, for example), are usually high-spin and tetrahedral. But the balance between all these factors seems to be quite subtle, since the complex $\left[\mathrm{Ni}\left(\mathrm{PPh}_{2} \mathrm{Et}\right)_{2} \mathrm{Br}_{2}\right]$ has been isolated in both forms, squareplanar and tetrahedral; they are found to be in equilibrium in solution and are thus very close in energy.


2-58 Square planar ( $d^{8}$ low-spin) Tetrahedral ( $d^{8}$ low-spin)

As we have already remarked at the beginning of this section, not all $\mathrm{ML}_{4}$ complexes adopt such high-symmetry structures. In this context, $d^{9}$ complexes, whose electronic configuration is 'intermediate' between one which favours a square-planar geometry (diamagnetic $d^{8}$ complexes) and others which favour a tetrahedral geometry (paramagnetic $d^{8}$ or diamagnetic complexes $d^{10}$ ), are particularly interesting. The two structural types are indeed found in the family of $d^{9} \mathrm{ML}_{4}$ complexes: for example, square-planar for $\left[\mathrm{M}(\mathrm{py})_{4}\right]^{2+}(\mathrm{M}=\mathrm{Cu}, \mathrm{Ag})(\mathrm{py}=$ pyridine $)$ but quasi-tetrahedral for $\left[\mathrm{Co}(\mathrm{CN})_{4}\right]^{4-},\left[\mathrm{Co}\left(\mathrm{PMe}_{3}\right)_{4}\right]$, and $\left[\mathrm{Ni}\left(\mathrm{PR}_{3}\right)_{3} \mathrm{X}\right]$ ( $\mathrm{X}=\mathrm{Cl}, \mathrm{Br}, \mathrm{I}$ ). A given complex can even adopt very different geometries depending on its environment. For example, 131 structures of the $\left[\mathrm{CuCl}_{4}\right]^{2-}$ anion have been published, ${ }^{\top}$ which differ in the nature of the associated cation. The whole range of geometries has in fact been observed, from the square plane to an almost ideal tetrahedron!



2-59
${ }^{8}$ For this geometrical parameter, the difference appears to be so small that it can depend on the environment! For example, in the solid-state structure of $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$ determined by X-ray diffraction, the axial bonds are slightly longer than the equatiorial, by about $0.008 \AA$ (D. Braga, F. Grepioni, and G. Orpen Organometallics 12, 1481 (1993)). But a study of the structure in the gas phase, by electron diffraction, gives the opposite result, with the equatorial bonds being longer by $0.01-0.03 \AA$ (B. W. McClelland, A. G. Robiette, L. Hedberg, and K. Hedberg Inorg. Chem. 40, 1358 (2001)).

### 2.5. Trigonal-bipyramidal $\mathrm{ML}_{5}$ complexes

In an $\mathrm{ML}_{5}$ complex which adopts a trigonal-bipyramidal (TBP) geometry, we distinguish the equatorial ligands ( $\mathrm{L}_{1}, \mathrm{~L}_{2}$, and $\mathrm{L}_{3}$ ) from those in axial positions ( $\mathrm{L}_{4}$ and $\mathrm{L}_{5}, \mathbf{2 - 5 9}$ ). The former are located at the vertices of the triangular base of the bipyramid, and they define the equatorial plane of the complex (the $x y$ plane in $2-59$ ), while the axial ligands are located at the vertices of the bipyramid. The angles between equatorial bonds are $120^{\circ}$ (trigonal base), but those between an equatorial and an axial bond are $90^{\circ}$.

## Comment

The equatorial and axial ligands are not equivalent by symmetry, since no symmetry operation of the complex's point group $\left(D_{3 h}\right)$ interchanges an equatorial ligand with an axial. Nor are they equivalent from a chemical point of view, and as a consequence, the $\mathrm{M}-\mathrm{L}_{e q}$ and $\mathrm{M}-\mathrm{L}_{a x}$ bond lengths can differ, even when all the ligands $L$ are identical. ${ }^{8}$

### 2.5.1. Characterization of the d block

With the orientation of the axes chosen in 2-59, the equatorial ligands are located in the $x y$ plane and the axial ligands on the $z$-axis, that is, at the intersection of the $x z$ and $y z$ planes.

The five ligands are thus placed in one or other of the nodal planes of the $x z$ (nodal planes $x y$ and $y z$ ) and $y z$ (nodal planes $x y$ and $x z$ ) orbitals. As a result, there cannot be any $\sigma$ interaction between the ligands and these two orbitals which stay nonbonding in the d-block of a TBP complex (2-60). By consulting the character table for the $D_{3 h}$ point group, we find that these two degenerate orbitals have $\mathrm{e}^{\prime \prime}$ symmetry.


2-60


2-61

A third orbital in the d block may readily be characterized. The $z^{2}$ orbital is destabilized by antibonding interactions with all five ligands. This is the highest-energy orbital in the d block, with the strongest antibonding interactions involving the axial ligands along the $z$-axis (2-61).
${ }^{9}$ The linear combinations of the ligand orbitals with which they interact also have $e^{\prime}$ symmetry, of course (see Chapter 6, § 6.6.4).

The $x y$ and $x^{2}-y^{2}$ orbitals have not yet been considered. The first is destabilized by antibonding interactions with the ligands $L^{2}$ and $L^{3}$ (2-62). No interaction is possible with the three other ligands which are located in the nodal plane $y z$. For the $x^{2}-y^{2}$ orbital, its two nodal planes bisect the $x$ - and $y$-axes and they contain the $z$-axis. It therefore cannot interact with the axial ligands placed on this axis, but antibonding interactions will develop with the three equatorial ligands, as none of them is located in either of the nodal planes. Two representations can be given for this orbital (2-63). In the first (2-63a), the orientation of the axes is the same as that used previously, whereas in the second $(\mathbf{2 - 6 3 b})$, the plane of the page is chosen as the equatorial plane $(x y)$. This second representation allows us to understand more easily the sign of the coefficients for the orbitals on the ligands $L_{2}$ and $L_{3}$ if the interaction with $x^{2}-y^{2}$ is to be antibonding (the sign for the orbital $\sigma_{1}$, on the ligand $\mathrm{L}_{1}$, is obvious since it points along the $y$-axis). It is clear from $\mathbf{2 - 6 3 b}$ that the overlap between the orbitals $\sigma_{2}$ and $\sigma_{3}$ and the part of the $x^{2}-y^{2}$ orbital concentrated along the horizontal axis $(y)$ is positive, whereas with the part of the orbital concentrated along the vertical axis $(x)$ it is negative. Now since the angle $L_{2}-M-L_{3}$ is $120^{\circ}$, the $M-L_{2}$ and $M-L_{3}$ bonds are closer to the $x$-axis than to the $y$-axis. The representation given in 2-63 therefore does indeed correspond to a negative value of the total overlap between the orbitals on the ligands $L_{2}$ and $L_{3}$ and the $x^{2}-y^{2}$ orbital on the metal centre, and thus to an antibonding interaction.


It should be noted that the antibonding $\mathrm{M}-\mathrm{L}_{i}$ character in this orbital is not equivalent for the three ligands. $\mathrm{L}_{1}$ is situated on the $y$-axis and therefore has a larger axial overlap than do the ligands $L_{2}$ and $L_{3}$, which are placed between the $x$ - and $y$-axes. Moreover, the coefficient for $\sigma_{1}$ is twice as large as those for $\sigma_{2}$ or $\sigma_{3}$ (Chapter $6, \S$ 6.6.4.2). These two factors lead to the $x^{2}-y^{2}$ orbital being antibonding mainly with the ligand $\mathrm{L}_{1}$, but $x y$ is antibonding only with the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{3}$ (2-62). While this analysis enables us to deduce that the $x y$ and $x^{2}-y^{2}$ orbitals are both antibonding in the $d$ block, an important additional piece of information is provided by examination of the character table for the $D_{3 h}$ group (Table 2.3). From the last column, we learn that the $x y$ and $x^{2}-y^{2}$ orbitals are degenerate by symmetry ( $\mathrm{e}^{\prime}$ symmetry). ${ }^{9}$

Table 2.3. Character table for the $\mathrm{D}_{3 h}$ group

| $D_{3 h}$ | $E$ | $2 C_{3}$ | $3 C_{2}$ | $\sigma_{h}$ | $2 S_{3}$ | $3 \sigma_{v}$ |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\mathrm{~A}_{1}^{\prime}$ | 1 | 1 | 1 | 1 | 1 | 1 |  |
| $\mathrm{~A}_{2}^{\prime}$ | 1 | 1 | -1 | 1 | 1 | -1 |  |
| $\mathrm{E}^{\prime}$ | 2 | -1 | 0 | 2 | -1 | 0 | $(x, y)$ |
| $\mathrm{A}_{1}^{\prime \prime}$ | 1 | 1 | 1 | -1 | -1 | -1 |  |
| $\mathrm{~A}_{2}^{\prime \prime}$ | 1 | 1 | -1 | -1 | -1 | 1 | $z$ |
| $\left.\mathrm{E}^{2}-y^{2}, x y\right)$ |  |  |  |  |  |  |  |
|  | 2 | -1 | 0 | -2 | 1 | 0 |  |

The orbitals shown in 2-62 and 2-63 therefore make up a set of two degenerate orbitals, of $e^{\prime}$ symmetry. Their energy is of course higher than that of the $\mathrm{e}^{\prime \prime}$ orbitals ( $x z$ and $y z$ ), since these latter are strictly nonbonding.

The character table also shows that the $p_{x}$ and $p_{y}$ orbitals on the central atom, like $x y$ and $x^{2}-y^{2}$, have $\mathrm{e}^{\prime}$ symmetry (penultimate column). They can therefore lead to the formation of polarized degenerate $e^{\prime}$ orbitals in the d block, as their contribution polarizes the $x y$ and $x^{2}-y^{2}$ orbitals. A more accurate shape of these orbitals can be obtained by adding a $p_{x}$ contribution to $x y(\mathbf{2 - 6 4})$ and $p_{y}$ to $x^{2}-y^{2}(\mathbf{2 - 6 5})$, so that the overlap of the $p$ orbital with the ligand orbitals is positive (see Appendix A, § 2.2). The addition of $p_{x}$ to the $x y$ orbital polarizes this orbital in the direction away from the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{3}$, thereby decreasing the antibonding interactions with these ligands (2-64). For $x^{2}-y^{2}$, the polarization by $p_{y}$ decreases the amplitude of the orbital in the direction away from the ligand $\mathrm{L}_{1}(\mathbf{2 - 6 5})$ with which the antibonding interaction was the strongest. As in the examples that we have already studied (see § 2.3.1.2 and 2.3.2 for SBP ML 5 complexes), the participation of the $p$ orbitals leads to a reduction of the antibonding character of the d-block orbitals.



Figure 2.10. Representation of the d-block orbitals for an $\mathrm{ML}_{5}$ complex with a TBP geometry, the equatorial plane $(x y)$ being either perpendicular to the plane of the page (left-hand side), or in this plane (right-hand side). In this latter case, for greater clarity, the axial ligands are not shown.

${ }^{10}$ For a review of the structures of $\mathrm{ML}_{5}$ complexes, consult: S. Alvarez and M. Llunell J. Chem. Soc. Dalton Trans. 3288 (2000).

The d block of an $\mathrm{ML}_{5}$ complex with a TBP geometry is therefore made up of two degenerate nonbonding orbitals ( $x z$ and $y z$, with $\mathrm{e}^{\prime \prime}$ symmetry), two degenerate orbitals that are fairly weakly antibonding ( $x y$ and $x^{2}-y^{2}$, with $\mathrm{e}^{\prime}$ symmetry) and one very strongly antibonding orbital ( $z^{2}$, with $a_{1}^{\prime}$ symmetry).

These results are illustrated in Figure 2.10 for two different orientations of the bipyramid, the equatorial plane ( $x y$ ) being either perpendicular to the plane of the page (left-hand side) or in this plane (right-hand side).








### 2.5.2. Electronic structure

### 2.5.2.1. Diamagnetic $d^{8}$ complexes

The most common electron count for $\mathrm{ML}_{5}$ complexes with a TBP geometry is $d^{8}$, which corresponds to the double occupation of the degenerate $\mathrm{e}^{\prime \prime}$ and $\mathrm{e}^{\prime}$ orbitals (2-66). These are therefore 18 -electron complexes, since there are 10 additional electrons associated with the five M-L bonds.

The species $\left[\mathrm{M}(\mathrm{CO})_{5}\right](\mathrm{M}=\mathrm{Fe}, \mathrm{Ru}$, and Os$)$ are 'archetypal' TBP $\mathrm{ML}_{5}$ complexes. We may also cite the complexes $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]^{-}$, $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right], \quad\left[\mathrm{Fe}(\mathrm{CO})_{4}(\mathrm{CN})\right]^{-}, \quad\left[\mathrm{Co}\left(\mathrm{C}=\mathrm{N}-\mathrm{CH}_{3}\right)_{5}\right]^{+}$, $\left[\operatorname{Ir}\left(\mathrm{PR}_{3}\right)_{3}\left(\mathrm{CH}_{3}\right)\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$, and $\left[\mathrm{Ni}\left(\mathrm{P}(\mathrm{OEt})_{3}\right)_{5}\right]^{2+}$. It should be noted that the $d^{8}$ electron count has already been found to be favourable for ML5 complexes that adopt an SBP geometry, with the metal above the basal plane (§ 2.3.3.1.). The two structural types have indeed been characterized in the $d^{8}-M L_{5}$ family of complexes. ${ }^{10}$ In the case of the
${ }^{11}$ The transformation $\operatorname{TBP}(1) \rightarrow \mathrm{SBP} \rightarrow$ $\operatorname{TBP}(2)$ is nothing more than the Berry pseudorotation which allows the exchange of axial and equatorial ligands to take place in a TBP complex.
anion $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$, the SBP structure can even coexist with a slightly distorted TBP structure when the associated cation is $\left[\mathrm{Cr}(\mathrm{en})_{3}\right]^{3+}$ (en=ethylenediamine). ${ }^{11}$

### 2.5.2.2. Other cases

The complex $\left[\mathrm{Co}(\text { dpe })_{2} \mathrm{Cl}\right]^{2+} \quad($ dpe $=1,2-\mathrm{bis}($ diphenylphosphino $)$ ethane) with a $d^{7}$ electronic configuration (a 17-electron low-spin complex) is also known in two structural forms: a red isomer with an SBP geometry (the chloride ligand is in the apical position), and a green isomer with a TBP geometry (the chloride ligand is in an equatorial position). TBP complexes with more than 18 electrons are also known. They are weak-field complexes (the $z^{2}$ orbital is occupied) with a $d^{9}$ electronic configuration (19 electrons), such as $\left[\mathrm{CuCl}_{5}\right]^{3-},\left[\mathrm{CuBr}_{5}\right]^{3-}$, or $\left.[\mathrm{Cu} \text { (imidazole) })_{3} \mathrm{Cl}_{2}\right]$, or $d^{10}$ (20 electrons), such as $\left[\mathrm{CdCl}_{5}\right]^{3-}$ or $\left[\mathrm{HgCl}_{5}\right]^{3-}$.

### 2.6. Trigonal-planar $\mathrm{ML}_{3}$ complexes

A trigonal-planar $\mathrm{ML}_{3}$ complex can be formed by removing the two axial ligands from a TBP ML 5 complex (2-67). The d-block orbitals may therefore be readily deduced from those established in the preceding section.



2-67

### 2.6.1. Characterization of the d block

The variations in energy for the five $d$ orbitals are shown in Figure 2.11. The $x z$ and $y z$ orbitals ( $\mathrm{e}^{\prime \prime}$ ) of a TBP ML ${ }_{5}$ complex have zero coefficients on the axial ligands, as do $x y$ and $x^{2}-y^{2}\left(\mathrm{e}^{\prime}\right)$. The removal of these two ligands does therefore not change either the shapes or the energies of these four orbitals (Figure 2.11). But the $z^{2}$ orbital, which was strongly antibonding due to interactions with the axial ligands, is substantially stabilized by their removal, becoming almost nonbonding. Only three weak antibonding interactions with the ligands in the $x y$ plane are left, and the amplitude of the $z^{2}$ orbital in this plane is small. Note that the orbital symmetries are the same in $\mathrm{ML}_{5}$ (TBP) and $\mathrm{ML}_{3}$ trigonal-planar complexes, as both have $D_{3 h}$ symmetry.

Figure 2.11. Derivation of the d-block orbitals for a trigonal-planar $\mathrm{ML}_{3}$ complex from those of an ML5 complex with a TBP geometry.


2-69


### 2.6.2. 16 -electron $d^{10}$ complexes

The d-block for $\mathrm{ML}_{3}$ complexes with a trigonal-planar geometry thus contains two nonbonding degenerate orbitals ( $x z$ and $y z$ ), a very weakly antibonding orbital $\left(z^{2}\right)$, and two weakly antibonding degenerate orbitals ( $x y$ and $x^{2}-y^{2}$ ). These five orbitals are sufficiently low in energy to be doubly occupied. As a result, the large majority of these complexes have a $d^{10}$ electronic configuration (2-68), the simplest no doubt being $\left[\mathrm{PdH}_{3}\right]^{3-}$. We may also cite the tricarbonyl complexes $\left[\mathrm{M}(\mathrm{CO})_{3}\right]$ $(\mathrm{M}-\mathrm{Ni}, \mathrm{Pd}$, or Pt$)$, as well as $\left[\mathrm{Pt}\left(\mathrm{Ph}_{3} \mathrm{P}\right)_{3}\right],\left[\mathrm{Pt}\left(\eta^{2}\right.\right.$-ethylene $\left.)\left(\mathrm{PR}_{3}\right)_{2}\right]$, or [ $\left.\mathrm{Ni}\left(\eta^{2} \text {-ethylene }\right)_{3}\right]$, complexes in which the oxidation state of the metal is zero. There are also complexes of the metals from groups 11 (oxidation state I), such as $\left[\mathrm{Cu}(\mathrm{CN})_{3}\right]^{2-},\left[\mathrm{Cu}(\mathrm{SR})_{3}\right]^{2-}$, or $\left[\mathrm{Au}\left(\mathrm{PPh}_{3}\right)_{2} \mathrm{Cl}\right]$, and 12 (oxidation state II), such as $\left[\mathrm{HgI}_{3}\right]^{-}$and $\left[\mathrm{Hg}(\mathrm{SR})_{3}\right]^{-}$.

These are all 16 -electron complexes (six for the bonds and ten in the d block). The 'lack' of two electrons compared to the 18 -electron rule arises because a nonbonding orbital on the metal remains empty. As in the case of square-planar $\mathrm{ML}_{4}$ complexes, this is the $p$ orbital perpendicular to the molecular plane (2-69), which, although nonbonding, is too high in energy to be occupied.

### 2.7. Linear $\mathrm{ML}_{2}$ complexes

There are many ways in which a linear $\mathrm{ML}_{2}$ complex may be obtained from the complexes already studied. For example, four coplanar ligands can be removed in an octahedral complex (2-70).




Figure 2.12. Derivation of the d-block orbitals for a linear $\mathrm{ML}_{2}$ complex from the d orbitals of an octahedral $\mathrm{ML}_{6}$ complex.


### 2.7.1. Characterization of the d block

Consider the orbitals in the d block of an octahedron (Figure 2.12, lefthand side). Removal of the four ligands $\mathrm{L}_{3}-\mathrm{L}_{6}$ has no effect on the three nonbonding $\mathrm{t}_{2 g}$ orbitals which have zero coefficients on all the ligands. The antibonding $x^{2}-y^{2}$ orbital in the $\mathrm{e}_{g}$ block has antibonding interactions with the four ligands $\mathrm{L}_{3}-\mathrm{L}_{6}$, but zero coefficients on $\mathrm{L}_{1}$ and $\mathrm{L}_{2}$. After the removal of the ligands $\mathrm{L}_{3}-\mathrm{L}_{6}$, this orbital becomes strictly nonbonding (Figure 2.12). There are, therefore, four nonbonding orbitals in the d block of linear $\mathrm{ML}_{2}$ complexes, whose symmetries are $\pi_{g}(x z, y z)$ and $\delta_{g}\left(x^{2}-y^{2}, x y\right)$ in the $D_{\infty h}$ point group. These orbitals are therefore degenerate in pairs by symmetry, but in fact all four have the same energy (accidental degeneracy).

## Comment

Notice that the $\pi$-type orbitals have just a single nodal plane that contains the internuclear axis, whereas the $\delta$ orbitals have two.

We have not yet considered the second orbital in the $\mathrm{e}_{g}$ block of the octahedron, $z^{2}$. It is stabilized by the elimination of the four weak antibonding interactions in the $x y$ plane, but this stabilization should be rather small, as the two principal antibonding interactions, along the $z$-axis, are still present.

In fact, this fifth orbital is not very high in energy since it is polarized by mixing with the $s$ orbital of the metal ( $s$ and $z^{2}$ have the same symmetry, $\sigma_{g}$, in the $D_{\infty h}$ point group). As in the previous examples which involved polarization by a $p$ orbital (see also Appendix A, § 2.A3), the sign of the $s$ orbital is such that its interaction with the ligand orbitals is bonding (2-71). The effect of this mixing is to increase the amplitude of the $z^{2}$ orbital in the $x y$ plane, perpendicular to the internuclear axis, and to decrease it along this axis. The antibonding character of the polarized
$z^{2}$ orbital is therefore substantially reduced by the participation of the $s$ orbital, so that its energy is not very high (Figure 2.12).


### 2.7.2. Electronic structure

In the very large majority of linear (or quasi-linear) $\mathrm{ML}_{2}$ complexes, the five orbitals of the d block are doubly occupied ( $d^{10}$ electronic configuration, 2-72). The most common examples involve the elements of group $11(\mathrm{Cu}, \mathrm{Ag}, \mathrm{Au})$, where the metal is in the oxidation state I. Anions such as $\left[\mathrm{Cu}(\mathrm{Me})_{2}\right]^{-},\left[\mathrm{Cu}(\mathrm{Ph})_{2}\right]^{-},\left[\mathrm{Cu}(\mathrm{Mes})_{2}\right]^{-}(\mathrm{Mes}=$ $\left.2,4,6-\mathrm{Me}_{3} \mathrm{C}_{6} \mathrm{H}_{2}\right),\left[\mathrm{Ag}(\mathrm{CN})_{2}\right]^{-},\left[\mathrm{Ag}\left(\mathrm{C}\left(\mathrm{SiMe}_{3}\right)_{3}\right)_{2}\right]^{-},\left[\mathrm{Au}\left(\mathrm{C}_{6} \mathrm{~F}_{5}\right)_{2}\right]^{-}$, or $\left[\mathrm{AuCl}_{2}\right]^{-}$, cations such as $\left[\mathrm{Cu}\left(\mathrm{NH}_{3}\right)_{2}\right]^{+},\left[\mathrm{Ag}\left(\mathrm{NH}_{3}\right) 2\right]^{+}$, or $\left[\mathrm{Ag}(\mathrm{CO})_{2}\right]^{+}$, and neutral complexes such as $\left[\mathrm{AuMe}\left(\mathrm{PMe}_{3}\right)\right]$ or $[\mathrm{AuCl}(\mathrm{Mes})]$ are all known. There are also examples from group 12, such as $\left[\mathrm{Hg}(\mathrm{CN})_{2}\right]$.

Although the d block is full, these complexes have only 14 electrons. This electronic deficiency is linked to the presence of two $p$ orbitals on the metal (2-73) which remain empty, even though they are nonbonding, since they are too high in energy to be occupied. The small number of ligands, combined with this electron deficiency, makes these complexes very reactive. They can either bind other ligands, or exist as chains, in which the coordination number of the metal is higher than two, rather than as $\mathrm{ML}_{2}$ monomers.

We should mention a few complexes in which the d block is not completely filled. For example, complexes with two borylamide ligands $\left[\mathrm{N}(\mathrm{Mes})\left(\mathrm{BR}_{2}\right)\right]$ are known for $\mathrm{M}=\mathrm{Mn}\left(d^{5}\right), \mathrm{Fe}\left(d^{6}\right), \mathrm{Co}\left(d^{7}\right)$, and $\mathrm{Ni}\left(d^{8}\right)$, and they adopt an almost linear geometry. These are all high-spin complexes, which is understandable given the small separation of the energy levels in the d block.

### 2.8. Other complexes or $\mathrm{ML}_{\boldsymbol{n}}$ fragments

Other structures exist for $\mathrm{ML}_{n}$ complexes besides those considered so far, though they are less common for stable species. We shall study five
of them in this section: pyramidal $\mathrm{ML}_{3}$ complexes (2-74a), 'T-shaped' $\mathrm{ML}_{3}(\mathbf{2}-\mathbf{7 4 b})$, 'butterfly' $\mathrm{ML}_{4}$ (2-74c), bent $\mathrm{ML}_{2}$ (2-74d), and ML (2-74e).


These entities are interesting since they may be considered as 'fragments' of more important complexes, such as octahedral or square-planar. In this context, knowledge of their orbital structure is useful for two reasons:

1. They can be used as fragment orbitals in some complexes. For example, the main MO of the complex $\left[\left(\eta^{5}-\right.\right.$ cyclopentadienyl) $\left.\mathrm{Mn}(\mathrm{CO})_{3}\right](\mathbf{2 - 7 6})$ can be constructed by interaction of the $\pi$ MO of cyclopentadienyl with the $d$ orbitals of the metallic fragment $\left[\mathrm{Mn}(\mathrm{CO})_{3}\right]$ in a pyramidal geometry.
2. The $d$ orbitals of these species 'resemble' the orbitals of certain organic molecules or fragments, both in their shape and their electronic occupation. It is therefore important to be familiar with them if one wishes to establish a link between the electronic structures of transition metal complexes and of organic molecules (the isolobal analogy, see Chapter 5).

### 2.8.1. Pyramidal $\mathrm{ML}_{3}$ complexes

The pyramidal $\mathrm{ML}_{3}$ structure (2-74a), with $C_{3 v}$ symmetry, is not common for complexes with three ligands. But a few examples are known, such as the $d^{0}$ complexes $\left[\mathrm{M}\left(\mathrm{HC}\left(\mathrm{SiMe}_{3}\right)_{2}\right)_{3}\right](\mathrm{M}=\mathrm{Y}$ or La ), or the low-spin $d^{6}$ complexes $\left[\mathrm{M}(\mathrm{Mes})_{3}\right](\mathrm{M}=\mathrm{Rh}$ or Ir$)(\mathbf{2 - 7 5})$. Moreover, the pyramidal $\mathrm{ML}_{3}$ entity can be considered as a fragment in 18 -electron complexes of the type [(polyene)-ML ${ }_{3}$ ], such as $\left[\left(\eta^{4}\right.\right.$ cyclobutadiene $\left.) \mathrm{Fe}(\mathrm{CO})_{3}\right],\left[\left(\eta^{5}-\mathrm{Cp}\right) \mathrm{Mn}(\mathrm{CO})_{3}\right](\mathrm{Cp}=$ cyclopentadienyl) (2-76) or $\left[\left(\eta^{6}\right.\right.$-benzene $\left.) \mathrm{Cr}(\mathrm{CO})_{3}\right]$.


2-76


Figure 2.13. Derivation of the d-block orbitals for a pyramidal $\mathrm{ML}_{3}$ complex from the $d$ orbitals of an octahedral $\mathrm{ML}_{6}$ complex. The hybrid $s-p$ orbital lying above the $d$ block is also shown.

In one approach to the derivation of the d-block orbitals of this complex, we start from an octahedral complex with the geometry indicated in 2-77 and remove the three ligands $\mathrm{L}_{4}, \mathrm{~L}_{5}$, and $\mathrm{L}_{6}$. A pyramidal $\mathrm{ML}_{3}$ fragment is thereby obtained, in which the angles between the $\mathrm{M}-\mathrm{L}$ bonds are $90^{\circ}$. In the orientation that is chosen, the $z$-axis coincides with the $C_{3}$ axis of the resulting $\mathrm{ML}_{3}$ complex.




2-77

The $\mathrm{t}_{2 g}$ (nonbonding) and $\mathrm{e}_{g}$ (antibonding) orbitals of an octahedral complex with this orientation of the axes have already been established ( $\S 2.1 .2 .5$, Figure 2.5) and are presented on the left-hand side of Figure 2.13. The removal of three ligands does not change the three nonbonding orbitals, since the coefficients on these ligands were zero: they remain nonbonding and degenerate in the $\mathrm{ML}_{3}$ fragment.

The $z^{2}$ orbital has $\mathrm{a}_{1}$ symmetry (the totally symmetric representation of the $\mathrm{C}_{3 v}$ point group), while the two other orbitals are degenerate by symmetry (1e). The $e_{g}$ orbitals are stabilized by the elimination of half of the antibonding interactions that were present in the octahedron. However, they are still antibonding in the $\mathrm{ML}_{3}$ fragment (Figure 2.13) and make up the degenerate orbital set 2 e . As in several of the examples already treated, the antibonding character of these orbitals is reduced by mixing with the $p_{x}(\mathbf{2 - 7 8})$ or $p_{y}(\mathbf{2 - 7 9})$ orbitals, which polarizes them in the direction opposite to the ligands.




2-80
${ }^{12}$ There are also a few planar complexes in which two of the three $\mathrm{L}-\mathrm{M}-\mathrm{L}$ angles are larger than $120^{\circ}$ (Y-shaped structures). The preference for T- or Y-shaped structures will be analysed in greater detail in Chapter 4. For further information on the structures of $\mathrm{ML}_{3}$ complexes, the reader may consult: S. Alvarez Coord. Chem. Rev. 193-195, 13 (1999).


2-78


2-79
If the $d$ block were completely filled, we would have a 16 -electron complex. There would therefore be an empty nonbonding orbital on the metal centre. For a trigonal-planar $\mathrm{ML}_{3}$ complex, that would be the $p_{z}$ orbital perpendicular to the molecular plane (2-69). But in the present case, it is a hybrid orbital of $a_{1}$ symmetry, a linear combination of the $s$ and $p$ orbitals on the metal, which points in the direction opposite to the ligands (2-80). Due to the relative energies of the $n d,(n+1) s$, and $(n+1) p$ AO on the metal, it stays higher in energy than the 2e orbitals which are essentially of $d$ character on the metal and weakly antibonding. This orbital lies lower than the pure $p$ orbital of a trigonal-planar ML ${ }_{3}$ complex, thanks to the contribution from the $s$ orbital $\left(\varepsilon_{s}<\varepsilon_{p}\right)$.

Notice that the structure of the d block, with three strictly nonbonding orbitals, enables us to understand why this geometry is observed for low-spin $d^{6}$ complexes, such as that represented in 2-75.

### 2.8.2. 'T-shaped' $\mathrm{ML}_{3}$ complexes

There are a few rare $\mathrm{ML}_{3}$ complexes that adopt a non-trigonal planar geometry, in which one $\mathrm{L}-\mathrm{M}-\mathrm{L}$ angle is substantially larger than $120^{\circ}$ ( $C_{2 v}$ symmetry). ${ }^{12}$ This type of structure is sketched in 2-81, where the angle $\mathrm{L}_{1}-\mathrm{M}-\mathrm{L}_{3}$ has the limiting value of $180^{\circ}$, justifying the name of a T-shaped structure. The diamagnetic $d^{8}$ complexes $\left[\mathrm{Rh}\left(\mathrm{PPh}_{3}\right)_{3}\right]^{+}{ }_{(2-}$ $\mathbf{8 2})$ and $\left[\mathrm{Ni}(\mathrm{Mes})_{3}\right]^{-}$are typical examples. The T-shaped entity $\mathrm{ML}_{3}$ is also interesting as it can be used as a fragment in monometallic complexes such as Zeise's salt, $\left[\operatorname{Pt}(\mathrm{Cl})_{3}\left(\eta^{2} \text {-ethylene }\right)\right]^{-}(\mathbf{2 - 8 3})$, or bimetallic complexes of the type $\left[\mathrm{Pd}_{2} \mathrm{~L}_{6}\right]^{2+}(2-84)$.


2-81


2-82

Figure 2.14. Derivation of the d-block orbitals of a T-shaped $\mathrm{ML}_{3}$ complex from the $d$ orbitals of a square-planar $\mathrm{ML}_{4}$ complex.


2-83



The orbitals of the $d$ block may be derived from those of a square-
( 2.2 .1 ), by removing one of the ligands (2-85).
The orbitals of the d block may be derived from those of a
planar complex (§ 2.2.1), by removing one of the ligands (2-85).





2-85

Changes to the shapes and energies of the orbitals are shown in Figure 2.14.

The three strictly nonbonding orbitals in the square-planar complex (with $\mathrm{b}_{2 g}$ and $\mathrm{e}_{g}$ symmetries in the $D_{4 h}$ point group) are not affected by the removal of a ligand. They become the $a_{2}, b_{1}$, and $b_{2}$ nonbonding orbitals of the $\mathrm{ML}_{3}$ fragment. The $z^{2}$ orbital ( $\mathrm{a}_{1 g}$ ) is very weakly stabilized $\left(1 a_{1}\right)$, due to the disappearance of a small antibonding interaction, while the $x^{2}-y^{2}$ orbital, which is antibonding towards the four ligands in the square-planar complex $\left(\mathrm{b}_{1 g}\right)$, is stabilized and becomes the $2 \mathrm{a}_{1}$ orbital in the $\mathrm{ML}_{3}$ complex. This stabilization is enhanced by mixing with the $p_{y}$ orbital which polarizes the orbital in the opposite direction to the ligand $\mathrm{L}_{2}(\mathbf{2 - 8 6})$. But overall, this orbital stays fairly high in energy, due to the three antibonding interactions that are still present in the $x y$ plane, along the directions of the lobes of $x^{2}-y^{2}$.
 the removal of a ligand. They become the $\mathrm{a}_{2}, \mathrm{~b}_{1}$, and $\mathrm{b}_{2}$ nonbonding



$$
+
$$


$=$


2-86

The T-shaped $\mathrm{ML}_{3}$ fragment is thus characterized by the presence of four low-energy $d$ orbitals, instead of the five present in the trigonalplanar geometry ( $\S 2.6 .1$ ). As a result, the stable complexes mentioned at the beginning of this section, such as $\left[\mathrm{Rh}\left(\mathrm{PPh}_{3}\right)_{3}\right]^{+}$(2-82), are diamagnetic and have a $d^{8}$ electronic configuration, whereas the $d^{10}$ complexes, which are far more numerous, adopt a trigonal-planar geometry ( $\S 2.6 .2$ ).

### 2.8.3. 'Butterfly' $\mathrm{ML}_{4}$ complexes

In the 'butterfly' $\mathrm{ML}_{4}$ complex sketched in 2-74c, one of the $\mathrm{L}-\mathrm{M}-\mathrm{L}$ angles is $180^{\circ}$, where the two bonds form the body of the butterfly. The angle between the other two bonds, which form the (folded) wings of the butterfly, is much smaller. More generally, $\mathrm{ML}_{4}$ fragments are described as 'butterfly' when they have $C_{2 v}$ symmetry, which only implies the existence of two planes of symmetry and a $C_{2}$-axis (2-87). It is therefore possible for the two angles mentioned to have very variable values, so long as they are not equal (the symmetry would then be $D_{2 d}$, as for the intermediate points along the deformation shown in 2-53). This $\mathrm{ML}_{4}$ structure, adopted by some complexes such as $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$ or $\left[\operatorname{Ir}(\mathrm{H})_{2}\left(\mathrm{P}\left({ }^{\mathrm{t}} \mathrm{Bu}\right)_{2} \mathrm{Ph}\right)\right]^{+}(\mathbf{2 - 8 8})$, can also be regarded as a fragment in the analysis of the electronic structure of other species, such as the carbene complex $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\mathrm{CR}_{2}\right)\right](\mathbf{2 - 8 9})$ or the bimetallic complex $\left[\mathrm{Fe}_{2}(\mathrm{CO})_{8}\right]$ (2-90).


2-87


2-88


2-89


2-90
${ }^{13}$ In the character table for the $C_{2 v}$ group (Chapter 6, § 6.2.5.1), one does indeed find that the symmetries of the $x y, x z$, and $y z$ orbitals are $a_{2}, b_{1}$, and $b_{2}$. However, the name of the orbital in a given symmetry is not the same as that we have just indicated. For example, $x y$ is placed in the $\mathrm{a}_{2}$ representation whereas we indicated that it has $b_{2}$ symmetry. These differences arise because the axis system previously used for the octahedron is not the one conventionally adopted for the $C_{2 v}$ group. For example, the $C_{2}$ axis should be the $z$-axis, rather than the $y$-axis. In other words, the orbitals sketched on the right-hand side of Figure 2.15 do indeed correspond to the symmetries indicated, and only their names would be changed if one adopted the axis system appropriate for the $C_{2 v}$ group.

Figure 2.15. Derivation of the d-block orbitals for a 'butterfly' $\mathrm{ML}_{4}$ complex from the $d$ orbitals of an octahedral $\mathrm{ML}_{6}$ complex. The hybrid $s-p$ orbital $\left(2 \mathrm{a}_{1}\right)$ placed between the $\mathrm{b}_{2}$ and $3 \mathrm{a}_{1}$ orbitals of the d block is also indicated.

A butterfly $\mathrm{ML}_{4}$ fragment can be obtained by the removal of two cis ligands in an octahedral complex ( $\mathrm{L}_{5}$ and $\mathrm{L}_{6}$ in 2-91). The $\mathrm{L}_{1}-\mathrm{M}-\mathrm{L}_{2}$ and $\mathrm{L}_{3}-\mathrm{M}-\mathrm{L}_{4}$ angles are then $90^{\circ}$ and $180^{\circ}$, respectively. The $d$ orbitals that we shall establish are adapted to this particular geometry, so some changes may be anticipated if these angles have different values (see, for example, Exercise 2.9).





2-91
We have already established the $d$ orbitals of an octahedral complex in the orientation chosen above, with four ligands placed along the bisectors of the $x$ - and $y$-axes ( $\S 2.1 .2 .5$, Figure 2.4); they are shown in the left-hand side of Figure 2.15.

Removal of the ligands $\mathrm{L}_{5}$ and $\mathrm{L}_{6}$ has no effect on either the shape or the energy of the three nonbonding $t_{2 g}$ orbitals of the octahedron (zero coefficients on these ligands). These orbitals ( $\mathrm{a}_{2}(x z), \mathrm{b}_{1}(y z)$, and $\left.1 \mathrm{a}_{1}\left(x^{2}-y^{2}\right)\right)$ therefore remain degenerate, not, as in the octahedron, by symmetry, but 'accidentally' in the lower-symmetry point group $C_{2 v}$. In the $\mathrm{e}_{g}$ block, the $z^{2}$ orbital is only very slightly stabilized by the removal of the two weak antibonding interactions in the $x y$ plane. It therefore stays as a strongly antibonding orbital ( $3 a_{1}$ ). In contrast, the $x y$ orbital is strongly stabilized by the removal of two of the four antibonding interactions with the ligands located in the $x y$ plane $\left(\mathrm{b}_{2}\right){ }^{13}$ This stabilization is enhanced by mixing with the $p_{x}$ orbital, which polarizes it in


the direction opposite to the ligands $L_{1}$ and $L_{2}$, in a way which we have already seen several times (2-92).


2-92

There are therefore four nonbonding or weakly antibonding $d$ orbitals. If these orbitals are doubly occupied, the result, taking the four bonds into account, will be a 16 -electron complex. As in the examples that we have already met, this result is linked to the existence of a nonbonding orbital on the metal which is not of $d$ type. In the $\mathrm{ML}_{4}$ butterfly complex, this is an $s-p$ hybrid orbital of $a_{1}$ symmetry (2-93) which points along the $C_{2}$ symmetry axis of the complex, that is, between the two ligands which were removed in the original octahedral complex. Its energy, which is not very high due to the participation of the $s$ orbital, is lower than that of the most antibonding orbital in the d block, as shown in Figure 2.15.

### 2.8.4. Bent $\mathrm{ML}_{2}$ complexes

$\mathrm{ML}_{2}$ complexes usually adopt a linear (or essentially linear) geometry (see § 2.7). Knowledge of the orbital structure for strongly bent $\mathrm{ML}_{2}$ species is therefore useful mainly when one wishes to consider them as fragments in larger complexes. For example, the $\left[\left(\eta^{2}-\right.\right.$ ethylene) $\left.\mathrm{Ni}\left(\mathrm{PR}_{3}\right)_{2}\right]$ complex (2-94) can be described in terms of a bent $\mathrm{ML}_{2}$ entity interacting with a molecule of ethylene.

Starting from the butterfly $\mathrm{ML}_{4}$ fragment previously studied, a bent $M L_{2}$ fragment may be obtained by removing the ligands $L_{3}$ and $L_{4}(2-95)$.


The $d$ orbitals of this fragment are readily obtained from those of the $\mathrm{ML}_{4}$ fragment (Figure 2.16). The shape and energy of the three nonbonding orbitals ( $a_{2}, \mathrm{~b}_{1}$, and $1 \mathrm{a}_{1}$ ), and of the antibonding $\mathrm{b}_{2}$ orbital, are all unchanged (zero coefficients on $L_{3}$ and $L_{4}$ ). However, the two main antibonding interactions in the $z^{2}$ orbital are eliminated, so that

Figure 2.16. Derivation of the d-block orbitals for a bent $\mathrm{ML}_{2}$ complex from the $d$ orbitals of a 'butterfly' $\mathrm{ML}_{4}$ complex. The hybrid $s-p$ orbital ( $2 \mathrm{a}_{1}$ ) and the metal $p$ orbital (lying above the $d$ block) are also indicated.

$\mathrm{a}_{1}$
$\mathrm{b}_{1}$

2-96


2-97


2-98

this orbital becomes nearly nonbonding $\left(2 a_{1}\right)$. The five orbitals in the $d$ block are therefore nonbonding or weakly antibonding. Once the four electrons associated with the two M-L bonds are taken into account, a $d^{10}$ electronic configuration leads to a 14 -electron complex.

This lack of four electrons compared to the 18 -electron rule arises because there are two nonbonding orbitals on the metal, which are not of $d$-type (2-96). The first is an $s-p$ hybrid orbital $\left(3 \mathrm{a}_{1}\right)$, identical to that found for the 'butterfly' $\mathrm{ML}_{4}$ fragment (Figure 2.16, $2 \mathrm{a}_{1} \rightarrow 3 \mathrm{a}_{1}$ ). The second is the $p$ orbital perpendicular to the plane of the fragment $\left(2 b_{1}\right)$, which is strictly nonbonding. It is higher in energy than the $3 \mathrm{a}_{1}$ orbital, mainly because the $s$ contribution in this latter hybrid orbital lowers its energy $\left(\varepsilon_{s}<\varepsilon_{p}\right)$. Notice that these two orbitals have the same shape as the two nonbonding MO in bent $\mathrm{AH}_{2}$ molecules, where A is a maingroup element (Chapter 1, Figure 1.5), and that their energetic order is the same $\left(\varepsilon\left(\mathrm{a}_{1}\right)<\varepsilon\left(\mathrm{b}_{1}\right)\right)$.

### 2.8.5. ML complexes

A coordination number of 1 is rare for stable complexes. The examples that are known involve a very bulky ligand such as the aryl radical, 2, 4, 6- $\mathrm{Ph}_{3} \mathrm{C}_{6} \mathrm{H}_{2}$, which forms $d^{10} \mathrm{ML}$ complexes with copper and silver (2-97). The M-L unit can also be considered as a fragment in a complex with a higher coordination number, such as $\mathrm{W}-\mathrm{CO}$ in the complex $\left[\left(\eta^{2} \text {-acetylene }\right)_{3} \mathrm{~W}(\mathrm{CO})\right](\mathbf{2 - 9 8})$.

The orbitals of this fragment can be obtained in different ways, for example, by removing a ligand from a linear $\mathrm{ML}_{2}$ complex (2-99) whose orbitals have already been derived ( $\oint$ 2.7.1).

The four strictly nonbonding $d$ orbitals of the $\mathrm{ML}_{2}$ complex ( $\pi_{g}$ and $\delta_{g}$ ) remain unchanged. The $z^{2}$ orbital is stabilized by the elimination of

Figure 2.17. Derivation of the d-block orbitals for an ML complex from the $d$ orbitals of a linear $\mathrm{ML}_{2}$ complex. The nonbonding orbitals ( $s-p$ hybrid and pure $p$ ), located above the d block, are also shown ( $2 \sigma$ and $2 \pi$ ).


2-100


2-99

an antibonding interaction (Figure 2.17). This stabilization is increased by mixing with the $p_{z}$ orbital which polarizes $z^{2}$ in the direction opposite to the ligand (as happens for the $z^{2}$ orbital in an $\mathrm{ML}_{5}$ complex (SBP), $\mathbf{2 - 4 4}$ ). Notice that the removal of the inversion centre ( $\mathrm{D}_{\infty h} \rightarrow \mathrm{C}_{\infty v}$ ) leads to the disappearance of the $g$ and $u$ subscripts in the symmetry labels (Chapter 6, § 6.2.5.2).

The $d$ block of the M-L fragment therefore contains five low-energy orbitals that are all likely to be occupied. The rare stable complexes do indeed have a $d^{10}$ electronic configuration (2-97). They are therefore 12-electron complexes with three empty nonbonding orbitals on the metal, an $s-p$ hybrid orbital ( $\sigma$ symmetry) and two pure $p$ orbitals of $\pi$ symmetry that are higher in energy (2-100).

## Exercises

## 2.1

1. Give the $d^{n}$ electronic configuration of $\left[\mathrm{M}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ octahedral complexes from the first transition series $(\mathrm{M}=\mathrm{Sc}, \mathrm{Ti}, \mathrm{V}, \mathrm{Cr}, \mathrm{Mn}$, $\mathrm{Fe}, \mathrm{Co}, \mathrm{Ni}, \mathrm{Cu}, \mathrm{Zn}$ ).
2. If these are high-spin complexes, give their ground-state electronic configuration (just for the d block).
3. How many unpaired electrons are there in each case?

## 2.2

1. Give the $d^{n}$ electronic configuration of the complex $\left[\mathrm{Cr}(\mathrm{CO})_{6}\right]$.
2. Is this a high-spin or a low-spin complex? Deduce its ground-state electronic configuration (just for the d block).
3. Reduction of this complex leads to dissociation into $\left[\mathrm{Cr}(\mathrm{CO})_{5}\right]^{-}$ and CO , rather than to the formation of $\left[\mathrm{Cr}(\mathrm{CO})_{6}\right]^{-}$. Suggest an explanation.

## 2.3

Making use of the results given in Figure 2.3, give the approximate shapes and relative energies of the bonding MO for an octahedral complex whose ligands are located as shown in 2-29.

## 2.4

Give the shapes and relative energies of the d-block MO for a squareplanar $\mathrm{ML}_{4}$ complex whose ligands are located on the bisectors of the $x$ - and $y$-axes.

## 2.5

1. How many electrons are there around the metal in the $\left[\mathrm{Rh}\left(\mathrm{PPh}_{3}\right)_{3}\right]^{+}$complex whose geometry is given in 2-82?
2. Link this result to the existence of nonbonding MO, lying above the d block, in a 'T-shaped' ML 3 complex (see Figure 2.14, right-hand side, for this d block).
3. Give the shapes of these MO.

## 2.6

1. Construct the $d$-orbital correlation diagram linking a squareplanar $\left(\theta=90^{\circ}\right.$ ) to a pyramidal $\mathrm{ML}_{4}$ complex ( $\theta=125.5^{\circ}$, nodal cone angle for the $z^{2}$ orbital).

2. What happens to the shape and the energy of the nonbonding metal orbital that is not a $d$ orbital?

## Apical and basal bonds in $\mathrm{ML}_{5}$ complexes (SBP)



## 2.7

1. Give the shapes and relative energies of the five d-block orbitals for an $\mathrm{ML}_{5}$ complex with an SBP geometry, if the metal is in the basal plane.

In what follows, we shall only consider this geometry, even if the metal in some cases is above the basal plane.
2. In these complexes, we distinguish the basal ligands $\left(\mathrm{L}_{\mathrm{b}}\right)$ from the apical one $\left(L_{a}\right)$. Starting from the knowledge that the $M-L_{a}$ bond is either a little shorter than or roughly equal in length to the $\mathrm{M}-\mathrm{L}_{\mathrm{b}}$ bonds for low-spin $d^{6}$ complexes, rationalize the following experimental results for the different electronic configurations:
(a) In the low-spin complex $\left[\mathrm{Co}(\mathrm{CN})_{5}\right]^{3-}$, on the other hand, we find $\mathrm{Co}-(\mathrm{CN})_{\mathrm{a}}=2.01 \AA$ but $\mathrm{Co}-(\mathrm{CN})_{\mathrm{b}}=1.90 \AA$.
(b) This trend is reinforced in the low-spin complex $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$, where $\mathrm{Ni}-(\mathrm{CN})_{\mathrm{a}}=2.17 \AA$ but $\mathrm{Ni}-(\mathrm{CN})_{\mathrm{b}}=1.85 \AA$.
3. Which is the longer bond $\left(\mathrm{Mn}-\mathrm{Cl}_{\mathrm{a}}\right.$ or $\left.\mathrm{Mn}-\mathrm{Cl}_{\mathrm{b}}\right)$ in the high-spin complex $\left[\mathrm{MnCl}_{5}\right]^{2-}$ ?
4. What differences may be anticipated between a low-spin and a high-spin $d^{8}$ complex?

## Bond lengths in $d^{8}-\mathrm{ML}_{4}$ complexes

## 2.8

1. Indicate the splitting of the d block in an $\mathrm{ML}_{4}$ complex whose geometry is
(i) square-planar;
(ii) tetrahedral.
2. What is the most stable electronic configuration for a $d^{8}$ complex in each geometry?
3. It is observed experimentally that typical metal-ligand bond lengths for $d^{8}$ complexes are larger in tetrahedral complexes (see table below: units are $\AA$ ).

|  | Square-planar | Tetrahedral |
| :--- | :--- | :--- |
| $\mathrm{Ni}-\mathrm{P}$ | 2.14 | 2.28 |
| $\mathrm{Ni}-\mathrm{S}$ | 2.15 | 2.28 |
| $\mathrm{Ni}-\mathrm{Br}$ | 2.30 | 2.36 |

Suggest a rationalization.

1. Derive the d-block orbitals for a butterfly $\mathrm{ML}_{4}$ complex from those of a TBP $\mathrm{ML}_{5}$ complex.



2. Are these different from the orbitals derived in $\S 2.8 .3$ starting from an octahedral complex (consider only the four lowest-energy $d$ orbitals)?

### 2.10

Derive the $d$ orbitals for the bent $\mathrm{ML}_{2}$ fragment from those of a square-planar $\mathrm{ML}_{4}$ complex. Compare them with those obtained in § 2.8.4.

## Use of group theory

### 2.11

1. Construct all the MO of a tetrahedral complex, making use of symmetry-adapted combinations of ligand orbitals and the data in the character table for the $\mathrm{T}_{d}$ group (Chapter $6, \S 6.6 .2$ ), following the procedure adopted in Figure 2.2 for an octahedral complex.
2. Indicate the orbitals associated with the $\sigma$ bonds, the d-block orbitals, and the $\sigma^{*}$ antibonding orbitals.

### 2.12

Repeat the questions in 2.11 for a trigonal-planar $\mathrm{ML}_{3}$ complex (use the data in Chapter 6, $\S$ 6.6.3).

## Orbital polarization (Appendix A)

### 2.13

1. In a square-planar $\mathrm{ML}_{4}$ complex, which metal orbital of $s$ or $p$ type can polarize the $z^{2}$ orbital in the d block?
2. What are the consequences of this polarization for the shape of the orbital and for the metal-ligand interactions?

2.A1

$2 . A 2$
${ }^{14}$ Y. Jean and F. Volatron 'An Introduction to Molecular Orbitals', Oxford University Press, NY (1993), chapter 6.


Figure 2.A1. Three-orbital interaction diagram, with two on the metal ( $d$ and $p$ ) and one on the ligands $(l)$.

## Appendix A: polarization of the $d$ orbitals

In many complexes, some of the antibonding orbitals in the d block are polarized by mixing, either with a $p$ orbital or with the $s$ orbital of the metal (e.g. see $\mathbf{2 - 6 4}$ and $\mathbf{2 - 6 5}$ for $d-p$ mixing and 2-71 for $d-s$ mixing). The purpose of this appendix is to study the orbital interaction scheme that produces these polarizations in greater detail.

When an orbital in the $d$ block results simply from the antibonding interaction between a metal $d$ orbital and an orbital of the same symmetry located on the ligands $(\ell)$ (two-orbital interaction scheme, 2.A1), the $d$ orbital is not polarized. However, polarization does occur if there is a $p$ (or $s$ ) orbital on the metal that has the same symmetry as the $d$ and $l$ orbitals, which leads to a three-orbital interaction scheme (2.A2). In this case, while the $d$ and $p$ (or $s$ ) orbitals are of course orthogonal ( $S=0$ ), since they are located on the same atom, the overlaps between the $l$ orbital on the ligands and each of the metal orbitals are non-zero ( $\mathrm{S}_{l-d}$ and $S_{l-p}, 2 . A 2$ ).

## 2.A1. Three-orbital interaction diagram

We consider the case where two orbitals on the metal, for example, $d$ and $p$, have the same symmetry as an orbital $l$ on the ligands. In the interaction diagram (Figure 2.A1), the latter is placed at lower energy than the metal orbitals, due to the greater electronegativity of the ligands. The interaction between these three atomic orbitals leads to the formation of three $\mathrm{MO}\left(\phi_{1}, \phi_{2}\right.$, and $\left.\phi_{3}\right)$. When we consider their energies, $\phi_{1}$ is lower than $l$ and $\phi_{3}$ higher than $p$. The orbital $\phi_{2}$ ends up at an intermediate energy level, a little higher than the $d$ orbital with this combination of initial orbital energies. ${ }^{14}$

Each MO is a linear combination of the three orbitals $l, d$, and $p$. When we examine their shapes, we find that the interactions between $l$ and $d$, and between $l$ and $p$, are both bonding in the lowest-energy MO $\phi_{1}$ and both antibonding in the highest-energy $\mathrm{MO} \phi_{3}$. In the orbital $\phi_{2}$, there is an antibonding overlap $(l-d)$ but also a bonding overlap $(l-p)$ (2.A3).

$\phi_{1}$

$\phi_{2}$

$\phi_{3}$
2.A3
${ }^{15}$ The analysis that follows is also applicable to $\mathrm{ML}_{3}$ complexes with a trigonal-planar geometry.

Figure 2.A2. Fragment orbitals (metal or ligands) with e' symmetry in a TBP $\mathrm{ML}_{5}$ complex. The equatorial ligands are in the plane of the page ( $x y$ ), and the axial ligands are omitted for greater clarity.

## 2.A2. Polarization by a $p$ orbital: TBP ML5 complexes

The d block of an $\mathrm{ML}_{5}$ complex with a TBP geometry contains two strictly nonbonding orbitals with $\mathrm{e}^{\prime \prime}$ symmetry ( $x z, y z$ ), one strongly antibonding orbital with $\mathrm{a}_{1}$ symmetry $\left(z^{2}\right)$, and two weakly antibonding orbitals with $\mathrm{e}^{\prime}$ symmetry $\left(x y, x^{2}-y^{2}\right)($ see $\S 2.5 .1$, Figure 2.10$)$. We shall examine these last two in greater detail.

In the $\mathrm{D}_{3 h}$ point group appropriate for this complex, the $p_{x}$ and $p_{y}$ orbitals have the same symmetry ( $\mathrm{e}^{\prime}$ ) as $x y$ and $x^{2}-y^{2}$ (see Table 2.3, $\S$ 2.5.1). Two combinations of orbitals on the equatorial ligands, $l_{x}$ and $l_{y}$, also have this symmetry (Chapter 6, § 6.6.4.2). These six orbitals are shown in Figure 2.A2, where the plane of the page is the equatorial plane ( $x y$ ). ${ }^{15}$ Notice that it is possible to separate each pair of degenerate orbitals into an $\mathrm{e}_{x}^{\prime}$ component, which is antisymmetric with respect to the $y z$ plane, and an $\mathrm{e}_{y}^{\prime}$ component which is symmetric with respect to this plane. The first group includes the $x y, p_{x}$, and $l_{x}$ orbitals, while the $x^{2}-y^{2}, p_{y}$, and $l_{y}$ orbitals belong to the second. No interaction is possible between orbitals that belong to different groups, since they have different symmetry properties with respect to the $y z$ plane.

$x y\left(\mathrm{e}_{x}^{\prime}\right)$

$p_{x}\left(\mathrm{e}_{x}^{\prime}\right)$

$p_{y}\left(\mathrm{e}_{y}^{\prime}\right)$


$l_{y}\left(\mathrm{e}_{y}^{\prime}\right)$

The interactions between orbitals with $\mathrm{e}^{\prime}$ symmetry therefore give two equivalent blocks, one involving the three $\mathrm{e}_{x}^{\prime}$ orbitals, the other the three $\mathrm{e}_{y}^{\prime}$ orbitals (Figure 2.A3). The MO $1 \mathrm{e}_{x}^{\prime}, 2 \mathrm{e}_{x}^{\prime}$, and $3 \mathrm{e}_{x}^{\prime}$, are in the first, and $1 \mathrm{e}_{y}^{\prime}, 2 \mathrm{e}_{y}^{\prime}$, and $3 \mathrm{e}_{y}^{\prime}$ are in the second.

Application of the rules presented in § 2.A1 leads to the following orbital combinations in the $1 \mathrm{e}_{x}^{\prime}$ (2.A4), $2 \mathrm{e}_{x}^{\prime}$ (2.A5), and $3 \mathrm{e}_{x}^{\prime}$ (2.A6) MO.

The $1 \mathrm{e}_{x}^{\prime}$ orbital (2.A4) belongs to the group of bonding MO in the ML ${ }_{5}$ complex, characterized by (i) bonding metal-ligand interactions, reinforced here by the polarization of the $d$ orbital; (ii) a lower energy than that of the ligand orbitals; and (iii) larger coefficients on the ligands than on the metal. This orbital does not belong to the d block of the complex. The $3 \mathrm{e}_{x}^{\prime}$ orbital (2.A6) is one of the antibonding MO in

Figure 2.A3. Interaction diagram for orbitals with $\mathrm{e}^{\prime}$ symmetry in a TBP ML 5 complex.

the complex, and the $d-p$ mixing in the metal reinforces its antibonding character. This is the highest-energy MO (Figure 2.A3); it is more concentrated on the $p_{x}$ metal orbital than on $x y$ and is also not part of the d block. Among the three MO that we are considering, the one that belongs to the d block is therefore the one with intermediate energy, $2 \mathrm{e}_{x}^{\prime}$ (2.A5), since it is mainly concentrated on the initial orbital with intermediate energy, $x y$. The contribution of the $p_{x}$ orbital to this MO leads to its overlap with the ligand $l_{x}$ orbital being bonding (a result that we had always assumed in Chapter 2). As a result, the $x y$ orbital is polarized in the direction opposite to the ligands, and its antibonding character is reduced. A similar analysis of the interactions of the $\mathrm{e}_{y}^{\prime}$ orbitals shows that the second d-block orbital with $\mathrm{e}^{\prime}$ symmetry is the $2 \mathrm{e}_{y}^{\prime}$ orbital, which is degenerate with $2 \mathrm{e}_{x}^{\prime}$, and in which the $l_{y}$ ligand orbital mixes in an antibonding way with $x^{2}-y^{2}$, but in a bonding way with $p_{y}$ (2.A7).

2.A4

2.A5

$2 . A 6$

2.A7

## 2.A3. Polarization by the $s$ orbital: linear $\mathrm{ML}_{2}$ complexes

It is possible for a $d$ orbital to have the same symmetry as the metal $s$ orbital. In particular, this happens for linear $\mathrm{ML}_{2}$ complexes, in which both the $z^{2}$ and the (higher-energy) $s$ orbitals have $\sigma_{g}$ symmetry. A threeorbital interaction scheme is therefore appropriate for these two orbitals and the ligand orbital $l$ with the same symmetry (Figure 2.A4), giving three $\mathrm{MO}, 1 \sigma_{g}, 2 \sigma_{g}$, and $3 \sigma_{g}$.

As in the preceding example, mixing between $z^{2}$ and $s$ reinforces the bonding character of the $1 \sigma_{g} \mathrm{MO}$ and the antibonding character of the $3 \sigma_{g} \mathrm{MO}$. The orbital with intermediate energy, $2 \sigma_{g}$, is the only one of these three that belongs to the d block. In this orbital, there is antibonding mixing between $z^{2}$ and $l$, but mixing of bonding type between $s$ and $l$ (2.A8). The $s$ contribution decreases the amplitude of $z^{2}$ along the $z$-axis but increases it in the $x y$ plane. The polarization of the $z^{2}$ orbital therefore reduces the antibonding metal-ligand interactions.

Figure 2.A4. Fragment orbitals (metal and ligands) with $\sigma_{g}$ symmetry in a linear $\mathrm{ML}_{2}$ complex.


$z^{2}$


## 2.A8

Without this polarization, the energy of a $z^{2}$ orbital that is antibonding with the ligands located on the $z$-axis is very high. Polarization of the orbital in the $x y$ plane, which does not contain any ligands, lowers its energy sufficiently for it to be occupied in stable linear $\mathrm{ML}_{2}$ complexes ( $d^{10}$ complexes, § 2.7.2).

## 2.A4. Polarization by the $s$ orbital and by a $p$ orbital: SBP $\mathrm{ML}_{5}$ complexes

In some cases, it is possible for a $d$ orbital to be polarized by the $s$ orbital and simultaneously by a $p$ orbital. This happens, for example, to the $z^{2}$ orbital in $\mathrm{ML}_{5}$ complexes with an SBP geometry ( $\S$ 2.3.1.2). This orbital does indeed have the same symmetry as the $s$ and $p_{z}$ orbitals ( $\mathrm{a}_{1}$ in the $\mathrm{C}_{4 v}$ point group). The $z^{2}$ orbital of the d block has antibonding interactions with the ligands ( $z^{2}-l$, 2.A9); the $s$ and $p$ contributions mix in so that:
(1) the contribution from $p_{z}$ is bonding with the apical ligand (located on the $z$-axis);.
(2) the contribution from $s$ is bonding with the basal ligands (located in a plane perpendicular to the $z$-axis) (2.A9).


Notice that we had simplified the problem in § 2.3.1.2, by only considering mixing with the $p_{z}$ orbital. It is indeed this latter which is
responsible for the polarization towards the vacant site of the octahedron that is the most important feature of the $z^{2}$ orbital in SBP ML ${ }_{5}$ complexes.

## Appendix B: orbital energies

When deriving the d block associated with different ligand fields, we obtained nonbonding, 'weakly' antibonding, or 'strongly' antibonding orbitals. These ideas can be made more precise by undertaking calculations which yield numerical values for orbital energies. We have used the extended Hückel method for model complexes of the $\mathrm{MH}_{n}$ type, the $\mathrm{MH}_{n}$ distance being fixed at $1.7 \AA$. The values obtained are given in eV ; while they should only be considered as indications, they can help to give an idea of the energy-level splittings with respect to the initial nonbonding level $\left(\varepsilon_{d}\right)$ in the main ligand fields. Information on the shapes of the orbitals may be obtained from the figures indicated in Chapter 2.


Octahedral complex $\mathrm{CrH}_{6}$ $\varepsilon_{d}(\mathrm{Cr})=-11.22 \mathrm{eV}$
(see Figure 2.3)

SBP complex $\mathrm{CrH}_{5}$
(metal in the basal plane)
$\varepsilon_{d}(\mathrm{Cr})=-11.22 \mathrm{eV}$
(see Figure 2.7)

$$
=\mathrm{b}_{1 g} \quad-8.60
$$



Square-planar complex $\mathrm{NiH}_{4}$ $\varepsilon_{d}(\mathrm{Ni})=-12.99 \mathrm{eV}$
(see Figure 2.6)

Tetrahedral complex $\mathrm{NiH}_{4}$ $\varepsilon_{d}(\mathrm{Ni})=-12.99 \mathrm{eV}$
(see Figure 2.9)


SBP complex $\mathrm{FeH}_{5}$
(metal out of the basal plane) $\varepsilon_{d}(\mathrm{Fe})=-12.70 \mathrm{eV}$
(see Figure 2.8)


TBP complex $\mathrm{FeH}_{5}$ $\varepsilon_{d}(\mathrm{Fe})=-12.70 \mathrm{eV}$ (see Figure 2.10)

'Butterfly' complex $\mathrm{FeH}_{4}$
$\varepsilon_{d}(\mathrm{Fe})=-12.70 \mathrm{eV}$
(see Figure 2.15)

Trigonal-planar complex $\mathrm{CuH}_{3}$ $\varepsilon_{d}(\mathrm{Cu})=-14.00 \mathrm{eV}$
(see Figure 2.11)

'T-shaped' complex $\mathrm{FeH}_{3}$ $\varepsilon_{d}(\mathrm{Fe})=-12.70 \mathrm{eV}$
(see Figure 2.14)

Linear complex $\mathrm{CuH}_{2}$
$\varepsilon_{d}(\mathrm{Cu})=-14.00 \mathrm{eV}$
(see Figure 2.12)

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## $\pi$-type interactions

In the preceding chapter, we derived the structure of the d block for several different $\mathrm{ML}_{n}$ complexes in which the ligands were assumed to use only a single orbital to form the metal-ligand bond. This orbital was either nonbonding or essentially so, and oriented towards the metal centre. Since the overlap between this ligand orbital and a $d$ orbital on the metal is of the axial type, their interaction leads to the formation of a molecular orbital (MO) which characterizes a $\sigma$ bond. As we have already seen in Chapter 1 ( $§$ 1.5.2), some ligands have another orbital (or even two) that can in principle contribute to the bonding with the metal. The orientation of this orbital is usually perpendicular to the axis defined by the $\sigma$ bond, so the resulting overlap with a $d$ orbital on the metal is lateral, and the interaction of $\pi$ type. Two examples are shown below (3-1), in which the ligand orbitals are pure $p$ atomic orbitals.

Among the ligands that can in principle have a $\pi$ interaction with a metal centre, it is important to distinguish two types, depending on the electronic occupation of the orbital that is concerned on the ligand:

- if it is doubly occupied, the ligand is said to be a $\pi$-donor
- if it is empty, then the ligand is a $\pi$-acceptor.

This nomenclature is of course linked to the capacity of the ligand to give or receive electrons through the $\pi$ interaction with the metal.

## Comment

The notation ' $\pi$ ' that is used here is almost never strictly correct according to group theory, where it reserved for doubly degenerate orbitals in linear molecules (Chapter 6). It is, however, widely used to refer to local symmetry; the expression ' $\pi$ interaction' is used when the two orbitals share a common nodal plane and have lateral overlap (3-1).

The examples treated in this chapter involve octahedral complexes. However, the procedure we have used is general and can be applied to all ligand fields. Several illustrations may be found in the exercises at the end of the chapter.

## 3.1. $\pi$-donor ligands: general properties

### 3.1.1. The nature of the $\pi$ orbital on the ligand

A $\pi$-type orbital on a ligand can be doubly occupied only if it is sufficiently low in energy. We may therefore be concerned with:

1. A nonbonding $p$ orbital on a very electronegative atom. Although nonbonding, this orbital is very low in energy and characterizes a lone pair on the atom linked to the metal. For example, consider the $p$ lone pair on the nitrogen atom of an amido group that can interact with one of the metal $d$ orbitals (Scheme 3-2 where the geometry around nitrogen is assumed to be planar).

lone pair



3-2
2. A bonding orbital that characterizes a $\pi$ bond between two atoms of the ligand, as in the case of the imino ligand (3-3) with the doubly occupied $\pi_{\mathrm{NC}}$ orbital. Notice that a $\pi$-bonding orbital on the ligand is automatically accompanied by the corresponding $\pi^{*}$-antibonding orbital. A ligand of this type can therefore act simultaneously as a $\pi$-donor, with its bonding orbital, and as a $\pi$-acceptor, thanks to its empty $\pi^{*}$ orbital. It is categorized as a $\pi$-donor if the former interaction is dominant. A heteronuclear bond provides a favourable case, if, as in the imino ligand (3-3), the more electronegative atom is bonded to the metal. The $\pi$ orbital in such cases is mainly concentrated on the atom bonded to the metal (large overlap), whereas the $\pi^{*}$ orbital is mainly concentrated on the more distant centre (small overlap, see Chapter 1, $\S 1.3 .2$ ).

${ }^{1}$ The polarization of this orbital, in the direction opposite to the metal, significantly weakens its $\sigma$-type overlap with the orbitals of the same symmetry on the metal. This orbital therefore has a negligible influence on the shapes and energies of the MO of the complex, as it remains localized on the ligand.


### 3.1.2. 'Single-face' and 'double-face' $\pi$-donors

Some ligands possess a single nonbonding doubly occupied $p$ orbital, such as the amino ligand (3-2), or a single $\pi$-bonding orbital near to the metal centre, such as the imino ligand (3-3). These are said to be 'single-face' $\pi$-donors.

## $p$ lone pairs


$\sigma$ lone pair


3-4

Other ligands, however, possess two orbitals of this type, in perpendicular planes: these are 'double-face' $\pi$-donors. The halogens are the most common example ( $\mathrm{F}, \mathrm{Cl}, \mathrm{Br}, \mathrm{I}$ ). These are X-type ligands, with seven valence electrons. Formation of the $\sigma$ bond involves the unpaired electron on the halogen and a metal electron; as a result, the oxidation state of the metal is increased by one (Chapter 1, § 1.1.1.2). There are therefore three lone pairs which remain on the halogen bonded to the metal (Chapter 1, § 1.5.2.3): a $\sigma$-type lone pair for which the $\mathrm{M}-\mathrm{X}$ internuclear axis is a symmetry axis, and two lone pairs characterized by nonbonding $p$ orbitals which 'point' perpendicular to the internuclear axis (more strictly, their axes of revolution are perpendicular to this axis, 3-4). Due to its symmetry, the $\sigma$ lone pair cannot participate in a $\pi$-type interaction. ${ }^{1}$ However, the $p$ lone pairs interact with the metal $d$ orbitals through two $\pi$-type interactions which take place in perpendicular planes (3-5). Thus, a halogen linked to a metal centre is a 'double-face' $\pi$-donor.

Analogous interactions occur when the $p$ orbitals are replaced by two $\pi$-bonding orbitals. This is the case when a ligand with a triple bond (one $\sigma$ and two $\pi$ ) binds to the metal in the $\eta^{1}$ mode, as shown in 3-6 for the nitrile ligand. As we have already noted for 'single-face' $\pi$ donors, the $\pi$-donor character of the ligand increases as the orbital coefficient becomes larger on the atom bound to the metal. This therefore

Figure 3.1. Sketch of the interaction between the doubly occupied orbital on a $\pi$-donor ligand $(X)$ and an empty $d$ orbital on the metal centre.
requires the more electronegative atom to be bound to the metal in the $\eta^{1}$ mode.


### 3.1.3. Perturbation of the $d$ orbitals: the general interaction diagram

We now wish to discover how the orbitals of the d block that we have previously obtained, by analysing only the $\sigma$ interactions (Chapter 2), are modified by the presence of a $\pi$ donor. These MO, either nonbonding or antibonding, are mainly concentrated on the $d$ orbitals of the metal, as this is a rather electropositive element. The orbital involved on a $\pi$-donor ligand characterizes either a lone pair located on a very electronegative centre, or a $\pi$ bond. In all cases, the $p$ or $\pi$ orbitals of a $\pi$-donor ligand are lower in energy than the $d$-block orbitals obtained by considering only the $\sigma$ interactions.

This relative position of the orbital energies has important consequences for the way in which the d block is perturbed by interaction with a $\pi$ donor. We shall consider the simplest case, where a nonbonding doubly occupied ligand $p$ orbital interacts with a nonbonding metal $d$ orbital that is assumed to be empty (Figure 3.1). Due to the relative energies of the interacting orbitals, their bonding combination is mainly concentrated on the ligand, whereas the largest coefficient in the antibonding combination is on the metal. Which of these two orbitals is the one that we shall consider a member of the $d$ block? To make our choice, we use the same criterion as that previously adopted when examining $\sigma$-type interactions (see, for example, Chapter 2,

§ 2.1.2.4): it is the molecular orbital that is mainly concentrated on the metal $d$ orbital, that is, the antibonding combination of the interacting orbitals. The bonding molecular orbital, as we have seen, is mainly concentrated on the ligand.

Rule: the $\pi$ interaction between a d-block orbital and the orbital of a $\pi$ donor leads to a destabilization of the d-block orbital, by mixing with the ligand orbital in an antibonding sense.

Figure 3.1 also enables us to obtain further insight into the concept of a $\pi$-donor ligand. Before the interaction, the $d$ orbital on the metal is empty but the $p$ orbital localized on the $\pi$ donor contains two electrons. After the interaction, these electrons occupy the bonding molecular orbital which is partially delocalized onto the metal. This delocalization of the occupied MO on to the two centres results in a partial electron transfer from the ligand to the metal centre, a result which is quite consistent with the label of $\pi$ donor attached to the ligand. The interaction is stabilizing, since it involves two electrons (Chapter 1 § 1.3.2).

However, the $d$ orbital on the metal centre is not necessarily always empty. In fact, the occupation of this orbital depends on the $d^{n}$ electronic configuration of the complex, and it is quite possible for it to be doubly occupied. In that case, the $\pi$ interaction is destabilizing on balance, since it involves four electrons, and moreover the electron transfer from the ligand to the metal in the bonding MO is now compensated by electron transfer from the metal to the ligand in the antibonding MO (3-7): overall, there is no longer any electron transfer between the two centres.


3-7


3-8
In this case, therefore, the ligand can no longer play the role of a $\pi$ donor, since there is now no empty $d$ orbital on the metal into which some of its electron density can be transferred. However, this ligand is still described as a $\pi$ donor, to indicate that it possesses at least one $\pi$-type orbital that is capable of electron transfer to a metal centre, so long as the latter can act as an acceptor. Notice that the rule given above about the consequences of the interaction with a $\pi$-donor ligand on the d block is valid for any $d^{n}$ electronic configuration of the complex: since the ligand orbital is at lower energy than the $d$ orbital with which it interacts, the new $d$ orbital, whether it is empty or occupied, is always destabilized by an antibonding interaction with the ligand $\pi$ orbital.

### 3.1.4. A first example: the octahedral complex $\left[\mathrm{ML}_{5} \mathrm{Cl}\right]$

Consider an octahedral complex with one Cl ligand (or any other halogen) which is a double-face $\pi$-donor (§ 3.1.2) (3-8); the five other ligands only have $\sigma$ interactions with the metal.
${ }^{2}$ In particular, the orbitals derived from the $\mathrm{e}_{g}$ block of the regular octahedron, where we find the antibonding $\sigma$-type interactions with the ligands, are no longer degenerate in the lower-symmetry complex [ $\mathrm{ML}_{5} \mathrm{Cl}$ ].



In order to construct the d block of this complex, we start from the orbitals that were obtained by only considering $\sigma$ interactions, then we perturb them by introducing the $\pi$ interactions. The d block derived from these $\sigma$ interactions will be assumed identical to that already established for a fully octahedral complex $\left(\mathrm{ML}_{6}\right)$, that is, split into two degenerate groups ( $\mathrm{t}_{2 g}$ and $\mathrm{e}_{g}$, Chapter 2, § 2.1.2.4). This is of course a simplification, since in an $\left[\mathrm{ML}_{5} \mathrm{Cl}\right]$ complex, the $\sigma$ interaction created by one of the L ligands is different from that involving the $\mathrm{Cl}^{2}{ }^{2}$

### 3.1.4.1. Perturbation of the block

By examining the symmetry properties of the $d$ orbitals of the regular octahedron as well as those of the $\pi$ donor, we can predict which ones will be perturbed by the $\pi$ interactions. In the case of the [ML5 Cl ] complex (3-8), it is convenient to use the symmetry planes $x z\left(P_{1}\right)$ and $y z\left(P_{2}\right)(3-9)$.

The $x y$ orbital is antisymmetric with respect to the two planes (AA), $x z$ is symmetric with respect to $P_{1}$ but antisymmetric with respect to $P_{2}$ (SA), while $y z$ is antisymmetric with respect to $P_{1}$ but symmetric with respect to $P_{2}(\mathbf{A S})$. The antibonding orbitals, $x^{2}-y^{2}$ and $z^{2}$, are symmetric with respect to both planes ( $\mathbf{S S}$ ). On the chlorine atom, the lone pairs ( $p_{x}$ and $p_{y}$ ) have $\mathbf{S A}$ and AS symmetries, respectively (3-10).


We can now construct the interaction diagram between the orbitals of the d block and those of the $\pi$-donor ligand (Figure 3.2), by combining orbitals with the same symmetry, $p_{x}$ with $x z$, and $p_{y}$ with $y z$.

Figure 3.2. Interaction diagram showing the perturbation of the d block of an octahedral complex ( $\sigma$ interactions only, left-hand side) by the two orbitals of a double-face $\pi$-donor ligand ( Cl , for example, right-hand side). The electronic occupation shown corresponds to a $d^{0}$ electronic configuration.

[^2]${ }^{4}$ Notice that in the point group of this complex ( $C_{4 v}$ ), the degenerate MO have e symmetry, while $x y, x^{2}-y^{2}$, and $z^{2}$ have $\mathrm{b}_{2}$, $\mathrm{b}_{1}$, and $\mathrm{a}_{1}$ symmetries, respectively.


These two interactions are equivalent in the sense that they involve orbitals that are related by a rotation of $90^{\circ}: x z$ and $y z$ on the one hand, $p_{x}$ and $p_{y}$ on the other. They therefore lead to the formation of two degenerate bonding MO and two antibonding MO that are also degenerate. As we showed in § 3.1.3, the antibonding MO $\left(x z-\lambda p_{x}\right.$ and $y z-\lambda p_{y}$, mainly concentrated on the metal $(\lambda<1)$ ), ${ }^{3}$ are the ones which belong to the d block of the $\left[\mathrm{ML}_{5} \mathrm{Cl}\right]$ complex. The presence of a chloride ligand therefore lifts the degeneracy of the three initial orbitals in the $t_{2 g}$ block of the octahedron, destabilizing two of them while leaving the third unchanged.

The two remaining orbitals in the $\operatorname{dblock}\left(x^{2}-y^{2}\right.$ and $\left.z^{2}\right)$ are identical to the initial MO derived by considering only the $\sigma$ interactions, since, by symmetry, they are not involved in the $\pi$ interactions. ${ }^{4}$ The two lowest-energy MO ( $p_{x}+\lambda x z$ and $p_{y}+\lambda y z$ ) are mainly localized on the chlorine and characterize the two $p$ lone pairs that are stabilized by a bonding interaction with the $d$ orbitals of the same symmetry. The occupation of the $d$ block naturally depends on the exact nature of the $\left[\mathrm{ML}_{5} \mathrm{Cl}\right]$ complex; Figure 3.2 corresponds to a complex with a $d^{0}$ electronic configuration.

### 3.1.4.2. The influence of the electronic configuration on the $\pi$-donor character of the Cl ligand

The interaction scheme given in Figure 3.2 does not depend on the $d^{n}$ electronic configuration of the $\left[\mathrm{ML}_{5} \mathrm{Cl}\right]$ complex: starting from the d block of the octahedron, two of the three nonbonding orbitals will always be destabilized while the three other orbitals are unaffected. However, the electron transfer resulting from the $\pi$ interactions does depend on the electronic occupation of the $d$ block.

5 In this description, the carbene is considered as an L ligand, with two electrons in the $2 a_{1}$ orbital but where the $p$ orbital of $b_{1}$ symmetry is empty (see Figure 1.5, Chapter 1, for the MO of a bent $\mathrm{AH}_{2}$ species). In some cases, it can also be described as an $\mathrm{X}_{2}$ ligand, with one electron in each of the $a_{1}$ and $b_{1}$ orbitals (see Chapter 4, § 4.3).

In a $d^{0}$ complex (Figure 3.2), the two lone pairs $p_{x}$ and $p_{y}$ are partially delocalized on to the metal after the interaction (bonding MO $p_{x}+\lambda x z$ and $p_{y}+\lambda y z$ ), and each interaction results in some electron transfer from the ligand to the metal. Since all the other orbitals are empty, there is no transfer in the opposite sense: in an $\left[\mathrm{ML}_{5} \mathrm{Cl}\right]$ complex with a $d^{0}$ electronic configuration, the chloride does indeed play the role of a double-face $\pi$-donor. A $d^{2}$ complex behaves in the same way: the two electrons in the d block occupy the nonbonding $x y$ orbital after the interaction, and therefore remain completely localized on the metal.

For the $d^{3}-d^{6}$ electronic configurations, the $x z$ and $y z$ orbitals, which were nonbonding before the interaction, are progressively filled, leading to some electron transfer in the opposite direction, from the metal to the ligand. So these orbitals, which were localized on the metal, become partially delocalized onto the chloride ligand after the $\pi$ interaction (antibonding MO $x z-\lambda p_{x}$ and $y z-\lambda p_{y}$ ). In the limiting case of a low-spin $d^{6}$ complex, in which the three lowest $d$ levels are all doubly occupied, the $\pi$ interactions no longer lead to any electron transfer from the chloride ligand to the metal. The $\mathrm{Cl} \rightarrow$ metal transfers in the bonding MO $p_{x}+\lambda x z$ and $p_{y}+\lambda y z$ are balanced by the metal $\rightarrow \mathrm{Cl}$ transfers in the antibonding MO $x z-\lambda p_{x}$ and $y z-\lambda p_{y}$. In other words, the metal in the $d^{6}$ electronic configuration has no empty orbital available to accept the $\pi$ electrons of the chloride ligand. Thus, the only consequence of the presence of a double-face $\pi$-donor ligand is the destabilization of two of the three nonbonding levels in the octahedron.

## 3.2. $\pi$-acceptor ligands: general properties

### 3.2.1. The nature of the $\pi$ orbital on the ligand

$\pi$-acceptor ligands are characterized by the existence of at least one empty orbital that can have a $\pi$ interaction with a $d$ orbital on the metal. If this orbital is empty, it must be fairly high in energy. It may be:

1. A $p$ orbital on a fairly electropositive atom, such as the boron $p$ orbital in boryl complexes or the carbon $p$ orbital in carbene complexes ${ }^{5}$ (3-11).



empty $p$ orbital
2. A $\pi^{*}$ antibonding orbital associated with a $\pi$ bond between two atoms of the ligand. This orbital is therefore delocalized onto two centres, but its overlap with the metal $d$ orbital is large if it is mainly concentrated on the atom that is bonded to the metal. This is what happens when the metal is less electronegative than the second centre involved in the double bond (Chapter 1, §1.2.2), as is the case, for example, for the formyl ligand $\mathrm{HC}=\mathrm{O}(\mathbf{3 - 1 2})$.


3-12

### 3.2.2. 'Single-face' and 'double-face' $\pi$-acceptors

As for $\pi$-donor ligands, we distinguish $\pi$-acceptor ligands which possess a single empty $p$ or $\pi^{*}$ orbital (single-face acceptors) from those which have two, pointing in two perpendicular directions (double-face acceptors). The boryl, carbene (3-11), and formyl (3-12) ligands, for example, are all in the first category.

Among double-face $\pi$-acceptors, we may cite the $\mathrm{B}-\mathrm{H}$ ligand, with its two nonbonding empty $p$ orbitals (Chapter $1, \S$ 1.5.2.2). There is another which is particularly important in organometallic chemistry, the carbonyl ligand, $\mathrm{C} \equiv \mathrm{O}$, which we shall study in this section. The principal molecular orbitals of CO that are involved in the interaction with a metal centre are shown in Figure 3.3 (see Chapter 1, § 1.5.2.4). They are the three highest-energy occupied MO, the $\pi_{\mathrm{CO}}$ bonding orbitals and the nonbonding $\sigma_{\mathrm{C}}$ orbital that essentially characterizes the lone pair on the carbon atom, and the two lowest-energy empty MO, the $\pi_{\mathrm{CO}}^{*}$ antibonding orbitals. The $\pi_{\mathrm{CO}}$ orbitals are mainly concentrated on the oxygen, which is the more electronegative centre, and the $\pi_{\mathrm{CO}}^{*}$ orbitals


Figure 3.3. Electronic structure of the carbonyl ligand $(\mathrm{C} \equiv \mathrm{O})$ : the three highest occupied and the two lowest unoccupied orbitals.
on the carbon. The nonbonding $\sigma_{\mathrm{C}}$ orbital, which is the highest occupied molecular orbital (HOMO) on the ligand, forms a $\sigma$ bond with the metal. As a result, the carbonyl ligand usually binds to a transition metal in the $\eta^{1}$ mode, with a linear $\mathrm{M}-\mathrm{C} \equiv \mathrm{O}$ arrangement that maximizes the $\sigma$ overlap. The $\pi$ interactions, that involve both the $\pi_{\mathrm{CO}}$ bonding MO and the $\pi_{\mathrm{CO}}^{*}$ antibonding MO, are added to this $\sigma$ interaction. This example illustrates a problem posed by all the ligands with a double or triple bond that are coordinated to the metal in the $\eta^{1}$ mode (i.e. through one of the atoms involved in the multiple bond): using their occupied $\pi$ orbital(s), they can act as a $\pi$ donor, whereas their empty $\pi^{*}$ orbitals can allow them to act as a $\pi$ acceptor. Depending on the relative strength of these two interactions, the ligand can act overall either as a $\pi$ donor or as a $\pi$ acceptor.

To analyse this probem in the case of the carbonyl ligand (see Appendix C for more details), we consider a linear $\mathrm{M}-\mathrm{C} \equiv \mathrm{O}$ arrangement and one of the $d$ orbitals on $\mathrm{M}(y z, 3-13)$ that can interact with the $\pi$ system of the carbonyl ligand. It is easy to see that this $d$ orbital has a non-zero overlap with a pair of $\pi$ and $\pi^{*}$ orbitals on CO, those that are oriented parallel to the $z$-axis, leading to a three-orbital interaction. The $d \leftrightarrow \pi$ and $d \leftrightarrow \pi^{*}$ interactions (3-13), which involve orbitals with different energies, are proportional to $S^{2} / \Delta \varepsilon$ (Chapter 1, $\S 1.3 .2$ ). They are


3-13
therefore strengthened by a small orbital energy difference $(\Delta \varepsilon)$ and by a large overlap $(S)$. The overlap term clearly favours the interaction with $\pi_{\mathrm{CO}}^{*}\left(S_{\pi}^{*}>S_{\pi}\right.$, due to the relative sizes of the coefficients on carbon in the $\pi_{\mathrm{CO}}^{*}$ and $\pi_{\mathrm{CO}} \mathrm{MO}$ ). Moreover, it can be shown (see appendix) that the $d$ orbital is closer in energy to the antibonding orbital than to the bonding
$\operatorname{orbital}\left(\Delta \varepsilon_{\pi}^{*}<\Delta \varepsilon_{\pi}, \mathbf{3 - 1 3}\right)$ for almost all the transition elements. In the three-orbital interaction scheme shown in 3-13, both factors, overlap and energy difference, therefore favour the $d \leftrightarrow \pi^{*}$ interaction. We may therefore simplify the diagram, by neglecting the relatively weak interaction ( $d \leftrightarrow \pi$ ) and retaining only the dominant interaction $\left(d \leftrightarrow \pi^{*}\right.$ ). This latter
 interaction, which involves an empty orbital on the carbonyl ligand, therefore gives it $\pi$-acceptor character.

In this simplified model, we therefore consider that the $\pi$ interactions with the carbonyl ligand involve only the two $\pi_{\text {CO }}^{*}$ orbitals; it is thus a double-face $\pi$-acceptor (3-14).

### 3.2.3. Perturbation of the $d$ orbitals: the general interaction diagram

When the unoccupied orbital on the acceptor ligand is a $\pi^{*}$ antibonding MO, its energy level is almost always higher than that of the nonbonding or weakly antibonding orbitals in the $d$ block. ${ }^{6}$ The same remark may be made about a nonbonding $p$ orbital on a very electropositive centre, such as boron in boryl complexes (3-11). But the relative orbital energies are less clear when the centre bearing the empty $p$ orbital is moderately electronegative, such as carbon in the carbene complexes $\left[\mathrm{L}_{n} \mathrm{M}-\mathrm{CR}_{2}\right](\mathbf{3 - 1 1})$. They then depend on the nature of the metal and of the substituents on the carbene ligand (see Chapter 4, § 4.3). At this stage of our analysis, we shall only consider the more common case where the energy level of the $p$ or $\pi^{*}$ orbitals on the $\pi$-acceptor ligand is higher than that of the nonbonding or weakly antibonding orbitals in the $d$ block.

As an example, we consider a $\pi$-acceptor ligand whose nonbonding $p$ orbital interacts with a doubly occupied $d$ orbital (Figure 3.4). Due to the relative energies of the initial orbitals, the bonding combination is mainly concentrated on the metal and the antibonding combination on the ligand. As in the preceding examples, the new d-block orbital is mainly concentrated on the metal, that is, in this case, the bonding combination of the initial orbitals.

Rule: the $\pi$ interaction between a d-block orbital and the orbital of a $\pi$-acceptor ligand leads to a stabilization of the d-block orbital, by mixing with the ligand orbital in a bonding sense.

Figure 3.4 shows the resulting $\pi$ electron transfer involving the ligand and the metal. Before the interaction, the two electrons were localized in the metal $d$ orbital. After the interaction, these electrons occupy the bonding molecular orbital which is partially located on

Figure 3.4. Sketch of the interaction between a doubly occupied $d$ orbital and the unoccupied orbital of a $\pi$ acceptor (X).

the ligand. This delocalization is thus accompanied by some electron transfer from the metal to the $\pi$-acceptor ligand.

### 3.2.4. A first example: the octahedral complex $\left[\mathrm{ML}_{5} \mathrm{CO}\right]$

### 3.2.4.1. Interaction diagram

We consider an octahedral complex with one carbonyl ligand, a doubleface $\pi$-acceptor, and five other ligands that only have $\sigma$ interactions with the metal (3-15).

To construct the d block of this complex, we shall follow the same procedure as that used for a complex with a $\pi$ donor (§ 3.1.4): we start from the orbitals for an ideal octahedral $\mathrm{ML}_{6}$ complex in which there are only $\sigma$ interactions, then we perturb these orbitals by 'switching on' the $\pi$ interactions with the carbonyl ligand.

We have already established (see 3-10) the symmetries of the metal $d$ orbitals with respect to the planes of symmetry $P_{1}$ and $P_{2}(\mathbf{3 - 1 5}): x y(\mathbf{A A})$, $x z(\mathbf{S A}), y z(\mathbf{A S}), x^{2}-y^{2}$ and $z^{2}(\mathbf{S S})$. Following the analysis developed in $\S$ 3.2.2, only the empty antibonding $\pi_{\mathrm{CO}}^{*}$ orbitals are treated. They are oriented parallel to the $x$ - and $y$-axes, and have $\mathbf{S A}$ and AS symmetries, respectively (3-16).

We can now construct the diagram for the interaction between the orbitals of the $d$ block and the $\pi_{\mathrm{CO}}$ orbitals on the carbonyl ligand (Figure 3.5), by combining the $\pi_{x}^{*}$ and $x z$ orbitals, with SA symmetry, and the $\pi_{y}^{*}$ and $y z$ orbitals, with AS symmetry. These are two equivalent interactions, since they involve orbitals which are related by a rotation of $90^{\circ}:(x z, y z)$ on the one hand and $\left(\pi_{x}^{*}, \pi_{y}^{*}\right)$ on the other. They therefore lead to the formation of two degenerate bonding MO and two antibonding MO that are also degenerate. As shown in § 3.2.3, the bonding $\mathrm{MO}\left(x z+\lambda \pi_{x}^{*}\right.$ and $\left.y z+\lambda \pi_{y}^{*}\right)$, which are mainly concentrated on the metal, are the orbitals which belong to the $d$ block of the [ $\mathrm{ML}_{5} \mathrm{CO}$ ] complex. The presence of a CO ligand therefore lifts the degeneracy of

Figure 3.5. Interaction diagram showing the perturbation of the $d$ block of an octahedral complex (left-hand side, $\sigma$ interactions only) by the two $\pi_{\mathrm{CO}}^{*}$ orbitals of a carbonyl ligand (right-hand side). The electronic occupation shown corresponds to a complex with a $d^{6}$ electronic configuration.

the three orbitals in the original $\mathrm{t}_{2 g}$ block of the octahedron, stabilizing two of them but leaving the third $(x y)$ unchanged. The two remaining orbitals in the d block $\left(x^{2}-y^{2}\right.$ and $\left.z^{2}\right)$ are identical to the original MO derived for $\sigma$ interactions only, since they are not affected by the $\pi$ interactions. Finally, the two highest-energy $\mathrm{MO}\left(\pi_{x}^{*}-\lambda x z\right.$ and $\left.\pi_{y}^{*}-\lambda y z\right)$ are mainly concentrated on the carbonyl ligand. The occupation of the d block clearly depends on the exact nature of the [ML; CO ] complex; Figure 3.5 corresponds to a complex with a $d^{6}$ electronic configuration. Before the interaction, the four electrons stabilized by the $\pi$ interactions occupied pure metal $d$ orbitals ( $x z$ and $y z$ ). After the interaction, these orbitals are partially delocalized onto the carbonyl ligand, showing the $\pi$-acceptor character of this species.

## Comment

The relative position of the $\sigma$-antibonding $\left(x^{2}-y^{2}, z^{2}\right)$ and the $\pi$-antibonding $\left(\pi_{x}^{*}-\lambda x z\right.$ and $\left.\pi_{y}^{*}-\lambda y z\right)$ orbitals may be inverted with respect to that shown in Figure 3.5. It depends among other things on the energy gap between the nonbonding ( $x y, x z$, and $y z$ ) and the antibonding ( $x^{2}-y^{2}, z^{2}$ ) orbitals created by the $\sigma$ interactions.

### 3.2.4.2. The metal-carbonyl bond: donation and back-donation interactions

The $\sigma$ bond between the carbonyl ligand and the metal centre is formed by the $\sigma_{\mathrm{C}}$ orbital that characterizes the lone pair on the carbon atom (Figure 3.3) and by a $d$ orbital on the metal ( $z^{2}$ in the orientation shown
in 3-15). This interaction, which leads to a transfer of electrons from the ligand to the metal, is called the donation interaction (3-17a). The $\pi$ interactions described in Figure 3.5 produce a transfer of electrons in the opposite direction, from the metal to the ligand. These are back-donation interactions ( $\mathbf{3 - 1 7 b}$ ). The carbonyl ligand is therefore simultaneously a $\sigma$-donor and a $\pi$-acceptor.


3-17a Donation


3-17b Back donation

The $\pi$ interactions reinforce the metal-carbon bond (bonding interactions $\mathrm{M}-\mathrm{C}$ ) but weaken the CO bond (antibonding interactions $\mathrm{C}-\mathrm{O}$ ) $(\mathbf{3 - 1 7 b})$. This electronic reorganization can be represented by the Lewis structures shown in 3-18.


Experimentally, the metal-carbon bond is observed to be substantially shorter than the value that would be expected for a $\sigma$-only bond, by some $0.2-0.3 \AA$. The result for the CO bond is less clear, as the distances measured for carbonyl complexes fall in the very narrow range of $1.14-1.15 \AA$. These values are only slightly longer than that found for the free ligand $(1.13 \AA)$. The equilibrium CO distance is therefore almost insensitive to the transfer of electrons, an observation that can at least partially be explained by the fact that the difference in length between a triple and a double bond is also relatively small (1.13 and $1.23 \AA$ A , respectively). Infrared spectroscopy (IR) provides a far better probe, from the value of the IR absorption frequency associated with the stretching of the CO bond $\left(v_{\mathrm{CO}}\right)$. This frequency is related to the strength of the bond and is very sensitive to the electronic population in the $\pi_{\mathrm{CO}}^{*}$ orbitals. For example, for the isolated ligand it decreases from 2143 to $1489 \mathrm{~cm}^{-1}$ on passing from the ground state to the excited electronic state in which an electron has been excited from the nonbonding $\sigma_{\mathrm{C}}$ orbital to the antibonding $\pi_{\mathrm{CO}}^{*}$ orbital (Figure 3.3). This is therefore a very sensitive indicator which enables the transfer of metal $\pi$ electrons to one or several carbonyl ligands to be demonstrated. Thus, $v_{\mathrm{CO}}$ decreases from $2143 \mathrm{~cm}^{-1}$ in isolated CO to $2000 \mathrm{~cm}^{-1}$ in $\left[\mathrm{Cr}(\mathrm{CO})_{6}\right]$, a complex with a $d^{6}$ electronic configuration, and even to $1860 \mathrm{~cm}^{-1}$ in $\left[\mathrm{V}(\mathrm{CO})_{6}\right]^{-}$, another $d^{6}$ complex which, due to its anionic nature, is
particularly effective for the transfer of electron density to $\pi$-acceptor ligands.

### 3.3. Complexes with several $\pi$-donor or $\pi$-acceptor ligands

In the first example that we considered, there was only a single ligand able to perturb the d block of the complex through $\pi$ interactions. When several ligands of this type are present, their individual effects are added, destabilizing the $d$ orbitals in the case of $\pi$-donors (antibonding mixing in the d block) or stabilizing them in the case of $\pi$-acceptors (mixing of a bonding type in the d block). However, as we shall see in the following examples, the effect produced by $n$ ligands of a given type is not necessarily $n$ times larger than that produced by a single ligand. The symmetry properties of the complex may prevent some orbitals, either occupied for $\pi$ donors or empty for $\pi$ acceptors, from interacting with the d-block orbitals.

### 3.3.1. The trans- $\left[\mathrm{ML}_{4} \mathrm{Cl}_{2}\right]$ octahedral complex



Consider an octahedral complex of the type [ $\mathrm{ML}_{4} \mathrm{Cl}_{2}$ ], with two doubleface $\pi$-donor ligands ( Cl ) in trans positions and four other ligands which only have $\sigma$ interactions with the metal. Three of the symmetry elements in the complex are shown in 3-19: the planes $P_{1}(x z)$ and $P_{2}(y z)$ and the inversion centre $i$, located on the central atom.

### 3.3.1.1. $\pi$-type orbitals on the ligands

Four lone pairs of $\pi$-type are available, described by the $p_{x}$ and $p_{y}$ orbitals on each chlorine atom. It must be noticed that the two Cl ligands are equivalent by symmetry, that is, they are interchanged by at least one symmetry element (e.g. the horizontal plane $x y$ or the inversion centre $i$ ).


Rather than consider each of the $p_{x}$ and $p_{y}$ orbitals individually, it is therefore preferable to use linear combinations of these orbitals that properly

7 The bonding or antibonding character is weak, due to the large separation between the two chlorine atoms. It is perhaps better to replace 'bonding/antibonding' by 'in-phase/ out-of-phase', to emphasize only the way in which the initial atomic orbitals are combined.
reflect the symmetry of the complex (Chapter $6, \S 6.4$ ). In the case of two equivalent atoms, these are simply 'bonding' and 'antibonding' ${ }^{7}$ combinations of the initial orbitals. These are usually indicated by ( $\pi_{x}$, $\pi_{y}$ ) and ( $\pi_{x}^{*}, \pi_{y}^{*}$ ), respectively (3-20), making reference to the nature of the overlap $(\pi)$ and the orientation of the orbitals ( $x$ - or $y$-axes).

Electronically speaking, these four orbitals are doubly occupied, since they are formed from atomic orbitals that describe $\pi$-type lone pairs on each of the chloride ligands. They therefore provide a delocalized description of these lone pairs, which is adapted to the symmetry of the complex.

Before constructing the diagram for the interaction between the $\pi$ orbitals on the ligands and the $d$ orbitals on the metal centre, we must analyse the symmetry properties of these orbitals. We shall describe two methods, in which we use either some selected symmetry elements, or the full set of these elements and the machinery of group theory.

### 3.3.1.2. Taking symmetry into account: initial analysis

To characterize orbital symmetry, we may, as in the examples already treated (§ 3.1.4 and 3.2.4), use the planes $P_{1}$ and $P_{2}(\mathbf{3 - 1 9})$. But we must add a third element to enable us to distinguish the symmetry of the bonding and antibonding combinations that we have constructed above. This additional element might be the $x y$ plane, or it might be the inversion centre $i$. With respect to this latter element, the $\pi_{x, y}$ orbitals are antisymmetric (A) but the $\pi_{x, y}^{*}$ orbitals are symmetric (S) (see 3-21 for the $\pi_{y}$ and $\pi_{y}^{*}$ orbitals).


When we take all three of these elements into account ( $P_{1}, P_{2}$, and $i$ ), the symmetry labels are SAA for $\pi_{x}$, ASA for $\pi_{y}, \mathbf{S A S}$ for $\pi_{x}^{*}$, and ASS
for $\pi_{y}^{*}$; on the metal, they are AAS for $x y$, SAS for $x z$, ASS for $y z$, and SSS for both $x^{2}-y^{2}$ and $z^{2}$ (3-22).

This initial analysis of the symmetry properties shows us that two of the four lone pairs of the chloride ligands, those described by the bonding combinations $\pi_{x}$ and $\pi_{y}$, cannot interact with the $d$-block orbitals, due to their different symmetry properties $(S=0)$. In the same way, the $x y$, $x^{2}-y^{2}$, and $z^{2}$ orbitals on the metal cannot, by symmetry, interact with the ligand orbitals. The only interactions that are possible $(S \neq 0)$ occur between $\pi_{x}^{*}$ and $x z$ on the one hand (symmetry SAS), and between $\pi_{y}^{*}$ and $y z$ on the other (symmetry ASS).


### 3.3.1.3. Taking symmetry into account: the use of group theory

Rather than select only some symmetry elements present in the system being studied, one can consider all of them by using group theory. This is the most rigorous method-there is no danger of forgetting a symmetry element that might be important for a particular interaction-and it allows us, particularly in high-symmetry systems, to anticipate certain special properties such as the existence of orbitals that are degenerate by symmetry. The symmetry elements of the complex studied here are characteristic of the $D_{4 h}$ point group (Chapter 6, § 6.2.2 and 6.6.1).

If we consider the two $\pi^{*}$ antibonding combinations of the lone pairs, we notice that they are related by a rotation of $90^{\circ}$ about the $z$-axis, which is a symmetry element of the system (a $C_{4}$-axis). The same applies for the two $\pi$-bonding combinations, and also for the $x z$ and $y z$ orbitals on the metal. These are pairs of orbitals that are degenerate by symmetry in the $D_{4 h}$ point group. If we consult the character table for

Table 3.1. Character table for the $D_{4 h}$ point group

| $D_{4 h}$ | $E$ | $2 C_{4}$ | $C_{2}$ | $2 C_{2}^{\prime}$ | $2 C_{2}^{\prime \prime}$ | $i$ | $2 S_{4}$ | $\sigma_{h}$ | $2 \sigma_{v}$ | $2 \sigma_{d}$ |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\mathrm{~A}_{1 g}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |  |
| $\mathrm{~A}_{2 g}$ | 1 | 1 | 1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 |  |
| $\mathrm{~B}_{1 g}$ | 1 | -1 | 1 | 1 | -1 | 1 | -1 | 1 | 1 | -1 | $x^{2}+y^{2}, z^{2}$ |
| $\mathrm{~B}_{2 g}$ | 1 | -1 | 1 | -1 | 1 | 1 | -1 | 1 | -1 | 1 | $x^{2}-y^{2}$ |
| $\mathrm{E}_{g}$ | 2 | 0 | -2 | 0 | 0 | 2 | 0 | -2 | 0 | 0 |  |
| $\mathrm{~A}_{1 u}$ | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 |  |
| $\mathrm{~A}_{2 u}$ | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 | 1 | 1 | $z$ |
| $\mathrm{~B}_{1 u}$ | 1 | -1 | 1 | 1 | -1 | -1 | 1 | -1 | -1 | 1 |  |
| $\mathrm{~B}_{2 u}$ | 1 | -1 | 1 | -1 | 1 | -1 | 1 | -1 | 1 | -1 |  |
| $\mathrm{E}_{u}$ | 2 | 0 | -2 | 0 | 0 | -2 | 0 | 2 | 0 | 0 | $(x, y)$ |

this group (Table 3.1), we notice that it contains two two-dimensional representations, referred to as $\mathrm{E}_{u}$ and $\mathrm{E}_{g}$, which differ by their behaviour with respect to the inversion centre of the complex ( $i$, located on the central metal): $\mathrm{E}_{g}$ is symmetric (its character in the $i$ column is positive), but $\mathrm{E}_{u}$ is antisymmetric (a negative character). The bonding combinations $\left(\pi_{x}, \pi_{y}\right)$, which are antisymmetric with respect to the inversion centre (3-21), have $\mathrm{e}_{u}$ symmetry, but the antibonding combinations ( $\pi_{x}^{*}$, $\pi_{y}^{*}$ ), which are symmetric with respect to the inversion centre, have $\mathrm{e}_{g}$ symmetry. The character table also gives us the symmetries of the $d$ orbitals on the metal centre, in the last column: $\mathrm{a}_{1 g}$ for $z^{2}, \mathrm{~b}_{1 g}$ for $x^{2}-y^{2}, \mathrm{~b}_{2 g}$ for $x y$, and $\mathrm{e}_{g}$ for $(x z, y z)$.

In the case being examined, we therefore come to the same conclusion as that established in the preceding section from a limited number of symmetry elements: the only interactions that occur concern the $\left(\pi_{x}^{*}, \pi_{y}^{*}\right)$ orbitals on the ligands, and $(x z, y z)$ on the metal centre, which constitute two degenerate pairs of orbitals with $\mathrm{e}_{g}$ symmetry in the $D_{4 h}$ point group.

### 3.3.1.4. The interaction diagram

We are now in a position to construct the diagram which describes the $\pi$ interactions in a trans-[ $\left.\mathrm{ML}_{4} \mathrm{Cl}_{2}\right]$ complex (Figure 3.6). We shall adopt the notations of group theory to represent the orbitals' symmetries. As in the preceding examples ( $\S 3.1 .4$ and 3.2.4), we suppose that the d block of the complex, before the interaction, is similar to that of a regular octahedral complex, with three nonbonding degenerate orbitals ( $x y, x z$, and $y z$ with the axes defined in 3-19) and two degenerate antibonding orbitals, $x^{2}-y^{2}$ and $z^{2}$ (left-hand side of Figure 3.6). The first three orbitals, which formed the $t_{2 g}$ block in the regular octahedron, have $\mathrm{b}_{2 g}(x y)$ and $\mathrm{e}_{g}(x z, y z)$ symmetries in the $D_{4 h}$ point group. The symmetries of the antibonding orbitals derived from the $\mathrm{e}_{g}$ block in the octahedron become

Figure 3.6. Interaction diagram showing the perturbation of the $d$ block of an octahedral complex ( $\sigma$ interactions only, left-hand side) by the lone pairs of two double-face $\pi$-donor ligands ( Cl , for example, right-hand side) in trans positions. The electronic occupation shown corresponds to a complex with a $d^{0}$ electronic configuration.

$\mathrm{a}_{1 g}$ for $z^{2}$ and $\mathrm{b}_{1 g}$ for $x^{2}-y^{2}$. The occupation of these orbitals depends on the $d^{n}$ electronic configuration of the complex, the example given in Figure 3.6 corresponding to a $d^{0}$ complex whose $d$ block is empty. We find the $\pi$-type orbitals of the chlorine atoms on the right-hand side of the figure (3-20). The $\mathrm{e}_{u}$ bonding combinations are placed slightly lower in energy than the $\mathrm{e}_{g}$ antibonding combinations, though the energy difference is only slight, due to the large separation between the two centres. The energies of these four orbitals are therefore close to that of a pure $p$ orbital on chlorine. As a result, they are placed at a lower energy level than the nonbonding $d$ orbitals on the metal. Notice that these $\pi$ and $\pi^{*}$ orbitals are all doubly occupied, and that they offer a delocalized description of the four $\pi$-type lone pairs on the chloride ligands (two per ligand).

For symmetry reasons, as we have already established, the two lone pairs described by the orbitals with $\mathrm{e}_{\mathfrak{u}}$ symmetry (3-20) cannot interact with the d-block orbitals. The only interactions that are possible $(S \neq 0)$ occur between the lone pairs with $\mathrm{e}_{g}$ symmetry and the $(x z, y z)$ orbitals on the metal (Figure 3.6). This interaction stabilizes four electrons from the lone pairs (the $\pi_{x}^{*}+\lambda x z$ and $\pi_{y}^{*}+\lambda y z$ orbitals, mainly concentrated on the Cl ligands) but destabilizes two of the three octahedral nonbonding orbitals $\left(x z-\lambda \pi_{x}^{*}\right.$ and $y z-\lambda \pi_{y}^{*}$, mainly concentrated on the metal). This destabilisation is larger than that produced by a single chloride ligand ([ $\left.\mathrm{ML}_{5} \mathrm{Cl}\right]$ complex, Figure 3.2) since there are now two $\pi$-type antibonding interactions for each of these orbitals. Two of the lone pairs, and therefore four electrons, are not affected by these interactions.

The new d block of the complex, as always, contains five orbitals that are mainly composed of the $d$ orbitals on the metal: $x y$, nonbonding, $x z-\lambda \pi_{x}^{*}$ and $y z-\lambda \pi_{y}^{*}, \pi$ antibonding, and $x^{2}-y^{2}$ and $z^{2}, \sigma$ antibonding.

This example shows us that for symmetry reasons, the $\pi$ interactions are not simply additive when the number of ligands is increased. Two lone pairs are involved in both the monochloro octahedral complex (Figure 3.2) and the trans dichloro complex, although four lone pairs are a priori available in the latter case. The difference between the two complexes concerns the magnitudes of the destabilization of the $x z$ and $y z$ orbitals of the d block and the stabilization of the lone pairs that take part in the interactions, both of which are larger in the dichloro compound. It must also be noted that for a given type of complex (octahedral, for example) and a given number of $\pi$-donor (or -acceptor) ligands, the symmetry properties of the complex, and therefore the metal-ligand interactions which occur, also depend on the arrangement of these ligands. This point is illustrated in Exercise 3.4, which concerns an octahedral $\left[\mathrm{ML}_{4} \mathrm{Cl}_{2}\right]$ complex in which the two chloride ligands are now in cis positions.

### 3.3.2. The trans-[ $\left.\mathrm{ML}_{4}(\mathrm{CO})_{2}\right]$ octahedral complex

This complex has the same symmetry as the preceding one (the $D_{4 h}$ point group). Since the carbonyl groups are $\pi$-acceptors, we consider the two empty $\pi_{x}^{*}$ and $\pi_{y}^{*}$ orbitals on each, and these are combined in pairs to form the symmetry-adapted orbitals for the complex (3-23). The bonding combinations, $\left(\pi_{x}^{*(+)}\right.$ and $\left.\pi_{y}^{*(+)}\right)$, which are antisymmetric with respect to the inversion centre, have $\mathrm{e}_{u}$ symmetry, whereas the antibonding combinations, $\left(\pi_{x}^{*(-)}\right.$ and $\left.\pi_{y}^{*(-)}\right)$, which are symmetric with respect to this operation, have $\mathrm{e}_{g}$ symmetry.


3-23
Since the overlaps between the $\pi_{x}^{*}$ (or $\pi_{y}^{*}$ ) orbitals are very small, due to the large distance between the ligands, the energies of the bonding

Figure 3.7. Interaction diagram showing the perturbation of the d block of an octahedral complex ( $\sigma$ interactions only, left-hand side) by the $\pi^{*}$ orbitals of two double-face $\pi$-acceptor carbonyl ligands (right-hand side) in trans positions. The electronic occupation shown corresponds to a complex with a $d^{6}$ electronic configuration.

( $\mathrm{e}_{u}$ ) and antibonding ( $\mathrm{e}_{g}$ ) combinations are close to that of a $\pi^{*}$ orbital in an isolated carbonyl ligand. As a result, they are higher in energy than the nonbonding $d$ orbitals on the metal, with the $\mathrm{e}_{u}$ orbitals being very slightly more stable than the $\mathrm{e}_{g}$ (Figure 3.7, right-hand side). These four orbitals are, of course, empty. On the left-hand side of the figure, we find the $d$-block orbitals, assumed to be those of a regular octahedron without any $\pi$ interactions. The occupation of these orbitals depends on the electronic configuration of the complex that is considered. In Figure 3.7, we are concerned with a $d^{6}$ complex. As in the preceding complex, only two of the four $\pi$-type orbitals on the ligands-the $\mathrm{e}_{g}$ orbitals-and two nonbonding orbitals on the metal ( $x z$ and $y z$, which also have $\mathrm{e}_{g}$ symmetry) can interact. This interaction therefore stabilizes those two d-block orbitals $\left(x z+\lambda \pi_{x}^{*(-)}\right.$ and $y z+\lambda \pi_{y}^{*(-)}$, mainly concentrated on the metal), whereas the antibonding combinations are mainly concentrated on the ligands $\left(\pi_{x}^{*(-)}-\lambda x z\right.$ and $\left.\pi_{y}^{*(-)}-\lambda y z\right)$. The new d block therefore consists of the following orbitals: $x z+\lambda \pi_{x}^{*(-)}$ and $y z+\lambda \pi_{y}^{*(-)}$, which are $\pi$-bonding, $x y$ (nonbonding), and $x^{2}-y^{2}$ and $z^{2}$, which are $\sigma$-antibonding.

In the case of a $d^{6}$ complex, four electrons from the $d$ block are stabilized. The stabilization is larger than that produced in the complex with just one carbonyl ligand (Figure 3.5), since there are two $\pi$-bonding interactions in each orbital stabilized instead of only one.

### 3.3.3. Construction of the d-block orbitals 'by hand'

The relative energies of the $\pi$-type ligand orbitals and the metal orbitals control the nature of the two interactions that occur in the d block: $(i)$ an
antibonding interaction with the occupied orbital of a $\pi$-donor, which destabilizes the $d$ orbital; (ii) a bonding interaction with the empty orbital of a $\pi$-acceptor which stabilizes the $d$ orbital. The shape of the perturbed $d$ orbitals can easily be obtained when there is a single ligand with a $\pi$ system: one only needs to combine the $d$ orbital and the $\pi$-type orbital with which it can overlap, in a bonding manner for an acceptor but in an antibonding sense for a donor. When there are two ligands of this type, we need first to construct the symmetry-adapted $\pi \mathrm{MO}$ for the complex, and then decide which of these could interact with the d orbitals. Once this is done, the same rules are applied for the perturbation of the $d$ orbitals by the symmetry-adapted orbitals on the ligands: an antibonding mixing with destabilization for $\pi$-donors, but a bonding mixture with stabilization for $\pi$-acceptors (Figures 3.6 and 3.7). In the two complexes we have studied, the shapes of the symmetry-adapted orbitals are obvious ( $\$ 3.3 .1$ and 3.3.2), but it can be much more difficult to obtain them in other cases (see, for example, Chapter 6, § 6.6.6 and Exercise 6.13). We now consider whether it is always necessary, for each type of complex (octahedral, TBP, square-planar, etc.) and for each type of substitution, to determine the symmetry-adapted combinations of ligand orbitals before being able to discover the shapes of the perturbed $d$ orbitals. In principle, it is indeed necessary to work in stages, along those lines, but we are now going to show, starting from the two preceding examples, how the main thrust of the information can be obtained more quickly.

### 3.3.3.1. The trans- $\left[\mathrm{ML}_{4} \mathrm{Cl}_{2}\right]$ and trans- $\left[\mathrm{ML}_{4}(\mathrm{CO})_{2}\right]$ complexes revisited

The $d \leftrightarrow \pi$ interactions due to a single Cl ligand were studied in the monochloro complex ([ $\left.\mathrm{ML}_{5} \mathrm{Cl}\right]$, Figure 3.2). If we now re-examine the d-block orbitals of the dichloro complex trans-[ $\mathrm{ML}_{4} \mathrm{Cl}_{2}$ ] (Figure 3.6), we notice, but after the event, that we could have obtained them directly by considering successively the contributions expected for each of the two ligands, that is, by proceeding in two stages, as shown in 3-24.


The interactions of the lone pairs $p_{x}^{(1)}$ and $p_{y}^{(1)}$ of the first substituent $\mathrm{Cl}_{1}$ leave the $x y, x^{2}-y^{2}$, and $z^{2}$ orbitals unchanged (zero overlaps),
whereas $x z$ and $y z$ are destabilized by antibonding interactions with $p_{x}^{(1)}$ and $p_{y}^{(1)}$, respectively (non-zero overlaps). Analogous observations can be made for the second substitution by $\mathrm{Cl}_{2}$, since the same overlaps are involved when we consider the $p_{x}^{(2)}$ and $p_{y}^{(2)}$ lone pairs. This two-stage decomposition of the construction of the d-block orbitals for the trans[ $\mathrm{ML}_{4} \mathrm{Cl}_{2}$ ] complex is illustrated in 3-25 for the three initially nonbonding orbitals of the octahedron. The $x y$ orbital which cannot interact with any of the lone pairs is unchanged, whereas the $x z$ and $y z$ orbitals are destabilized in the same way (two equivalent antibonding interactions).
$x y$



$x z$

$\xrightarrow{\mathrm{Cl}_{1}}$


$y z$


3-25

On can proceed in the same way for the trans $-\mathrm{ML}_{4}(\mathrm{CO})_{2}$ complex, the only difference being that the mixing with the $\pi_{C O}^{*}$ orbitals occurs in a bonding sense ( $\pi$ acceptor) (3-26).


### 3.3.3.2. mer- $\left[\mathrm{ML}_{3}(\mathrm{CO})_{3}\right]$ complexes

We shall now apply this method 'by stages' to a new octahedral complex with three carbonyl ligands and three other ligands that only have $\sigma$ interactions with the metal. The arrangement chosen for the ligands is indicated in 3-27. Of the two possible isomers for an $\left[\mathrm{ML}_{3}(\mathrm{CO})_{3}\right]$ complex, this one is the mer isomer, in contrast to the fac isomer in which the carbonyl ligands are placed at the three vertices that make up a face of the octahedron. On the $\mathrm{C}_{1} \mathrm{O}_{1}$ and $\mathrm{C}_{3} \mathrm{O}_{3}$, carbonyl ligands, the $\pi^{*}$ orbitals point parallel to the $x$ - and $z$-axes, whereas on the central ligand $\mathrm{C}_{2} \mathrm{O}_{2}$ they are parallel to the $x$-and $y$-axes (3-27). We now examine how each of the octahedral orbitals can overlap with the $\pi^{*}$ orbitals on the carbonyl ligands.



3-27

The $x^{2}-y^{2}$ and $z^{2}$ orbitals have zero overlap with all the $\pi^{*}$ orbitals, as in the preceding complexes. These orbitals ( $\sigma$ antibonding) are therefore not perturbed by the $\pi$ interactions. It is only the nonbonding $x y$, $x z$, and $y z$ orbitals, derived from the $t_{2 g}$ block of the regular octahedron, that will be perturbed.

We consider first the $x y$ orbital (3-28). Its overlap with $\pi_{x}^{*(1)}$ is nonzero, but it is zero with $\pi_{z}^{*(1)}$, located in a nodal plane of $x y$. As a result, $x y$ is stabilized by a bonding interaction with $\pi_{x}^{*(1)}$. This analysis allows us to anticipate an equivalent interaction with the $\pi_{x}^{*(3)}$ orbital on the symmetry-related ligand $\mathrm{C}_{3} \mathrm{O}_{3}$. But the two orbitals $\pi_{x}^{*(2)}$ and $\pi_{y}^{*(2)}$ on the central ligand $\mathrm{C}_{2} \mathrm{O}_{2}$ are both located in the nodal planes of $x y$ ( $x z$ and $y z$, respectively), and their overlap with $x y$ is zero (3-28). We conclude that $x y$ is stabilized by two bonding interactions, with $\pi_{x}^{*(1)}$ and $\pi_{x}^{*(3)}{ }^{(3-31 a)}$.


The $y z$ orbital can interact with $\pi_{z}^{*(1)}$ on $\mathrm{C}_{1} \mathrm{O}_{1}$, but not with $\pi_{x}^{*(1)}$ which is located in one of its nodal planes (xy) (3-29). The same remarks may be made about the group $\mathrm{C}_{3} \mathrm{O}_{3}$. On the central carbonyl $\mathrm{C}_{2} \mathrm{O}_{2}$, $\pi_{y}^{*(2)}$ is the only orbital whose overlap with $y z$ is non-zero. The $y z$ orbital is therefore stabilized by three bonding interactions, with $\pi_{z}^{*(1)}, \pi_{y}^{*(2)}$, and $\pi_{z}^{*(3)}(\mathbf{3 - 3 1 b})$. Notice that the contributions of $\pi_{z}^{*(1)}$ and $\pi_{z}^{*(3)}$ are equal in magnitude, by symmetry, but the contribution of $\pi_{y}^{*(2)}$ on the central carbonyl $\mathrm{C}_{2} \mathrm{O}_{2}$ may be different, since it is not exchanged with the two other groups by any symmetry element in the complex.

$S=0$

$\boldsymbol{S} \neq \mathbf{0}$

$S=0$

$\boldsymbol{S} \neq \mathbf{0}$

3-29

The $x z$ orbital cannot interact with any of the orbitals on the $\mathrm{C}_{1} \mathrm{O}_{1}$ and $\mathrm{C}_{3} \mathrm{O}_{3}$ groups, and can combine only with the $\pi_{x}^{*(2)}$ orbital on the central group $\mathrm{C}_{2} \mathrm{O}_{2}(\mathbf{3 - 3 0})$. Therefore, the $x z$ orbital is stabilized by $a$ single bonding interaction, with $\pi_{x}^{*(2)}{ }_{(3-31 c)}$.


3-30

We see, therefore, that the three nonbonding orbitals derived from the $t_{2 g}$ block of the regular octahedron are perturbed in different ways in the substitution pattern that we have studied (3-31). All the $d \leftrightarrow \pi^{*}$ interactions concern the same lateral overlap between a pure $d$ orbital and a $\pi_{\mathrm{CO}}^{*}$ orbital, but the number of interactions changes from one orbital to another.


3-31a


3-31b


3-31c

As a consequence, the degeneracy of the three orbitals is lifted in this substitution pattern. We can anticipate that $y z$, stabilized by three bonding interactions with the $\pi_{\mathrm{CO}}^{*}$ orbitals, is lower in energy than $x y$ (two bonding interactions), which in turn is lower than $x z$ (only a single bonding interaction) (3-32). Part of the complete diagram for the interaction of these three $d$ orbitals and the $\pi_{\mathrm{CO}}^{*}$ orbitals has therefore been obtained. Moreover, it is the most important part, since it gives us the three new $d$ orbitals that may be occupied in this strong-field octahedral complex. It is straightforward to derive the missing parts of this diagram. Three antibonding combinations, mainly concentrated on the carbonyl groups, correspond to the three bonding combinations described in 3-31. We note that three linear combinations of the $\pi_{\mathrm{CO}}^{*}$ orbitals have been used in these interactions. Since initially there were six $\pi_{\text {CO }}^{*}$ orbitals, two per ligand, three other combinations have not interacted; they remain entirely localized on the ligands, at an energy level close to that of an isolated $\pi_{\text {CO }}^{*}$ orbital.


3-32

### 3.3.3.3. A rigorous or an approximate method?

The method we have used has consisted in analysing successively the possible interactions for each of the ligands with a $\pi$ system. This method, which can be applied in the same way to any type of complex with any substitution pattern (see Exercises 3.1-3.5), enables us to obtain rather quickly the shapes of the $d$ orbitals as perturbed by the $\pi$ interactions, as well as their relative energies. The use of group theory, which requires us in each case to construct the symmetry-adapted orbitals on the ligands, is more time-consuming. But is there really a fundamental difference between the two methods? The application of group theory is based on the symmetry elements of the complex as a whole; this is the rigorous way to proceed. However, the construction 'by hand' that has been proposed in the two preceding sections uses local symmetry elements shared by a $d$ orbital and a particular $\pi^{*}$ orbital, to deduce whether interaction is possible between them (Schemes 3-28-3-30). The two methods fuse when the symmetry elements of the complex and the local
symmetry elements used are identical. This is the case, for example, for the trans- $\left[\mathrm{ML}_{4} \mathrm{Cl}_{2}\right]$ and trans- $\left[\mathrm{ML}_{4}(\mathrm{CO})_{2}\right]$ complexes ( $\S$ 3.3.3.1) and for the $x y$ and $y z$ orbitals of the $m e r-\left[\mathrm{ML}_{3}(\mathrm{CO})_{3}\right]$ complex. However, in this latter case, we used a local symmetry element that is not a symmetry element of the complex. Use of the horizontal plane $x y$ allows us to show that the overlap between $\pi_{x}^{*(1)}$ and $x z$ is zero ( $\mathbf{3 - 3 0}$, left-hand side), and we concluded that $\pi_{x}^{*(1)}$ makes no contribution to the $x z$ orbital (3-31c). This conclusion is not rigorous, since the $x y$ plane is not a symmetry element for the complex as a whole. Fortunately, calculations show that the participation of $\pi_{x}^{*(1)}$ in this orbital is very small (as is that of $\pi_{x}^{*(3)}$, by symmetry), precisely because the local overlap between $\pi_{x}^{*(1)}$ and $x z$ is zero. The approximate description of the $d$ orbitals given in 3-31 is thus essentially correct.

### 3.3.4. $\left[\mathrm{MCl}_{6}\right]$ and $\left[\mathrm{M}(\mathrm{CO})_{6}\right]$ octahedral complexes

When the six double-face $\pi$-donor or $\pi$-acceptor ligands are identical, octahedral symmetry is preserved. We therefore find the usual splitting pattern in the d block, into two groups of degenerate orbitals, whose symmetries are $t_{2 g}$ and $e_{g}$. We shall now use the method set out in the preceding section to determine the shapes of the MO and the energetic consequences of the $\pi$ interactions.

### 3.3.4.1. The d block in $\left[\mathrm{MCl}_{6}\right]$ and $\left[\mathrm{M}(\mathrm{CO})_{6}\right]$ complexes

In the $\left[\mathrm{MCl}_{6}\right]$ complex, each orbital in the $t_{2 g}$ group ( $x y, x z$, and $y z$ ) can interact with four of the $p$ orbitals that characterize lone pairs on the ligands (3-33). All the interactions are antibonding ( $\pi$-donor ligands); they destabilize the three orbitals in an equivalent way, so these remain degenerate.


3-33

It is easy to verify that the orbitals in the $\mathrm{e}_{\mathrm{g}}$ group ( $z^{2}$ and $x^{2}-y^{2}$ ) have zero overlap with all the nonbonding $p$ orbitals on the ligands, so their energy and shape remain unchanged. The symmetry properties of the octahedral complexes $\left[\mathrm{MCl}_{6}\right]$ therefore allow only three linear combinations of the nonbonding $p$ orbitals on the chloride ligands, those shown in 3-33, to interact with the d-block orbitals. In other words,


Figure 3.8. Perturbation of the d block (in the centre, $\sigma$ interactions only) of an octahedral complex $\mathrm{ML}_{6}$ by the $\pi$ interactions with double-face $\pi$-donor ligands ( $\mathrm{L}=\mathrm{Cl}$, for example) on the left, and double-face $\pi$-acceptor ligands ( $\mathrm{L}=\mathrm{CO}$, for example) on the right.
only three lone pairs out of the twelve (two per ligand) can transfer electron density to the metal centre; we shall return to this point in §3.5, Figure 3.10.

By an exactly analogous analysis, it can be shown that in the case of six carbonyl ligands, the three orbitals in the $t_{2 g}$ block are stabilized to the same extent by four bonding interactions with the $\pi_{\mathrm{CO}}^{*}$ orbitals (3-34), though the two $e_{g}$ orbitals are not affected.


3-34

### 3.3.4.2. Strong-field and weak-field complexes

The energetic consequences of the $\pi$ interactions for the d-block orbitals in the $\left[\mathrm{MCl}_{6}\right]$ and $\left[\mathrm{M}(\mathrm{CO})_{6}\right]$ complexes are shown in Figure 3.8. If everything else is equal, the $\sigma$ interactions in particular, the presence of $\pi$-donor ligands decreases the energy separation ( $\Delta E$ ) between the $\mathrm{t}_{2 g}$ and $\mathrm{e}_{g}$ orbitals (weak field), whereas $\pi$-acceptor ligands increase this separation (strong field). This trend is consistent with the spectrochemical series presented below, which, from measurements of the $d-d$ transition energies ( $\mathrm{t}_{2 g} \rightarrow \mathrm{e}_{g}$ ), ranks ligands according to the strength of the field that they create (Chapter $2, \S 2.1 .2 .6$ ). It is clear that $\pi$-donor ligands, such as the halogens, are found at the beginning of the series (weak field), whereas $\pi$-acceptor ligands, such as carbonyl or cyanide, are at the end (strong field). For example, $\Delta E$ decreases from 26,600 to $15,060 \mathrm{~cm}^{-1}$ passing from $\left[\mathrm{Cr}(\mathrm{CN})_{6}\right]^{3-}$ (six $\pi$-acceptor ligands) to $\left[\mathrm{CrF}_{6}\right]^{3-}$ (six $\pi$-donor ligands), even though these two complexes have the same $d^{3}$ electronic configuration.

$$
\begin{align*}
\mathrm{I}^{-} & <\mathrm{Br}^{-}<\mathrm{Cl}^{-}<\mathrm{F}^{-}<\mathrm{OH}^{-}<\mathrm{O}^{2}<\mathrm{H}_{2} \mathrm{O}<\mathrm{NH}_{3} \\
& <\mathrm{NO}_{2}^{-}<\mathrm{CH}_{3}^{-}<\mathrm{C}_{6} \mathrm{H}_{5}^{-}<\mathrm{CN}^{-}<\mathrm{PR}_{3}<\mathrm{CO} \tag{3.1}
\end{align*}
$$

Notice, however, that the $\pi$ interactions are not the only factor which influences the $t_{2 g}-e_{g}$ energy separation. We must remember that the $\sigma$ interactions create the primary $d$-orbital splittings, and that as these interactions become stronger (phosphine ligands $\mathrm{PR}_{3}$, metal-carbon bonds in organometallic complexes, for example), the antibonding $\mathrm{e}_{g}$ levels are pushed to ever higher energy.

[^3]

3-36

## 3.4. $\pi$ complexes: the example of ethylene

In all of the preceding examples, the ligands that possess a $\pi$ system are bound to the metal centre by a single atom (coordination mode $\eta^{1}$ ). The situation is different when several ligand atoms are bound in an equivalent way to the metal centre. In particular, this is the case for $\pi$ complexes (Chapter 1, § 1.1.1.3), in which the $\pi$ system of the ligand 'points' in the direction of the metal, rather than being perpendicular to the metal-ligand bond. The description of the metal-ligand bonds now requires us to take all the $\pi$ orbitals of the ligand into account, both occupied and empty, as well as the orbitals of appropriate symmetry on the metal. As an example, we shall first consider the simplest complex from this family, that in which a molecule of ethylene is bound to the metal in the $\eta^{2}$ mode.

### 3.4.1. Orbital interactions: the Dewar-Chatt-Duncanson model ${ }^{8}$

### 3.4.1.1. $\sigma$ or $\pi$ interactions?

We consider first the $\pi$-bonding orbital on ethylene. It can interact with a metal orbital, for example, $z^{2}(\mathbf{3 - 3 5})$. The resulting overlap is intermediate between an axial overlap, normally associated with a $\sigma$ interaction, and a lateral overlap which corresponds to a $\pi$ interaction. Notice that there is no nodal surface that is common to the two orbitals, and that their overlap is not changed by a rotation of the ligand around the $z$ axis. As a consequence, even though the ligand orbital is of $\pi$ type, the metal-ligand interaction in which it is involved has the characteristics of an ordinary $\sigma$ interaction (see 3-17a, for example). The $\pi^{*}$ orbital can also interact with a metal $d$ orbital, for example, $x z(3-36)$. As in the 'traditional' $\pi$ interactions described in preceding sections, the two orbitals have a common nodal plane $(y z)$, and a rotation of the ligand around the $z$-axis decreases their overlap, to the extent that it is completely eliminated for a rotation of $90^{\circ}$. The interaction that involves the $\pi^{*}$ orbital can therefore be described as a $\pi$ interaction.

### 3.4.1.2. The Dewar-Chatt-Duncanson model

Ethylene behaves as an L-type ligand thanks to its doubly occupied $\pi$ orbital. It can transfer electron density to the metal (a donation interaction) by interaction with an empty orbital on the metal centre, whose symmetry is suitable ( $z^{2}$, for example, 3-37a). This is a stabilizing twoelectron interaction, which stabilizes the $\pi$ level of the ligand. Due to the relative energies of the two initial orbitals, the occupied MO is mainly concentrated on the ligand (a bonding orbital).



There is a second interaction that involves the empty antibonding $\pi^{*}$ orbital and the metal $d$ orbital with the same symmetry $(x z)$, which is lower in energy than the $\pi^{*}$ orbital ( $\mathbf{3 - 3 7 b}$ ). If this latter orbital is doubly occupied, this interaction is stabilizing, and it leads to a transfer of electron density from the metal to the ligand. This is therefore a backdonation interaction, where ethylene plays the role of a $\pi$ acceptor, using its empty $\pi^{*}$ orbital. The doubly occupied orbital, mainly concentrated on the metal, is part of the $d$ block of the complex; it can be described as a metal $d$ orbital that is stabilized by a bonding interaction with the $\pi^{*}$ orbital on ethylene.


These two stabilizing interactions taken together constitute the Dewar-Chatt-Duncanson model of the bond between an olefin and a metal centre.

### 3.4.2. Electronic structure of a $d^{6}$ complex $\left[\mathrm{ML}_{5}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$

### 3.4.2.1. The complete interaction diagram

Consider a pseudo-octahedral $d^{6}$ complex of the type [ML5 $\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)$ ] (3-38), where the five L ligands are supposed to have only $\sigma$ interactions with the metal centre. If this complex is decomposed into a $d^{6}$ fragment $\mathrm{ML}_{5}$, with a square-base pyramidal (SBP) geometry, and an ethylene fragment, the interaction between the orbitals on the two fragments enables us to analyse the electronic factors that are at the origin of the ethylene-metal bond.

On the metallic fragment $\mathrm{ML}_{5}$, we consider the four low-energy $d$ orbitals in this type of structure (Chapter 2, § 2.3.1), together with the $\pi$ and $\pi^{*}$ orbitals on the ethylene ligand. The $\pi$ orbital is doubly occupied, as are the three strictly nonbonding $d$ orbitals of the $d^{6}$ fragment $\mathrm{ML}_{5}$.

Figure 3.9. Diagram for the interaction between the $d$ orbitals of a $d^{6} \mathrm{ML}_{5}$ fragment and the $\pi$ and $\pi^{*}$ orbitals on an ethylene ligand (the Dewar-Chatt-Duncanson model).


Given the orbitals' symmetry properties with respect to the planes $P_{1}$ and $P_{2}$ that are defined in 3-38, we obtain the interaction diagram shown in Figure 3.9. The $\pi$ orbital (SS) is stabilized by a bonding interaction with the polarized $z^{2}$ orbital ( $\mathbf{S S}$ ) on the metallic fragment. The doubly occupied $y z$ orbital (SA) is stabilized by a bonding interaction with the $\pi^{*}$ orbital (SA). The two other $d$ orbitals ( $x y(\mathbf{A A})$ and $x z(\mathbf{A S})$ ) are not affected by the interaction. ${ }^{9}$ The two stabilizing interactions, donation ( $\pi \rightarrow z^{2}$ ) and back-donation ( $y z \rightarrow \pi^{*}$ ), are thus indeed those that were described in a general way in the preceding section.

The MO of the complex that are shown in Figure 3.9 can therefore be described in the following ways: (i) the SS bonding MO is a bond orbital; (ii) the SA weakly bonding MO and the AS and AA nonbonding MO are the three orbitals that are derived from the $\mathrm{t}_{2 g}$ block of a regular octahedron. If they are doubly occupied, the electronic configuration is $d^{6}$; (iii) the SA antibonding MO is essentially antibonding on the ethylene ligand; (iv) lastly, the SS antibonding MO, largely composed of the $z^{2}$ orbital, is mainly antibonding for the ethylene ligand and the trans ligand L . In fact, it is one of the orbitals that are derived from the antibonding $\mathrm{e}_{g}$ block of a regular octahedron. The other orbital, $x^{2}-y^{2}$, does not appear on this diagram since we did not consider it on the initial ML ${ }_{5}$ fragment, as its energy is too high.

### 3.4.2.2. A molecular ethylene complex or a metallacyclopropane?

The donation interaction ( $\pi \rightarrow z^{2}$ ) reduces the electron density in the $\pi$-bonding orbital of the ethylene ligand. The carbon-carbon

${ }^{10}$ There is, however, a difference in the relative magnitudes of the coefficients on the metal and on the ligand in the orbital whose symmetry is written SA. In the molecular ethylene complex, this orbital is concentrated mostly on the metal, whereas in the metallacyclopropane, it is mainly on the ligand. We shall return to this point in $\oint 4.5$ of Chapter 4, which is devoted to the mechanism of oxidative addition.
bond is therefore weakened (3-39). Now the back-donation interaction $\left(y z \rightarrow \pi^{*}\right)$ transfers electron density into the $\pi^{*}$ antibonding orbital, which also weakens this bond (3-40). However, these two interactions are both bonding between the carbon atoms and the metal centre, so they contribute to the formation of metal-carbon bonds.

The question is thus open: is it appropriate to speak of an ethylene complex, in which the ligand is clearly identified as molecular ethylene (3-41a), or is it better to describe it as a metallacyclopropane, with a $\mathrm{C}-\mathrm{C}$ single bond and two $\mathrm{M}-\mathrm{C}$ bonds (3-41b)? These are in fact two mesomeric forms of the same complex; the metal oxidation state, in particular, is different in these two forms, being higher by two in the metallacycle. The change from one limiting form to the other represents the oxidative addition reaction of ethylene, which leads to the metallacyclopropane. The real complex is usually intermediate between these two limiting forms. Typically, the carbon-carbon bond length is around $1.43 \AA$ (instead of 1.32 or $1.54 \AA$ for a standard double or single bond, respectively), and the ethylene ligand becomes non-planar: the $\mathrm{CH}_{2}$ groups become pyramidal, in the direction opposite to the metal centre, though not as extensively as in a real three-membered ring.

From the electronic structure point of view, there are similarities between the two $\mathbf{S S}$ and $\mathbf{S A}$ bonding molecular orbitals of the $\pi$ complex (3-42 and 3-43, left-hand side) and the corresponding orbitals of the metallacyclopropane. The latter can be represented schematically by considering the in-phase and out-of-phase combinations of two localized orbitals, each of which characterizes a $\sigma_{\mathrm{MC}}$ bond (3-42 and 3-43, righthand side). ${ }^{10}$ The first combination corresponds to the orbital $\pi+z^{2}$ of the molecular ethylene complex, the second to the orbital $y z+\pi^{*}$.


3-42

SA



Figure 3.10. The $\pi$ system of the cyclopentadienyl ligand $\left(\mathrm{Cp}^{-}\right)$.

Similarly, there is a resemblance between these orbitals and two of the three orbitals that characterize the $\mathrm{C}-\mathrm{C}$ bonds in cyclopropane itself (3-44, symmetry SS and 3-45, symmetry SA).

In fact, depending on the nature of the metallic fragment $\mathrm{ML}_{n}$ which interacts with the ethylene ligand and on the substituents on this ligand, a whole range of interactions can be observed, ranging from weak interactions, which correspond to a complex of molecular ethylene, to strong interactions which lead to a structure close to a metallacyclopropane. So the best answer to the general question posed at the beginning of this section 'a $\pi$ complex of molecular ethylene or a metallacyclopropane?' seems in fact to be . . . both!

### 3.4.3. Metallocenes $\mathrm{Cp}_{2} \mathrm{M}$

### 3.4.3.1. The $\pi$ system of the cyclopentadienyl ligand

The shapes and the relative energies of the five $\pi$ orbitals in cyclopentadiene $(\mathrm{Cp})$ are presented in Figure 3.10, where the electronic occupation corresponds to the anion $\mathrm{Cp}^{-}$(the ionic model). A symbolic representation of the orbitals, which shows only the nodal positions between the carbon atoms, is given on the right of the figure. As a consequence of the three occupied $\pi \mathrm{MO}, \mathrm{Cp}^{-}$can be characterized as an $\mathrm{L}_{3}$ ligand.

The donation and back-donation interactions of the Dewar-ChattDuncanson model can occur, thanks to the three occupied and the two empty MO, respectively.

### 3.4.3.2. $\mathrm{Cp}_{2} \mathrm{M}$ complexes: the ferrocene example

The three occupied MO on each $\mathrm{Cp}^{-}$ligand combine in-phase $\left(\pi^{(+)}\right)$ and out-of-phase $\left(\pi^{(-)}\right)$, leading to the formation of six occupied MO that are delocalized onto the two ligands. It is easy to verify that $\pi_{1}^{(+)}$




Figure 3.11. The d block of a $\mathrm{Cp}_{2} \mathrm{M}$ complex, where the electronic occupation $\left(d^{6}\right)$ corresponds, for example, to ferrocene $\left[\mathrm{Cp}_{2} \mathrm{Fe}\right]$.
has the proper symmetry to interact with the $s$ orbital on the metallic centre (3-46a). Similarly, $\pi_{1}^{(-)}$can interact with $p_{z}(\mathbf{3 - 4 6 b}), \pi_{2}^{(+)}$with $p_{y}(\mathbf{3 - 4 6 c}), \pi_{3}^{(+)}$with $p_{x}(\mathbf{3 - 4 6 d}), \pi_{2}^{(-)}$with $y z(\mathbf{3 - 4 6 e})$ and $\pi_{3}^{(-)}$with $x z(\mathbf{3 - 4 6}(\mathbf{f}))$. In this way, six occupied bonding MO are formed that characterize six bonds, as in an octahedral complex.


The participation of two $d$ orbitals in the MO that describe the bonds leads to the presence of two antibonding orbitals in the d block ( $y z$ and $x z$ ), while the other $d$ orbitals ( $z^{2}, x^{2}-y^{2}$, and $x y$ ) make up a block of three nonbonding or nearly nonbonding orbitals. In fact, $x^{2}-y^{2}$ and $x y$ are stabilized by bonding interactions with the $\pi^{*}$ orbitals of appropriate symmetry on the Cp rings (Figure 3.11).

The d block of a $\mathrm{Cp}_{2} \mathrm{M}$ complex therefore has the same characteristics as that of an octahedral complex. In ferrocene, $\left[\mathrm{Cp}_{2} \mathrm{Fe}\right]$, whose electronic configuration is $\left(d^{6}\right)$, the three nonbonding $d$ orbitals are occupied and this complex can be described as pseudo-octahedral, with 18 electrons.

### 3.4.4. $\mathrm{Cp}_{2} \mathrm{ML}_{n}$ complexes

There are many complexes of the type $\left[\mathrm{Cp}_{2} \mathrm{ML}_{n}\right](n=1,3)$, where M is a transition metal towards the left of the periodic classification ( $\mathrm{M}=\mathrm{Ti}, \mathrm{V}, \mathrm{Zr}, \mathrm{Hf}, \mathrm{Mo}$, for example). In these complexes, the $\mathrm{Cp}_{2} \mathrm{M}$ fragment is bent, rather than linear as in the metallocenes. The other ligands are all located in the plane that is perpendicular to that defined by the M atom and the centres of the two Cp rings (3-47 and 3-48). If $n=3$, this arrangement imposes very small $L-M-L$ angles (about


Figure 3.12. Energy changes for the three lowest-energy orbitals of the d block for a $\mathrm{Cp}_{2} \mathrm{M}$ complex passing from a linear to a bent arrangement of the Cp ligands.
$60^{\circ}$ in $\left.\mathrm{Cp}_{2} \mathrm{Nb}(\mathrm{H})_{3}, 3-47\right)$. Since each Cp ligand forms three bonds with the metal, the coordination number of these complexes is seen to be between $7(n=1)$ and $9(n=3)$, which is quite rare for the metals of the first three transition series.

### 3.4.4.1. The bent $\mathrm{Cp}_{2} \mathrm{M}$ fragment

Consider the three low-energy $d$ orbitals previously established for a linear $\mathrm{Cp}_{2} \mathrm{M}$ complex (Figure 3.11, the $z^{2}, x^{2}-y^{2}$, and $x y$ orbitals). In the $C_{2 v}$ point group, these orbitals have $\mathrm{a}_{1}\left(z^{2}\right.$ and $\left.x^{2}-y^{2}\right)$ and $\mathrm{b}_{2}(x y)$ symmetries. When the $\mathrm{Cp}_{2} \mathrm{M}$ fragment is bent (Figure 3.12), the $\mathrm{b}_{2}$ orbital is destabilized, since the stabilizing interactions with the $\pi^{*}$ orbitals on the Cp rings decrease and the repulsive interactions between the $\mathrm{M}-\mathrm{Cp}$ bonds increase. The same behaviour could be expected for the $1 \mathrm{a}_{1}$ orbital $\left(x^{2}-y^{2}\right)$, but it mixes with the $2 \mathrm{a}_{1}$ orbital $\left(z^{2}\right)$, with the result that its energy stays roughly constant, whereas the $2 \mathrm{a}_{1}$ orbital is strongly destabilized (an interaction between two orbitals of the same symmetry). This mixing between $z^{2}$ and $x^{2}-y^{2}$ which accompanies the bending of the fragment leads to a polarization of $x^{2}-y^{2}$ along the $x$-axis, and of $z^{2}$ in the $x y$ plane, as shown to the right of Figure 3.12.

It will be useful to consider another representation of these orbitals, which shows their amplitudes and nodal properties in the symmetry plane $x y$ which interchanges the two Cp ligands (3-49). The other ligands $(\mathrm{L})$ in $\left[\mathrm{Cp}_{2} \mathrm{ML}_{n}\right]$ complexes lie in this plane.


### 3.4.4.2. $\mathrm{Cp}_{2} \mathrm{ML}_{n}$ complexes $(n=2,3)$

We consider first the complex $\left[\mathrm{Cp}_{2} \mathrm{MoCl}_{2}\right]$ (3-48). It can be decomposed into a metallic fragment, $\left[\mathrm{Cp}_{2} \mathrm{Mo}^{2+}\right.$, and two chloride ligands, $\mathrm{Cl}^{-}$(the ionic model). The metallic fragment has a $d^{2}$ electronic configuration: the $1 \mathrm{a}_{1}$ orbital described above is occupied and the two other $d$ orbitals ( $\mathrm{b}_{2}$ and $2 \mathrm{a}_{1}$ ) are empty (Figure 3.13, left-hand side). The two $\mathrm{Cl}^{-}$ligands


Figure 3.13. Construction of the MO of the complex $\left[\mathrm{Cp}_{2} \mathrm{MoCl}_{2}\right]$ from those on the $\left[\mathrm{Cp}_{2} \mathrm{Mo}\right]^{2+}$ and $2 \mathrm{Cl}^{-}$fragments.
supply four electrons that occupy two MO, which are in-phase ( $\mathrm{a}_{1}$ ) and out-of-phase ( $\mathrm{b}_{2}$ ) combinations of two $p$ orbitals (Figure 3.13, righthand side). The formation of the MO of the complex from those on the fragments therefore involves an interaction between $\mathrm{b}_{2}$ and $\mathrm{a}_{1}$ pairs of orbitals. For the latter, the dominant interaction involves the $2 \mathrm{a}_{1}$ orbital on the metallic fragment. In fact, the $\mathrm{a}_{1}$ orbital on the chlorides is concentrated in the nodal planes of the $1 \mathrm{a}_{1}$ orbital on the metallic fragment, so the overlap between these two orbitals is very weak.

There is thus only one approximately nonbonding orbital in the d block, and it is doubly occupied. With a $d^{2}$ electronic configuration, and taking account of the eight metal-ligand bonds, $\left[\mathrm{Cp}_{2} \mathrm{MoCl}_{2}\right]$ is therefore an 18 -electron complex. 16 -electron $d^{0}$ complexes are also known in this family ( $n=2$ ), such as $\left[\mathrm{Cp}_{2} \mathrm{ZrCl}_{2}\right]$; the nonbonding $d$ orbital is empty in them.

We now turn to a complex from the $\mathrm{Cp}_{2} \mathrm{ML}_{3}$ family, such as $\mathrm{Cp}_{2} \mathrm{NbH}_{3}(3-47)$. This can be decomposed into a metallic fragment, $\mathrm{Cp}_{2} \mathrm{Nb}^{3+}$, whose electronic configuration is $d^{0}$, and a fragment that groups together the three hydrides, $\mathrm{H}_{3}^{3-}$ (the ionic model), with six electrons in three orbitals. The symmetries of these three orbitals are well adapted to interact with the three $d$ orbitals of the metallic fragment $(\mathbf{3 - 5 0})$. In particular, there is a large overlap between the $1 \mathrm{a}_{1}$ orbital on the metallic fragment and the $2 a_{1}$ orbital of $\mathrm{H}_{3}^{3-}$.




After the interaction, three doubly occupied bonding MO are formed that are mainly concentrated on the hydrides (3-51a). They characterize the three $\mathrm{Nb}-\mathrm{H}$ bonds. The three antibonding combinations, $(\mathbf{3}-\mathbf{5 1 b})$ which belong to the d block, are empty.


There are therefore no nonbonding $d$ orbitals in this complex. With its $d^{0}$ electronic configuration and nine bonds around the metal, $\left[\mathrm{Cp}_{2} \mathrm{NbH}_{3}\right]$ is an 18-electron complex.

## 3.5. $\pi$ interactions and electron counting

The formal electron count, as it is usually performed for transition metal complexes, (Chapter 1, § 1.1), takes account of the electrons involved in the $\sigma$ interactions and the $n$ nonbonding electrons on the metal which do not participate in these interactions ( $d^{n}$ electronic configuration). But it ignores the consequences of the $\pi$ interactions, even though these are accompanied by transfers of electron density involving some d-block orbitals.

As a first illustration of this point, consider $\left[\mathrm{W}\left(\mathrm{PR}_{3}\right)_{6}\right]$ and [W $\left.(\mathrm{CO})_{6}\right]$; both are described as octahedral complexes whose electronic configuration is $d^{6}$, with six electrons in the three $t_{2 g}$ orbitals of the d block of the octahedron. A phosphine ligand, $\mathrm{PR}_{3}$, may, in a first approximation, be considered as a simple $\sigma$-donor. ${ }^{11}$ From this viewpoint, the occupied $\mathrm{t}_{2 g}$ orbitals of the $\left[\mathrm{W}\left(\mathrm{PR}_{3}\right)_{6}\right]$ complex are therefore pure atomic $d$ orbitals, entirely localized on the metal. However, since the carbonyl ligand is a $\pi$ acceptor, the $t_{2 g}$ orbitals of the [W(CO) 6 ] complex are partially delocalized on to the ligands (3-34), leading to a substantial transfer of electron density from the metal to the ligands. Although both formally have a $d^{6}$ electronic configuration, the shape of
${ }^{12}$ We noted, in Chapter 2, that complexes with a $d^{0}$ electronic configuration can exist if the ligands possess lone pairs (see $\S$ 2.1.3.2, 2.3.3.2, and 2.4.2.2)

Figure 3.14. Interaction diagram showing the perturbation of the $d$ block for an octahedral $\left[\mathrm{M}(\mathrm{Cl})_{6}\right]$ complex ( $\sigma$ interactions only, left-hand side) by the symmetry-adapted $\pi$ orbitals ( $\mathrm{t}_{2 g}$ ) on the six double-face $\pi$-donor ligands. The electronic occupation shown corresponds to a complex with a $d^{0}$ electronic configuration $\left(\left[\mathrm{W}(\mathrm{Cl})_{6}\right]\right.$, for example).
the orbitals reflects the fact that the electron density at the metal centre is higher in the $\left[\mathrm{W}\left(\mathrm{PR}_{3}\right)_{6}\right]$ complex than in $\left[\mathrm{W}(\mathrm{CO})_{6}\right]$.

Complexes with $\pi$-donor ligands that can transfer electron density to the metal centre are even more interesting. Consider the octahedral complex $\left[\mathrm{W}(\mathrm{Cl})_{6}\right]$, for example: according to the standard counting procedure, its electronic configuration is $d^{0}[\mathrm{~W}(\mathrm{VI})]$, and therefore it is a 12 -electron complex. But this description is substantially modified if we take account of the $\pi$ interactions. We have already seen that the three orbitals of the $t_{2 g}$ block are destabilized by antibonding interactions with the $p$ lone pairs on the chloride ligands (3-33). In this complex, therefore, all the orbitals of the $d$ block are antibonding, those destabilized by the $\sigma$ interactions (the $\mathrm{e}_{g}$ block) being higher in energy than those destabilized by the $\pi$ interactions (Figure 3.14). This antibonding character of all the d-block orbitals enables us to understand the stability of this $d^{0}$ complex rather more easily. ${ }^{12}$ Moreover, though the $\pi$ interactions destabilize three orbitals in the d block, they also stabilize the three orbitals mainly concentrated on the lone pairs of the chloride ligands (Figure 3.14). The $\mathrm{t}_{2 g}$ orbitals of the metal therefore contribute partially to the occupied MO. As a consequence, though it is accurate, following the normal definition of the $d$ block, to describe $\left[W(C l)_{6}\right]$ as a $d^{0}$ complex, the $t_{2 g}$ orbitals of the metal are in fact partially occupied, and the metal is much less poor in electrons than it appears at first sight.

In view of this analysis, it might seem legitimate for the $\pi$ electrons that are transferred, at least partially, to the $d$ orbitals of the metal centre to be included in the electron count. The complex $\left[\mathrm{W}(\mathrm{Cl})_{6}\right]$, previously described as a 12 -electron complex ( $\sigma$ electrons only), thus becomes an... 18 -electron complex if one takes account of the six $\pi$ electrons

${ }^{13}$ To understand this stability another way, notice that as a result of the $\pi$ interactions, there is a fairly large energy gap between the highest occupied ( $\pi$ bonding) and the lowest unoccupied MO ( $\pi^{*}$ antibonding).
'transferred' to the metal (Figure 3.14). It is not a coincidence if the ' 18 -electron rule' is now verified with this new counting procedure: once the $\pi$ interactions are considered, there are no empty nonbonding MO on the metal. To attain the normal count of 18 electrons around the metal, it is therefore sufficient for all the bonding MO ( $\sigma$ and $\pi$ ) to be doubly occupied; this electron count allows us to understand the stability of a complex which at first sight is very electron-deficient. ${ }^{13}$ The stability of the complex $\left[\mathrm{W}(\mathrm{Cl})_{6}\right]$ in an octahedral geometry contrasts strikingly with the instability of $\left[W(\mathrm{Me})_{6}\right]$ in this same geometry. The latter is also a $d^{0}$ complex, but one for which the $\pi$ interactions are negligible. It therefore possesses three MO based on $d$ orbitals that are nonbonding and unoccupied, and it can be described as a 'true' 12-electron complex. In fact, it adopts a different geometry, which is trigonal-biprismatic; we may conclude that the octahedral geometry of $\left[\mathrm{W}(\mathrm{Cl})_{6}\right]$ is controlled by $\pi$ interactions.

This analysis can be extended to other ligand fields, and is not limited to complexes whose d block is completely empty. In the presence of $\pi$-donor ligands, the apparent electron deficiency is compensated by interactions with the lone pairs on the ligands. Consider tetrahedral complexes such as $\left[\mathrm{Ti}(\mathrm{Cl})_{4}\right]$ or $\left[\mathrm{Zr}(\mathrm{Cl})_{4}\right]$, which are $d^{0}$ complexes of $\mathrm{Ti}(\mathrm{IV})$ or $\operatorname{Zr}(\mathrm{IV})$, and therefore complexes with only eight electrons! However, it can readily be shown that five of the eight orbitals that characterize the lone pairs on the four chloride ligands are able, by symmetry, to interact with the d-block orbitals (Chapter 2, § 2.4.1): two have e symmetry and three have $t_{2}$ symmetry. In this way, $10 \pi$ electrons are 'transferred' to the metal, so the complexes may be considered to have 18 electrons.

## Exercises

## $\pi$ interactions in an $M L_{4}$ complex

## 3.1

1. Indicate the shapes and the relative energies of the four lowest-energy d-block orbitals for a square-planar $\mathrm{ML}_{4}$ complex ( $\sigma$ interactions only).
2. How are these shapes and energies modified in the complexes?
(a) $\left[\mathrm{PtCl}_{4}\right]^{2-}$
(b) $\left[\mathrm{Ni}(\mathrm{CN})_{4}\right]^{2-}$
3. Repeat the question for Wilkinson's catalyst, $\left[\mathrm{Rh}\left(\mathrm{PPh}_{3}\right)_{3} \mathrm{Cl}\right]$; for simplicity, the triphenylphosphine ligands may be considered to be pure $\sigma$ donors.

## $\pi$ interactions in an $\mathrm{ML}_{5}$ complex (TBP)

3.2

1. Indicate the shapes and the relative energies of the four lowestenergy d-block orbitals for an $\mathrm{ML}_{5}$ complex in a TBP geometry ( $\sigma$ interactions only).
2. How are the shapes and energies of these orbitals modified if a carbonyl group is in an axial position (1a), or an equatorial position (1b)? (for the first structure, show that certain overlaps that are non-zero by symmetry are in fact very small, so they may be neglected).

1a

1b

2a

2b
3. Deduce the more favourable substitution site for $\pi$ interactions in diamagnetic $d^{4}$ and $d^{8}$ complexes.
4. Repeat question 2 for complexes substituted by a chloride ligand in an axial (2a) or equatorial position (2b).
5. Deduce the more favourable substitution site for $\pi$ interactions in diamagnetic $d^{4}$ and $d^{8}$ complexes $2 \mathbf{a}$ and $\mathbf{2 b}$.

## $\pi$ interactions in an octahedral $\mathrm{ML}_{6}$ complex

## 3.3

1. By writing a Lewis structure, show that $\mathrm{C} \equiv \mathrm{N}$ is an X -type ligand.
2. Assuming that the order of increasing energy for the MO is the same in CN as in $\mathrm{C} \equiv \mathrm{O}$,
(a) indicate the nature of the MO that has a $\sigma$ interaction with a metal centre;
(b) give the shape and the occupation of the $\pi$-type MO ;
(c) indicate whether this ligand is a $\pi$ acceptor or a $\pi$ donor;
(d) indicate whether the CN ligand binds to the metal through carbon or nitrogen in the $\eta^{1}$ coordination mode.
3. Give the shapes and relative energies of the three lowest-energy $d$ orbitals in an octahedral complex of the type $\left[\mathrm{M}(\mathrm{CN})_{6}\right]$.
4. In which of the two complexes $\left[\mathrm{Fe}(\mathrm{CN})_{6}\right]^{4-}$ and $\left[\mathrm{Fe}(\mathrm{CN})_{6}\right]^{3-}$ are the $\mathrm{Fe}-\mathrm{CN}$ bonds shorter?
5. Is this result predictable just from the oxidation state of the metal in the two complexes?

## 3.4

Consider an octahedral complex $\left[\mathrm{ML}_{4} \mathrm{Cl}_{2}\right]$ in which the Cl ligands are in cis positions: the $x$-axis bisects the $\mathrm{Cl}-\mathrm{M}-\mathrm{Cl}$ angle, as shown below.



1. Identify the two planes of symmetry in the complex $\left(P_{1}\right.$ and $\left.P_{2}\right)$.
2. Give the shapes of the three nonbonding orbitals for a regular octahedral complex in which the ligands only have $\sigma$ interactions with the metal. Indicate their symmetry properties with respect to $P_{1}$ and $P_{2}$.
3. Construct the four symmetry-adapted orbitals that characterize the $\pi$-type lone pairs on the Cl ligands, and give their symmetries with respect to $P_{1}$ and $P_{2}$.
4. How many lone pairs on the ligands are involved in $\pi$ interactions?
5. Place these seven orbitals (three $d$ orbitals on the metal and four lone-pair orbitals) on an energy-level diagram, and justify the relative positions:
(a) of the ligand and metal orbitals;
(b) of the four different ligand orbitals.
6. Construct the interaction diagram, and indicate which of the resulting MO belong to the d block.
7. How many electrons are there around the metal according to the traditional count in the case of a $d^{0}$ complex?
8. How many electrons are there once the $\pi$ transfers are taken into account?

## Pentagonal bipyramidal $\mathrm{ML}_{7}$ complex

3.5

Among the different geometries that can be adopted by an $\mathrm{ML}_{7}$ complex, the pentagonal-bipyramid (PBP), with two axial ligands ( $\mathrm{L}_{1}$ and $\mathrm{L}_{7}$, on the $z$-axis) and five equatorial ligands $\left(\mathrm{L}_{2}-\mathrm{L}_{6}\right.$, in the $x y$ plane), is one of the most common.


1. How many of the $d$ orbitals are strictly nonbonding? Give the shapes of these orbitals.
2. Consider PBP complexes with a $d^{4}$ electronic configuration in which there is one carbonyl ligand, in either an axial $\left(\mathrm{L}_{1}\right)$ or an equatorial position (choose the position $\mathrm{L}_{2}$ ).
(a) analyse the $\pi$-type interactions in each isomer;
(b) deduce the favoured site for a carbonyl ligand.

## The use of group theory

3.6

Solve Exercise 3.1 (question 2) by making use of the results of Exercise 6.13 (Chapter 6) on the symmetry-adapted orbitals for $\mathrm{ML}_{4}$ complexes with a $\pi$ system on the ligands, and the character table for the $\mathrm{D}_{4 \mathrm{~h}}$ point group (Table 6.18).

## 3.7

1. Give the shapes and relative energies of the five d-block orbitals for a trigonal-planar $\mathrm{ML}_{3}$ complex ( $\sigma$ interactions only).
2. Consider the complex $\left[\mathrm{Ni}(\mathrm{CO})_{3}\right]$ in this geometry. By making use of the results established in Chapter 6 (§ 6.6.6.2), give the symmetry-adapted linear combinations of the $\pi_{\mathrm{CO}}^{*}$ orbitals (the point group is $D_{3 h}$, whose character table is given in Table 6.21).
3. Construct the diagram for the interaction of these orbitals with those in the d block of question 1 . Deduce the shapes and relative energies of the d-block MO in $\left[\mathrm{Ni}(\mathrm{CO})_{3}\right]$.
4. Try to find these orbitals without using group theory (construction 'by hand', § 3.3). Show that some of the information provided by group theory cannot be readily obtained without its use.

## Appendix C: the carbonyl ligand, a double-face $\pi$-acceptor

In an octahedral complex with a carbonyl ligand (see § 3.2.2), two metal $d$ orbitals ( $x z$ and $y z, 3 \mathrm{C} .1$ ) can interact with the $\pi$-bonding and $\pi^{*}$ antibonding MO on the CO ligand: $\pi_{x}$ and $\pi_{x}^{*}$ with $x y, \pi_{z}$ and $\pi_{z}^{*}$ with $y z$, following two equivalent three-orbital interaction diagrams (see Appendix A, Chapter 2). By comparing the strength of the $d \leftrightarrow \pi$ and $d \leftrightarrow \pi^{*}$ interactions, we may deduce whether the carbonyl ligand is a $\pi$ donor or a $\pi$ acceptor. These interactions, which involve orbitals of different energy, are proportional to $S^{2} / \Delta \varepsilon$ (Chapter $1, \S$ 1.3.2); they therefore increase in strength with a decrease in the orbital energy difference, but with an increase in the overlap.

Table 3C.1. $d$-orbital energies for the metals of the first transition series, and the energy differences $d-\pi_{\mathrm{CO}}\left(\Delta \varepsilon_{\pi}\right)$ and $\pi_{\mathrm{CO}}^{*}-d\left(\Delta \varepsilon_{\pi^{*}}\right)$ calculated by the extended Hückel method (eV)

|  | $C c$ | $T i$ | $V$ | $C r$ | $M n$ | Fe | $C o$ | $N i$ | $C u$ |
| :--- | ---: | ---: | ---: | ---: | :---: | ---: | ---: | ---: | ---: |
| $\varepsilon_{d}$ | -8.5 | -10.8 | -11.0 | -11.2 | -11.7 | -12.6 | -13.2 | -13.5 | -14.0 |
| $\Delta \varepsilon_{\pi}$ | 7.2 | 4.9 | 4.7 | 4.5 | 4.0 | 3.1 | 2.5 | 2.2 | 1.7 |
| $\Delta \varepsilon_{\pi^{*}}$ | -0.6 | 1.7 | 1.9 | 2.1 | 2.6 | 3.5 | 4.1 | 4.4 | 4.9 |


3.C1

The comparison of $\Delta \varepsilon_{\pi}$ with $\Delta \varepsilon_{\pi^{*}}$ is not straightforward, since the $d$-orbital energy depends on the nature of the metal: these metal orbitals become considerably more stable as one moves from left to right in the periodic classification (Chapter 1, Table 1.3). One can obtain at least a qualitative indication of the relative values of $\Delta \varepsilon_{\pi}$ and $\Delta \varepsilon_{\pi^{*}}$ by using the extended Hückel method, ${ }^{14}$ for example, for the metals of the first transition series (Table 3C.1). The energy of the $d$ orbitals lies between $-8.5 \mathrm{eV}(\mathrm{Sc})$ and $-14.0 \mathrm{eV}(\mathrm{Cu}),{ }^{15}$ and the energies calculated for the $\pi_{\mathrm{CO}}$ and $\pi_{\mathrm{CO}}^{*} \mathrm{MO}$ are -15.7 and -9.1 eV , respectively, for a $\mathrm{C}-\mathrm{O}$ distance of $1.14 \AA$ (the average value in metal carbonyl complexes).

If we consider only the energetic criterion, the interaction with the $\pi_{\mathrm{CO}}^{*}$ orbitals is stronger than with $\pi_{\mathrm{CO}}$ for all the metals on the left of the table, up to and including manganese $\left(\Delta \varepsilon_{\pi^{*}}<\Delta \varepsilon_{\pi}\right)$, suggesting that the carbonyl ligand should behave as a $\pi$ acceptor. On the other hand, for the metals furthest to the right, and particularly for $\mathrm{Co}, \mathrm{Ni}$, and Cu , the energy separation involving the $\pi_{\mathrm{CO}}$ orbitals is smaller.

The second criterion that needs to be examined involves the overlap. The analysis here is simpler, since the $S_{\pi}$ and $S_{\pi^{*}}$ overlaps, for a given metal, are proportional to the magnitudes of the coefficients on carbon
in the $\pi_{\mathrm{CO}}$ and $\pi_{\mathrm{CO}}^{*} \mathrm{MO}$, respectively (we are neglecting here the overlap with that part of the MO located on the oxygen atom, since this is so much further from the metal centre). As a result of the polarization of the $\pi$ orbitals, the coefficient on carbon is far larger in the antibonding orbital ( 0.66 ) than in the bonding $(0.37)$. The ratio of the overlaps $S_{\pi^{*}} / S_{\pi}$ is therefore 1.78 . Since the interaction depends on the square of the overlap, we may conclude that the overlap term favours the $\pi$-acceptor character of the carbonyl ligand by a factor of about 3.2. For the metals located towards the left of the periodic table, the combination of a smaller energy difference and a much larger overlap ensures that the $d \leftrightarrow \pi_{\mathrm{CO}}^{*}$ interactions are far stronger than the $d \leftrightarrow \pi_{\mathrm{CO}}$ interactions, so the carbonyl ligand behaves as a double-face $\pi$-acceptor. For the metals located towards the right, the energy difference favours the $\pi$-donor character of the carbonyl ligand, but the overlap favours a $\pi$-acceptor behaviour. If we examine the extreme case of copper, (Table 3.C1), the overlap factor (about 3.2 in favour of the $d \leftrightarrow \pi_{\mathrm{CO}}^{*}$ interaction) is only slightly different from the energetic factor (about 2.9 in favour of the $d \leftrightarrow \pi_{\mathrm{CO}}$ interaction). This is the only case where the $\pi$-donor and $\pi$-acceptor characters are roughly balanced, and this arises because the $d$ orbitals on copper are very low in energy, compared to those on the other transition metals.

To conclude, an accurate treatment of the $\pi$ interactions that involve the carbonyl ligand requires both the occupied $\pi_{\mathrm{CO}}$ and empty $\pi_{\mathrm{CO}}^{*} \mathrm{MO}$ to be taken into account (a three-orbital interaction scheme, 3.C1). If one notes that the interactions involving the antibonding MO are dominant, one can simplify the description, by considering only the $\pi_{\text {CO }}^{*}$ orbitals on the carbonyl ligand. In this simplified model, three MO are considered for this ligand: the occupied orbital $\sigma_{\mathrm{C}}$, that characterizes the lone pair on carbon (Chapter 3, Figure 3.3), giving the ligand a $\sigma$-donor character, and the two empty $\pi_{\mathrm{CO}}^{*}$ orbitals, which confer a double-face $\pi$-acceptor character.

## Applications

A detailed knowledge of the electronic structure of transition metal complexes, and in particular of the shape and electronic occupation of the d-block orbitals, enables several problems related to their structure and reactivity to be studied. The examples discussed in this chapter illustrate a method for analysing these problems that usually relies on a study of orbital interactions between the ligand and the metal centre, and/or on a correlation diagram that links the orbitals of two different structures. The answers that one may hope to obtain from this type of analysis are qualitative rather than quantitative. For example, one can often determine which of two possible conformations of a complex is the more stable, and why, without being able to deduce the energy difference between them. Once the electronic factors that favour a particular structure are established, it may well also be possible to predict the type of changes that will follow from, for example, a change in the nature of the ligands. A qualitative interpretation of the structure and reactivity of complexes is exceedingly interesting for chemists, even if studies of a different type, that depend on accurate calculations, are necessary to provide theoretical data that may usefully be compared, quantitatively, with experimental results.

### 4.1. Conformational problems

Several conformations may sometimes be envisaged for a particular complex, depending on the orientation of a ligand with respect to the rest of the molecule. In such circumstances, the analysis for each conformation of the interactions between the ligand orbitals and those on the remaining fragment containing the metal often enables us to understand why one particular conformation is energetically favoured. We shall consider four examples in this section, using this 'fragment method', which involve mono- and bis-ethylene complexes and a molecular hydrogen complex; several other examples are presented at the end of the chapter in the form of exercises.

### 4.1.1. $d^{8}-\left[\mathrm{ML}_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$ complexes

Consider a trigonal-bipyramidal (TBP) complex in which an ethylene ligand occupies an equatorial site. The $d^{8}$ electronic configuration, which

leads to an 18 -electron complex, is the most common for this class of compounds (Chapter 2, § 2.5.2) and the complex $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$ is a well-known example.

Two limiting orientations may be envisaged for the ethylene ligand, depending on whether the carbon-carbon bond is perpendicular to the equatorial plane $(\mathbf{4 - 1 a})$ or in that plane $(\mathbf{4 - 1 \mathbf { b }})$. Both structures have $C_{2 v}$ symmetry and we shall use the planes of symmetry $P_{1}$ and $P_{2}(4-1)$ to analyse the symmetry of the most important orbitals.


The orbital structure of each conformation may be analysed as the result of the interaction between the occupied $\pi$ orbitals and the empty $\pi^{*}$ orbitals on ethylene on one hand, and the orbitals of a $d^{8} \mathrm{ML}_{4}$ fragment with a 'butterfly' geometry (Chapter 2, § 2.8.3) on the other. For this latter, we shall consider the four occupied orbitals of the d block (three nonbonding, one weakly antibonding) and the lowest-energy empty orbital, which is an $s-p$ hybrid orbital pointing towards the empty site of the TBP. The shapes of these orbitals are presented below (4-2) in the simplest case, where the four ligands are identical and only have $\sigma$ interactions with the metal. To simplify matters, the small contributions from the ligands to the two highest orbitals are not shown. The symmetries of the orbitals with respect to the planes $P_{1}$ and $P_{2}$ are also indicated in 4-2.

On passing from the perpendicular structure to the coplanar one (4-1), the orientation of the $\mathrm{ML}_{4}$ fragment does not change. However, there is a rotation of the ethylene ligand by $90^{\circ}$, and therefore also of the associated $\pi$ and $\pi^{*}$ orbitals. These are shown in $4-\mathbf{3 a}$ and $\mathbf{b}$, and their symmetries with respect to the planes $P_{1}$ and $P_{2}$ are given for each



4-3a (perpendicular)

SS



4-3b (coplanar)

[^4]conformation. ${ }^{1}$ We note already that the symmetry of the $\pi^{*}$ orbital changes from one conformation to the other, and this will prove to be crucial for the conformational preference for the complex.

The $\pi$ orbital on the ethylene ligand, whose symmetry is SS in both conformations, can interact with (i) the empty orbital and (ii) the occupied nonbonding orbital of the same symmetry on the $\mathrm{ML}_{4}$ fragment (4-2). The first of these interactions involves two electrons; it constitutes the donation interaction of the Dewar-Chatt-Duncanson model (Chapter 3, § 3.4.1.2) for the metal-olefin bond in this complex. It is easy to see that the overlap between these orbitals does not depend on the orientation of the olefin, due to the cylindrical symmetry of the empty orbital on the metal centre ( $4-4 \mathbf{a}$ and $\mathbf{b}$ ). Things are essentially the same for the second interaction: the overlap occurs mainly with the lobe of the $d$ orbital that points towards the $\pi$ orbital, and this lobe also has cylindrical symmetry with respect to the metal-olefin axis (4-5a and b). We may conclude that since the interactions which involve the $\pi$ orbital of the ethylene ligand are identical in the two conformations, they cannot contribute to a pronounced energetic preference for one of them.


We now consider the interaction of the empty $\pi^{*}$ orbital with the orbital of the same symmetry on the metal fragment. In the perpendicular conformation, the nonbonding orbital with SA symmetry is involved, but the antibonding AS orbital is concerned in the coplanar conformation (4-2). In both cases, this is a two-electron stabilizing interaction that leads to a transfer of electron density from the metal to the ligand. This is therefore the back-donation interaction in the Dewar-Chatt-Duncanson model (Chapter 3, § 3.4.1.2). As the orbital involved on the metallic fragment is not the same in the two conformations, the strength of this interaction depends on the conformation considered. Now the energetic stabilization created by an interaction is proportional to the square of the overlap and inversely proportional to the energy difference between the orbitals (Chapter 1, $\S$ 1.3.2). The antibonding $\pi^{*}$

Figure 4.1. Comparison of the back-donation interactions ( $d \rightarrow \pi^{*}$ ) in the perpendicular (on the left) and coplanar (on the right) conformations of a $d^{8}-\left[\mathrm{ML}_{4}\left(\eta^{2}\right.\right.$-ethylene $\left.)\right]$ complex (4-1). The energy levels and symmetries of the $d$ orbitals involved in the two cases are given in the centre of the diagram.

orbital is higher in energy than the d-block orbitals, which are nonbonding or, in one case, weakly antibonding (4-2 and Figure 4.1). The energy separation between the interacting orbitals is therefore smaller in the coplanar conformation, since, in this case, it is the highest-energy occupied orbital on the metallic fragment that is involved. Moreover, the overlap involves a $d$ orbital that is polarized towards the $\pi^{*}$ orbital in the coplanar conformation $\left(S_{\mathrm{c}}\right)$, but a pure $d$ orbital in the perpendicular case $\left(S_{\mathrm{p}}\right)$. Therefore, $S_{\mathrm{c}}>S_{\mathrm{p}}$, (4-6). In summary, a smaller energy separation and a larger overlap favour the back-donation interaction in the planar structure.


The greater electronic stabilization which follows ( $\Delta E_{\mathrm{c}}>\Delta E_{\mathrm{p}}$, Figure 4.1) leads to a preference for this conformation, which is indeed adopted in all known complexes of the type $d^{8}$-[ML ${ }_{4}\left(\eta^{2}\right.$-olefin $\left.)\right]$. The barrier to olefin rotation measured by nuclear magnetic resonance (NMR) is of the order of $10-15 \mathrm{kcal} \mathrm{mol}^{-1}$.

### 4.1.2. $d^{6}-\left[\mathrm{ML}_{5}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$ complexes: staggered or eclipsed conformation?

An octahedral complex with an ethylene ligand may adopt either the conformation in which the carbon-carbon bond eclipses the neighbouring $M-L$ bonds (4-7a), or the staggered conformation shown in 4-7b. This latter should be more stable on steric grounds. But experimentally, the eclipsed conformation is observed for complexes with a $d^{6}$ electronic configuration, such as $\left[\operatorname{Mo}\left(\mathrm{PR}_{3}\right)_{5}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$-hence the interest in this conformational problem.


The two conformations have $C_{2 v}$ symmetry, and as in the preceding example, we shall use the two planes of symmetry $P_{1}(y z)$ and $P_{2}(x z)$ to analyse the symmetries of the orbitals on the ethylene fragment ( $\pi$ and $\pi^{*}$ ) and of those on the $d^{6}$ metallic fragment ML5. For the latter, we shall consider the four lowest-energy $d$ orbitals (three doubly occupied nonbonding and the empty polarized $z^{2}$ orbital, see Chapter 2, § 2.3.1).

To pass from the eclipsed conformation (4-7a) to the staggered $(4-7 \mathbf{b})$, the orientation of the ethylene molecule has been fixed and a rotation of $45^{\circ}$ applied to the $\mathrm{ML}_{5}$ fragment. The orbitals of the ethylene fragment are therefore as shown in 4-8 for both conformations, whereas the reorientation of the metallic fragment leads to a change in the name and the symmetry of one of the its nonbonding orbitals (4-9a and $\mathbf{b}$ ): $x y(\mathbf{A A})$ becomes $x^{2}-y^{2}(\mathbf{S S})$ (Chapter 2, § 2.1.2.4 and 2.1.2.5).


In this system, the donation interaction of the Dewar-ChattDuncanson model involves the occupied $\pi$ orbital and the empty $z^{2}$ orbital, both with SS symmetry, in the two conformations. As the $z^{2}$ orbital has cylindrical symmetry, the overlap between $\pi$ and $z^{2}$ does not depend on the orientation of the ethylene ligand, so this interaction cannot lead to any conformational preference. The back-donation interaction involves the empty $\pi^{*}$ orbital and the occupied $y z$ orbital, both with symmetry SA, in the two conformations. This second


Figure 4.2. Interactions between the $\pi$ and $\pi^{*}$ orbitals on ethylene (in the centre) and the orbitals of the same symmetry on the $d^{6}-\mathrm{ML}_{5}$ fragment in eclipsed (left) and staggered (right) conformations.

Figure 4.3. Construction of the MO for the eclipsed and staggered conformations of the $d^{6}-\left[\mathrm{ML}_{5}\left(\eta^{2}\right.\right.$-ethylene $\left.)\right]$ complex, by interaction between the $\pi$ and $\pi^{*}$ orbitals on ethylene (centre) and the orbitals of the $d^{6}-\mathrm{ML}_{5}$ fragment (on the left for the eclipsed conformation, on the right for the staggered).
interaction is therefore also identical in the two conformations. So in this system, the conformational preference does not arise from either of the stabilizing interactions that have been invoked to describe the metal-ligand bond. If we consider the symmetries of the various orbitals in the eclipsed conformation, no additional interaction is possible, other than those we have already described. However, in the staggered structure, there is an interaction between the occupied $\pi$ orbitals and $x^{2}-y^{2}$, as both have $\mathbf{S S}$ symmetry. This destabilizing four-electron interaction, between the $\pi$ orbital on the olefin and one of the nonbonding $d$ orbitals on the metal (Figure 4.2), exists only in the staggered conformation; it is therefore the origin of the observed preference for the eclipsed conformation. The experimental barrier to olefin rotation is of the order of $10 \mathrm{kcal} \mathrm{mol}^{-1}$.

The electronic structure of the two conformations can be described more completely by constructing the MO that result from the interaction of the fragment orbitals. Those on ethylene are placed in the middle of Figure 4.3, while those of the metallic fragment are on the left (eclipsed) or on the right (staggered conformation). In the interests of clarity, only the doubly occupied MO of the full complex are shown. The back-donation interaction between the SA orbitals produces the same stabilization in both conformations. There is a stabilizing interaction (donation) that involves the $\mathbf{S S}$ orbitals in the eclipsed conformation (lefthand side). For the staggered conformation, the two-electron stabilizing interaction and the four-electron destabilizing interaction described above (Figure 4.2) are represented by a three-orbital interaction scheme ( $\mathbf{S S}$ symmetry) that involves four electrons.

The lowest-energy orbital in each conformation of the complex is mainly concentrated on the $\pi$ orbital of ethylene. This is a bonding metal-ligand orbital that does not belong to the d block. The three other occupied MO are mainly or entirely concentrated on the metal $d$ orbitals. They are the MO derived from the $t_{2 g}$ block of a regular octahedral

complex, the degeneracy being partially or completely lifted by the presence of the ethylene ligand. These three MO are all doubly occupied, since we are considering a complex with a $d^{6}$ electronic configuration. For each conformation, there is a d-block orbital that is stabilized by a bonding interaction with the $\pi^{*}$ orbital on ethylene and a nonbonding orbital. The preference for the eclipsed conformation arises from the third occupied orbital, which is nonbonding in the eclipsed structure but destabilized in the staggered case (SS). This destabilization should be linked to the analysis illustrated in Figure 4.2, which shows a repulsive interaction between SS orbitals in the staggered conformation.

In conclusion, we note that the highest-occupied molecular orbital (HOMO) 'rule', according to which the more stable of two structures of a given species is the one whose HOMO is at lower energy, is perfectly applicable to this conformational problem, and to the one treated in the preceding section (see Figure 4.1).

### 4.1.3. $d^{6}-\left[\mathrm{ML}_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)_{2}\right]$ complexes: coupling of two $\pi$-acceptor ligands

A new conformational problem arises when there are two ligands for which several orientations are possible, instead of just one: does the orientation of one of them have an influence on the other? Steric repulsions may, of course, be significant, but here we shall be interested in purely electronic factors that may induce a coupling between the orientations of the two ligands.

Octahedral complexes of the type $d^{6}-\left[\mathrm{ML}_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)_{2}\right]$, in which there are two ethylene ligands in trans positions, provide a characteristic example to illustrate this point. The problem of the orientation of each $\mathrm{C}=\mathrm{C}$ bond with respect to the rest of the complex has been treated in the preceding section, where it was shown that in the most stable conformation, these bonds eclipse the neighbouring $\mathrm{M}-\mathrm{L}$ bonds. If this result is accepted, there is still the question of the relative orientation of the two ethylene ligands. Two limiting structures can be imagined: the coplanar conformation (4-10a), which has $D_{2 h}$ symmetry, in which the two $\mathrm{C}=\mathrm{C}$ bonds are located in the same plane $(y z)$; and the perpendicular conformation ( $\mathbf{4 - 1 0 b}$ ), with $D_{2 d}$ symmetry, in which one of the ethylene molecules has turned by $90^{\circ}$ to be in the $x z$ plane. Experimentally, this second orientation is observed in this class of compounds, for example, in the complexes trans- $\left[\mathrm{Mo}\left(\mathrm{PMe}_{3}\right)_{4}\left(\mathrm{C}_{2} \mathrm{H}_{4}\right)_{2}\right]$ and trans$\left.[\mathrm{Mo} \text { (diphos) })_{2}\left(\mathrm{C}_{2} \mathrm{H}_{4}\right)_{2}\right]$ (diphos $=\mathrm{Ph}_{2} \mathrm{PCH}_{2} \mathrm{CH}_{2} \mathrm{PPh}_{2}$ ). The barrier to rotation about the metal-olefin bond has been estimated to be about $15 \mathrm{kcal} / \mathrm{mol}$ in this latter compound. Since the two ethylene molecules are well separated, by about $4 \AA$, it seems unlikely that steric factors are the origin of this conformational preference. It is therefore natural to
look for an electronic factor, by analysing the interactions between the ligands and the residual metallic fragment.


4-10a (coplanar)


4-10b (perpendicular)

We shall consider the $\pi$ and $\pi^{*}$ orbitals on the ethylene ligands in each conformation, as they interact with the orbitals of the squareplanar $d^{6}-\mathrm{ML}_{4}$ fragment in which the three nonbonding $d$ orbitals, $x y$, $x z$, and $y z$ (Chapter 2, § 2.2.1) are doubly occupied. Two other orbitals will also be considered in the following analysis, since they point towards the two ethylene molecules (4-11): the empty $z^{2}$ orbital which is weakly antibonding, and $p_{z}$ which is nonbonding but higher in energy (Chapter 2, § 2.2.2).



We consider first the interactions of the occupied $\pi$ orbitals. The situation is identical to that described in the preceding section for the eclipsed confirmation of a $d^{6}\left[\mathrm{ML}_{5}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$ complex: the overlap between a $\pi$ orbital and any of the occupied nonbonding orbitals on the metallic fragment is zero. However, there are two interactions with the empty $z^{2}$ and $p_{z}$ orbitals, as shown in 4-12 for the coplanar conformation. The in-phase combination of the $\pi_{1}$ and $\pi_{2}$ orbitals $\left(\pi^{(+)}\right)$interacts with $z^{2}(4-12 a)$, and the out-of-phase combination $\left(\pi^{(-)}\right)$with $p_{z}(4-12 b)$. These are two-electron interactions (ligand $\rightarrow$ metal donation), and after the interaction, only the bonding MO, mainly concentrated on the ligands, are doubly occupied. Due to the cylindrical symmetry of the $z^{2}$ and $p_{z}$ orbitals, there is no change to the overlaps involved if one of
the two ethylene molecules is rotated by $90^{\circ}$ to give the perpendicular conformation. The interaction between the two $\pi$ orbitals is changed, but the distance between the ligands is too large for this change to give any energetically significant consequences. The donation interactions do not, therefore, give rise to any pronounced conformational preference.


4-12a


4-12b

We now turn our attention to the interactions that involve the empty $\pi^{*}$ orbitals. As in the monoethylene complex (§ 4.1.2), the overlap for each of the $\pi^{*}$ orbitals with $z^{2}$ is zero by symmetry. The same is true for the $p_{z}$ orbital, which, like $z^{2}$, has cylindrical symmetry with respect to the $z$-axis. We must still consider the interactions with the nonbonding and doubly occupied $d$ orbitals (4-11), that is, the metal $\rightarrow$ ligands back-donation interactions. Following the method developed in $\S 3.3$ of the preceding chapter, we shall examine the overlap of the $\pi^{*}$ orbitals with each of the three nonbonding $d$ orbitals, remembering that a nonzero overlap leads to a bonding interaction $d \leftrightarrow \pi^{*}$ that stabilizes the $d$ orbital ( $\pi$-acceptor character of the ligand). In the planar conformation, the two $\pi^{*}$ orbitals, concentrated in the $y z$ plane, overlap with the $y z$ orbital which is therefore stabilized by two bonding interactions (4-13a). In other words, $y z$ is stabilized by the out-of-phase combination $\left(\pi^{*(-)}\right)$ of the orbitals $\pi_{1}^{*}$ and $\pi_{2}^{*}$. But the $x z$ and $x y$ orbitals, which are antisymmetric with respect to the $y z$ plane, have zero overlap with the $\pi^{*}$ orbitals, which are symmetric with respect to this plane. Their shapes and energies therefore stay unchanged (4-13b and c).


Figure 4.4. Diagram for the interaction between the nonbonding $d$ orbitals (in the centre) and the $\pi^{*}$ orbitals on the ethylene ligands in the coplanar ( $D_{2 h}$, on the left) and perpendicular ( $D_{2 d}$, on the right) conformations of an octahedral trans $\left[\mathrm{ML}_{4}\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right)_{2}\right]$ complex with a $d^{6}$ electronic configuration.

In the perpendicular conformation, the $\pi_{1}^{*}$ orbital still interacts with the orbital $y z(\mathbf{4 - 1 4 a})$, but $\pi_{2}^{*}$ now overlaps with the orbital $x z$ $(4-14 b)$, and $x y$ remains as a pure $d$ orbital (4-14c). In this conformation, two $d$ orbitals are therefore stabilized to an equivalent extent (bonding interactions), but the third is unchanged.


4-14a $\left(y z+\pi^{*}\right)$


4-14b $\left(x z+\pi^{*}\right)$


4-14c (xy)

These results are given as an interaction diagram in Figure 4.4; symmetry labels from the $D_{2 h}$ point group are used for the coplanar conformation (on the left), and from the $D_{2 d}$ point group for the perpendicular conformation (on the right). The $\pi^{*(+)}$ and $\pi^{*(-)}$ combinations in the coplanar structure have $b_{2 u}$ and $b_{3 g}$ symmetry, respectively, whereas the $\pi^{*}$ orbitals in the perpendicular structure form a degenerate pair of orbitals of e symmetry. Due to the large spatial separation between the ethylene ligands, the energies of the $\pi^{*(+)}$ and $\pi^{*(-)}$ combinations are very similar, and only slightly different from the energy of the degenerate $\pi^{*}$ orbitals in the perpendicular conformation. The three nonbonding $d$ orbitals on the metal, whose symmetries are $b_{3 g}$ $(y z), \mathrm{b}_{2 g}(x z)$, and $\mathrm{b}_{1 g}(x y)$ in the $D_{2 h}$ point group, and $\mathrm{e}(y z, x z)$ and $\mathrm{b}_{2}(x y)$ in the $D_{2 d}$ point group, respectively, are placed in the centre of the figure. Interaction between orbitals of the same symmetry leads, as we have shown above, to the stabilization of one $d$ orbital in the coplanar conformation (4-13a), whereas two of them are stabilized in the perpendicular conformation (4-14a and $\mathbf{b}$ ).


## Comment

In each conformation, the three doubly occupied MO are the orbitals derived from the $t_{2 g}$ block of a regular octahedral complex with a $d^{6}$ electronic configuration.

From the energetic point of view, four electrons are stabilized in the perpendicular conformation, but only two in the coplanar case. However, the stabilization of the orbital is larger for the coplanar conformation ( $\Delta E_{\mathrm{c}}>\Delta E_{\mathrm{p}}$ ), since there are two bonding interactions (4-13a), as compared to only one for the degenerate orbitals (4-14a and $\mathbf{b}$ ). In order to deduce the conformational preference, it is therefore necessary to know whether $\Delta E_{\mathrm{c}}$ is smaller than, equal to, or more than twice as large as $\Delta E_{\mathrm{p}}$.

Since these are interactions between orbitals with different energies, the stabilizations $\Delta E$ are proportional to $S^{2} / \Delta \varepsilon$, where $S$ is the overlap between the orbitals and $\Delta \varepsilon$ the energy difference between them (Chapter $1, \S$ 1.3.2). The term $\Delta \varepsilon$ is essentially the same for the conformations, since the energies of the $\pi^{*(-)}$ (coplanar conformation) and $\pi^{*}$ (perpendicular conformation) orbitals are almost identical. What can be said about the $S^{2}$ term? If $S_{d \pi^{*}}$ is the bonding overlap between a $d$ orbital and a $\pi^{*}$ orbital, the overlap between the fragment orbitals in the perpendicular conformation is $S_{\mathrm{p}}=S_{d \pi^{*}}$, and the total stabilization due to the four electrons (Figure 4.4, right-hand side) is proportional to $4 S_{d \pi^{*}}^{2}$. In the coplanar conformation, the overlap takes place between a $d$ orbital and the $\pi^{*(-)}$ orbital, which is written as $(1 / \sqrt{2})\left(\pi_{1}^{*}-\pi_{2}^{*}\right)$ in normalised form. The overlap involved, $S_{c}$, is therefore $(1 / \sqrt{2})\left(S_{d \pi^{*}}+S_{d \pi^{*}}\right)=\sqrt{2} S_{d \pi^{*}}$, and the stabilization of the two electrons (Figure 4.4, left-hand side) is proportional to $2\left(\sqrt{2} S_{d \pi^{*}}\right)^{2}=4 S_{d \pi^{*}}^{2}$. We therefore come to exactly the same result for the two conformations, so we can apparently deduce that there will be essentially free rotation for an ethylene ligand, but this conclusion is contradicted by the experimental results that were presented briefly at the beginning of this section, which indicate a barrier to rotation of about $15 \mathrm{kcal} \mathrm{mol}^{-1}$.

However, we must remember that the expression showing that the energetic stabilization is proportional to $S^{2} / \Delta \varepsilon$ holds only when there is a large energy difference $\Delta \varepsilon$ between the two interacting orbitals. If the energy gap is not large, then the stabilization is proportional to $S$ (Chapter 1, § 1.3.2, Note 8). In the case under discussion, the energy separation $d-\pi^{*}$ may be as little as about 2 eV for a molybdenum complex (an estimate obtained from calculations that use the extended Hückel method). It is therefore worthwhile to examine the other limiting case (small $\Delta \varepsilon$, stabilization proportional to $S$ ) to compare the two conformations. If we reuse the overlaps calculated above, we find that the stabilization is proportional to $2 \sqrt{2} S_{d \pi^{*}}$ in the coplanar
${ }^{2}$ G. J. Kubas, R. R. Ryan, B. I. Swanson, P. J. Vergamini, H. J. Wassermann J. Am. Chem. Soc. 106, 451 (1984). In the following 17 years, about 800 research articles appeared that are devoted to studies of the structures and reactivities of molecular hydrogen complexes.
conformation, but to $4 S_{d \pi^{*}}$ in the perpendicular conformation. The latter is now clearly favoured. This analysis shows us that we can anticipate a marked conformational preference for the perpendicular structure if the $d-\pi^{*}$ energy separation is not too large. As this gap increases, the energy difference between the two conformations should progressively decrease.

This preference for the perpendicular structure, which is more or less pronounced depending on the orbital energies, is an important result that can be generalized as follows: when two $\pi$-acceptor ligands are present, it is preferable for each to interact with a different occupied $d$ orbital, as happens in the perpendicular conformation, rather than for both to interact with the same $d$ orbital, as happens in the planar conformation.

### 4.1.4. Orientation of $\mathrm{H}_{2}$ in the 'Kubas complex' $\left[\mathrm{W}(\mathrm{CO})_{3}\left(\mathrm{PR}_{3}\right)_{2}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$

The preferential orientation of a ligand may also be caused by the presence of non-equivalent ligands on the metallic fragment. In this section, we shall treat what is known as a 'molecular hydrogen complex', that is, a complex in which a molecule of dihydrogen is bound to a metal centre in the $\eta^{2}$ mode. This molecule behaves as an L-type ligand, thanks to the two electrons in its $\sigma_{\mathrm{H}-\mathrm{H}}$ bond. The first example of this type of complex was characterized in 1984, by Kubas and co-workers. ${ }^{2}$ It is the $d^{6}$ octahedral complex $\left[\mathrm{W}(\mathrm{CO})_{3}\left(\mathrm{PR}_{3}\right)_{2}\left(\eta^{2}-\mathrm{H}_{2}\right)\right](4-15 a)$, where $\mathrm{R}={ }^{i} \mathrm{Pr}$, in which the length of the $\mathrm{H}-\mathrm{H}$ bond as measured by neutron diffraction $(0.82 \AA)$ is only slightly greater than that in an isolated $\mathrm{H}_{2}$ molecule ( $0.74 \AA$ ).

One of the remarkable aspects of the structure of this complex is linked to the non-equivalence of the ligands other than $\mathrm{H}_{2}\left(\mathrm{PR}_{3}\right.$ or CO$)$. The dihydrogen molecule is oriented parallel to the $\mathrm{W}-\mathrm{PR}_{3}$ bonds, even though at least one other conformation, that with $\mathrm{H}-\mathrm{H}$ parallel to the $\mathrm{W}-\mathrm{CO}$ bonds ( $\mathbf{4 - 1 5 \mathbf { b }}$ ), could be imagined. Is there an electronic factor that causes this conformation preference? Before we can answer this question, we have to analyse the orbital interactions that lead to the metal- $\mathrm{H}_{2}$ bond.

### 4.1.4.1. Orbital interactions in a $d^{6}-\left[\mathrm{ML}_{5}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$ complex

We consider first a model $d^{6}$ complex $\left[\mathrm{ML}_{5}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$ in which the five non $-\mathrm{H}_{2}$ ligands are identical and only have $\sigma$ interactions with the metal centre. We shall limit ourselves to the conformation in which the $\mathrm{H}-\mathrm{H}$ bond eclipses two neighbouring $\mathrm{M}-\mathrm{L}$ bonds (4-16).

Figure 4.5. Diagram of the interaction between the four lowest-energy $d$ orbitals for a $d^{6} \mathrm{ML}_{5}$ fragment (SBP) and the $\sigma_{\mathrm{H}_{2}}$ and $\sigma_{\mathrm{H}_{2}}^{*}$ orbitals of $\mathrm{H}_{2}$, that enables the MO of the $d^{6}$ model complex $\left[\mathrm{ML}_{5}\left(\eta^{2}-\mathrm{H}_{2}\right)\right](4-16)$ to be constructed.



The complex can be described as the result of interaction of an $\mathrm{H}_{2}$ molecule with a $d^{6}$ square-based pyramidal (SBP) ML ${ }_{5}$ metallic fragment. For the latter, we consider, as in $\S$ 4.1.2, the four lowest-energy $d$ orbitals (three of which are doubly occupied and nonbonding, the fourth being the polarized, empty $z^{2}$ orbital), and the $\sigma_{\mathrm{H}_{2}}$ and $\sigma_{\mathrm{H}_{2}}^{*}$ orbitals on $\mathrm{H}_{2}$ associated with the $\mathrm{H}-\mathrm{H}$ bond. Given the orbital symmetries with respect to the planes $P_{1}$ and $P_{2}(4-16)$, the interaction diagram shows two two-electron stabilizing interactions (Figure 4.5), between occupied $\sigma_{\mathrm{H}_{2}}$ and empty $z^{2}$ on the one hand ( $\mathbf{S S}$ symmetry), and between occupied $y z$ and empty $\sigma_{\mathrm{H}_{2}}^{*}$ on the other (SA symmetry).


$$
\text { 4-17a }\left(\mathrm{H}_{2} \rightarrow \mathrm{M}\right)
$$



4-17b $\left(\mathrm{M} \rightarrow \mathrm{H}_{2}\right)$

The first interaction, which leads to a transfer of electron density from dihydrogen to the metal, is of the donation type ( $4-17 \mathbf{a}$ ), whereas the second, which involves the empty orbital on $\mathrm{H}_{2}$, is a metal $\rightarrow$ dihydrogen back-donation interaction (4-17b).

## Comment

This bonding scheme is completely analogous to the Dewar-ChattDuncanson model for the interaction of an ethylene molecule with a metallic centre (Chapter 3, § 3.4), the $\pi_{\mathrm{CC}}$ and $\pi_{\mathrm{CC}}^{*}$ MO being replaced here by the $\sigma_{\mathrm{H}_{2}}$ and $\sigma_{\mathrm{H}_{2}}^{*}$ MO.

Notice that both interactions weaken the $\mathrm{H}-\mathrm{H}$ bond: donation reduces the population of the bonding orbital $\sigma_{\mathrm{H}_{2}}$, and back-donation increases the population of the antibonding orbital $\sigma_{\mathrm{H}_{2}}^{*}$. An increase of
${ }^{3}$ In solution, an equilibrium is established for the Kubas complex between the dihydrogen and dihydride forms, confirming that the molecular hydrogen complex is an intermediate in the oxidative addition reaction of $\mathrm{H}_{2}$.
the $\mathrm{H}-\mathrm{H}$ distance is indeed observed, to a greater or lesser extent, in all molecular hydrogen complexes. At the same time, both the occupied SS and SA MO are bonding between the metal and the hydrogen atoms (4-17), so $\mathrm{M}-\mathrm{H}$ bonds are being formed. This complex may therefore be viewed as a 'frozen structure' along the reaction pathway that leads to the rupture of the $\mathrm{H}-\mathrm{H}$ bond and the formation of two $\mathrm{M}-\mathrm{H}$ bonds (an oxidative addition reaction with the formation of a dihydride complex). ${ }^{3}$

### 4.1.4.2. Orientation of $\mathrm{H}_{2}$ in the Kubas complex

We shall consider two orientations for the $\mathrm{H}_{2}$ molecule. The $\mathrm{H}-\mathrm{H}$ bond is either parallel to the $\mathrm{W}-\mathrm{CO}$ bonds (4-18a), or parallel to the $\mathrm{W}-\mathrm{PR}_{3}$ bonds ( $\mathbf{4 - 1 8 b}$ ), the latter conformation being found experimentally.


As in the preceding examples ( $\S 4.1 .2$ and 4.1.3), the donation interaction does not depend on the orientation of the dihydrogen molecule, due to the cylindrical symmetry of the $z^{2}$ orbital (4-19a). The back-donation interaction involves either the $y z$ or the $x z$ orbital on the metallic fragment, in the first or second confirmations, respectively ( $\mathbf{4 - 1 9 b}$ and $\mathbf{c}$ ), and the $\sigma_{\mathrm{H}_{2}}^{*} \mathrm{MO}$ on dihydrogen. If all the ligands other than $\mathrm{H}_{2}$ are identical, the $x z$ and $y z$ orbitals are degenerate and there is therefore no difference between the two back-bonding interactions represented in 4$\mathbf{1 9 b}$ and $\mathbf{c}$. Both sketches therefore represent the eclipsed conformation of a $d^{6}\left[\mathrm{ML}_{5}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$ complex.


4-19a


4-19b


4-19c

In the Kubas complex, the presence of carbonyl ligands ( $\pi$ acceptors) and phosphines (which may, in a first approximation, be considered as
pure $\sigma$-donors), lifts the degeneracy of the nonbonding $d$ orbitals. As we have shown in Chapter 3 (§ 3.3.3.2), $y z$ is stabilized by three bonding interactions with $\pi_{\mathrm{CO}}^{*}, x y$ by two but $x z$ by only one (4-20); as a result, the energies of these orbitals follow the order $\varepsilon_{y z}<\varepsilon_{x y}<\varepsilon_{x z}$.



4-20
There is an important consequence of this energetic nonequivalence: in this complex, the back-donation interaction depends on the conformation that is adopted. In structure 4-18a, $\sigma_{\mathrm{H}_{2}}^{*}$ interacts with the lowest of the three $d$ orbitals $(y z)$, but with the highest $(x z)$ in 4-18b. Since the $\sigma_{\mathrm{H}_{2}}^{*}$ orbital is strongly antibonding, the energy gap ( $\Delta \varepsilon$ ) between it and any of the three $d$ orbitals is large. We may therefore assume that the stabilization created by the interaction is proportional to $S^{2} / \Delta \varepsilon$. The energetic factor $(\Delta \varepsilon)$ favours conformation $\mathbf{4 - 1 8 b}$, since the energy gap between the interacting orbitals is smaller (Figure 4.6). When we consider the overlap, it is important to realize that $y z$ is delocalized into three carbonyl groups, but $x z$ into only one (4-20). Since MO are always normalized, the coefficient of the $d$ orbital is larger for $x z\left(+1 \pi_{\mathrm{CO}}^{*}\right)$ than for $y z\left(+3 \pi_{\mathrm{CO}}^{*}\right)$, thereby favouring structure $\mathbf{4 - 1 8 \mathbf { b }}$ ( $S_{\sigma^{*}-x z}>S_{\sigma^{*}-y z}$, Figure 4.6). So both factors, energy gap and overlap, favour the conformation in which $\mathrm{H}-\mathrm{H}$ is parallel to the $\mathrm{W}-\mathrm{PR}_{3}$ bonds.

We note, in conclusion, that the conformation adopted by $\mathrm{H}_{2}$ in the Kubas complex is controlled by the back-donation interaction (4-15).

Figure 4.6. The back-donation interaction in conformations 4-18a (on the left, $\mathrm{H}-\mathrm{H}$ parallel to the $\mathrm{W}-\mathrm{CO}$ bonds) and $\mathbf{4 - 1 8 b}$ (on the right, $\mathrm{H}-\mathrm{H}$ parallel to the $\mathrm{W}-\mathrm{PR}_{3}$ bonds) of the complex $\left[\mathrm{W}(\mathrm{CO})_{3}\left(\mathrm{PR}_{3}\right)_{2}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]$.





4-21a


4-21b
${ }^{4}$ M. Brookhart, M. L. H. Green J. Organomet. Chem. 250, 395 (1983).

But since there is a large energy gap between the orbitals involved in this interaction, the conformational preference is rather weak, and the barrier to rotation of the $\mathrm{H}_{2}$ molecule is only about $2-3 \mathrm{kcal} \mathrm{mol}^{-1}$.

## 4.2. 'Abnormal' bond angles

The geometry of some complexes may be very different from one's initial thoughts, particularly for bond angles. We shall study two examples, to see which electronic factors may cause what appears to be a large distortion from an 'ideal' structure.

### 4.2.1. Agostic interactions

### 4.2.1.1. Introduction

In carbene complexes $\left[L_{n} M=C H R\right]$, the most common geometry is trigonal-planar about the carbon, with $\mathrm{M}-\mathrm{C}-\mathrm{H}_{\alpha}$ and $\mathrm{M}-\mathrm{C}-\mathrm{R}$ angles close to $120^{\circ}(4-21 a)$. But a large structural change occurs in some complexes: the $\mathrm{M}-\mathrm{C}-\mathrm{H}_{\alpha}$ angle decreases to about $80^{\circ}$, whereas the $M-C-R$ angle increases very considerably, to $160-170^{\circ}$ (4-21b). The deformation that is observed can therefore be described as a pivoting of the alkylidene ligand, which brings the $\mathrm{C}-\mathrm{H}_{\alpha}$ bond closer to the metal centre.

The same phenomenon can be observed with an alkyl ligand. For example, the methyl group can pivot, decreasing the $\mathrm{M}-\mathrm{C}-\mathrm{H}_{1}$ angle from the expected value of $109^{\circ}(4-22 a)$ to about $90^{\circ}(4-22 b)$.


4-22a


4-22b

All these complexes are electron-deficient $\left(N_{t}<18\right)$. The deformation that is observed has therefore often been interpreted as the consequence of the pressure to complete the valence shell of the metal, even if a substantial geometrical distortion is necessary, by introducing an additional interaction with a $\mathrm{C}-\mathrm{H}$ bond.

The group which brings the $\mathrm{C}-\mathrm{H}$ bond with its two electrons into the vicinity of the metal centre has been called 'agostic' by Brookhart and Green. ${ }^{4}$ These complexes can be considered as 'frozen' structures on the pathway for the insertion of a metal into a $\mathrm{C}-\mathrm{H}$ bond, in the same way that molecular hydrogen complexes were treated on the oxidative addition pathway that leads to a dihydride (§ 4.1.4).


5 This is an acceptable model for a qualitative analysis based on changes in overlaps, on the occupation of different orbitals, etc. But if more accurate calculations were to be undertaken for a quantitative study, then it would be necessary to study either the real complex, or a much less simplified model, such as $\left[\mathrm{Ti}(\mathrm{Cl})_{3}\left(\mathrm{PH}_{3}\right)_{2}\left(\mathrm{CH}_{3}\right)\right]$.

### 4.2.1.2. An agostic methyl group: the complex $\left[\mathrm{Ti}(\mathrm{Cl})_{3}\left(\mathrm{Me}_{2} \mathrm{P}-\left(\mathrm{CH}_{2}\right)_{2}-\mathrm{PMe}_{2}\right)\left(\mathrm{CH}_{3}\right)\right]$

The structure of the complex $\left[\mathrm{Ti}(\mathrm{Cl})_{3}\left(\mathrm{Me}_{2} \mathrm{P}-\left(\mathrm{CH}_{2}\right)_{2}-\mathrm{PMe}_{2}\right)\left(\mathrm{CH}_{3}\right)\right]$ obtained by neutron diffraction reveals an agostic methyl group, as the $\mathrm{Ti}-\mathrm{C}-\mathrm{H}_{1}$ angle is only $93.7^{\circ}(\mathbf{4 - 2 3})$. It is an essentially octahedral complex, in which the oxidation state of titanium is +4 . The electronic configuration is thus $d^{0}$, so the d block is completely empty! This complex is therefore severely electron-deficient, as $N_{t}=12$ (the electrons associated with the six metal-ligand bonds).

To simplify the problem, we shall analyse the model octahedral complex $\left[\mathrm{Ti}(\mathrm{H})_{5}\left(\mathrm{CH}_{3}\right)\right]^{2-}$, in which all the ligands except $\mathrm{CH}_{3}$ have been replaced by hydrogen atoms. The two negative charges are necessary to obtain a $d^{0}$ electronic configuration. Although this simplified complex has no experimental reality, it does possess the essential characteristics of the real complex 4-23: the same number of ligands, the same electronic configuration, and the presence of a methyl ligand that may be subject to an agostic distortion. ${ }^{5}$

We shall consider two structures in turn: the non-agostic case (4-24a) in which the $\mathrm{Ti}-\mathrm{C}-\mathrm{H}_{1}$ angle is $109^{\circ}$, and an agostic version (4-24b) where the $\mathrm{Ti}-\mathrm{C}-\mathrm{H}_{1}$ angle is close to $90^{\circ}$. To pass from one structure to the other, the methyl group is pivoted about the carbon atom, as in the experimental structure 4-23.


We shall analyse the interactions between the methyl group and the metallic centre by decomposing the complex $\left[\mathrm{Ti}(\mathrm{H})_{5}\left(\mathrm{CH}_{3}\right)\right]^{2-}$ into two fragments, $\mathrm{CH}_{3}^{-}$and $\left[\mathrm{TiH}_{5}\right]^{-}$. In the first, which is a pyramidal $\mathrm{AH}_{3}$ system, the lone pair on carbon is described by a hybrid orbital that is essentially nonbonding (see Chapter 1, Figure 1.4). This MO is the highest-energy occupied orbital; we shall refer to it as $n_{\mathrm{CH}_{3}}$, and at least initially, it is the only orbital that we shall consider on this fragment (Figure 4.7, right-hand side). The metallic fragment is of the $\mathrm{ML}_{5}$ type (SBP), where the metal is located in the base of the pyramid. As in the examples treated in $\S 4.1 .2$ and 4.1.4, we shall consider only the four lowest-energy $d$ orbitals on this fragment, three of which are nonbonding $(x y, x z$,

Figure 4.7. Orbital interaction diagram for the $\mathrm{CH}_{3}^{-}$and $\left[\mathrm{TiH}_{5}\right]^{-}$fragments in the non-agostic structure 4-24a.






and $y z)$ and one antibonding $\left(z^{2}\right)$. Since the electronic configuration of $\left[\mathrm{TiH}_{5}\right]^{-}$is $d^{0}$, these four orbitals are empty (Figure 4.7, left-hand side).

We consider first the non-agostic structure 4-24a, in which the local $C_{3}$ axis of the methyl group coincides with the $\mathrm{Ti}-\mathrm{C}$ bond. The orbital interaction diagram is presented in Figure 4.7. The orbital $n_{\mathrm{CH}_{3}}$ cannot interact with any of the nonbonding orbitals on the metallic fragment, since it is of $\sigma$-type and located in one of the nodal planes of the $x y$, $x z$, and $y z$ orbitals. However, its overlap with $z^{2}$ is very large, since these two orbitals both point along the $z$-axis and they are polarized towards each other. Their interaction produces a bonding MO, which is doubly occupied and mainly concentrated on the ligand; it characterizes the $\sigma_{\mathrm{Ti}}-_{\mathrm{C}}$ bond. The antibonding combination makes up one of the antibonding MO of the d block in this octahedral complex (the $z^{2}$ orbital).

The pivoting of the methyl group (4-24b) moves its local $C_{3}$ axis, so that the $n_{\mathrm{CH}_{3}}$ orbital is no longer oriented along the $z$-axis (4-25). The overlap between $n_{\mathrm{CH}_{3}}$ and $z^{2}$ therefore decreases, as does the electronic stabilization associated with this interaction. If this were the only factor to consider, this motion would therefore be energetically unfavourable, and no agostic distortion would be expected.


However, there is another significant change in the orbital interactions: the overlap between $n_{\mathrm{CH}_{3}}$ and $y z$, which was zero in the non-distorted structure, becomes substantial in the agostic structure (4-26). Since the $y z$ orbital is empty (a $d^{0}$ fragment), a new stabilizing interaction is created which now favours the agostic structure. Note that no interaction can take place between $n_{\mathrm{CH}_{3}}$ and $x y$ or $x z(S=0)$, whatever the geometry.


Two factors have therefore been revealed; one is unfavourable for the distortion, but the other favours it. How, then, can we explain why the second factor dominates the first? We need to consider the relative energies of the orbitals concerned. The energy gap $\Delta \varepsilon$ between $n_{\mathrm{CH}_{3}}$ and $y z$ (nonbonding) is much smaller than that between $n_{\mathrm{CH}_{3}}$ and $z^{2}$ (antibonding; see Figure 4.7). Now the electronic stabilization increases if the difference in energy between the interacting orbitals decreases. The agostic distortion therefore progressively replaces the ( $n_{\mathrm{CH}_{3}} \leftrightarrow z^{2}$ ) interaction by a stronger one ( $n_{\mathrm{CH}_{3}} \leftrightarrow y z$ ). Despite what might have been expected, this distortion reinforces the $\mathrm{Ti}-\mathrm{C}$ bond.

At this point in our analysis, we must note that the interpretation of the agostic distortion we have suggested does not really match the notion that it is due to an additional interaction between a hydrogen atom, or a $\mathrm{C}-\mathrm{H}$ bond, and the metal centre. Notice, however, that a weak M... H bonding interaction does appear in the MOs represented in 4-26 for the agostic structure (right-hand side). Moreover, one of the bonding MO of $\mathrm{CH}_{3}$ (often referred to as $\pi_{\mathrm{CH}_{3}}$ ), that mainly characterizes the $\mathrm{C}-\mathrm{H}_{1}$ bond, does have an increased interaction with the empty $y z$ orbital, due to the substantial overlap between $1 s_{\mathrm{H}_{1}}$ and $y z(4-27)$. This interaction


S 1
leads to a transfer of electron density from the ligand to the metal, weakening the $\mathrm{C}-\mathrm{H}_{1}$ bond (loss of bonding electrons) but creating an incipient $\mathrm{Ti}-\mathrm{H}$ bond (gain of bonding electrons). This factor adds to the one previously discussed, further stabilizing the agostic structure.

To conclude, the important point is that the agostic distortion occurs because there is a low-energy, empty orbital ( $y z$ in our example) in the nondistorted complex that can be used to make the distortion favourable. Both characteristics are important: (i) empty, because if it were occupied, the two-electron stabilizing interactions shown in 4-26 and 4-27 (righthand sides) would become repulsive four-electron interactions, which would therefore oppose the distortion. We are reminded of the need for an electron-deficient metal, that was noted at the beginning of this section; (ii) low in energy, so that the favourable factor described in $\mathbf{4 - 2 6}$ dominates the unfavourable factor described in 4-25. This aspect is also very important, since no agostic distortion of the methyl group is observed, for example, in the complex $\left[\mathrm{Ti}(\mathrm{Cl})_{3}\left(\mathrm{CH}_{3}\right)\right]$, even though it is extremely electron-deficient (formally, only eight electrons!). It can be shown that the empty orbital that could be used in an agostic distortion is in fact too high in energy in this complex to be useful; the three $\mathrm{Ti}-\mathrm{C}-\mathrm{H}$ angles remain equal.

### 4.2.2. $d^{6} \mathrm{ML}_{5}$ complexes: a 'T-shaped' or 'Y-shaped' geometry?

$d^{6} \mathrm{ML}_{5}$ complexes are electron-deficient species (16e) that often appear as reaction intermediates. They are indeed obtained from octahedral $d^{6} \mathrm{ML}_{6}$ complexes (18e) by the loss of a ligand, which is the first stage in many organometallic reactions. Despite this high reactivity, many of them have been isolated and geometrically characterized.

For a long time, the only coordination mode known for these complexes was the SBP. The group VI ( $\mathrm{Cr}, \mathrm{Mo}, \mathrm{W}$ ) pentacarbonyl complexes were the first to be observed, and many others with this geometry were later characterized with different metals ( $\mathrm{Ru}(\mathrm{II}), \mathrm{Rh}(\mathrm{III}), \mathrm{Ir}(\mathrm{III}), \mathrm{etc})$ and very varied ligands (alkyl, silyl, phosphine, halide, etc.). In all these complexes, the angle between two basal transoid ligands lies between $160^{\circ}$ and $180^{\circ}$; three examples are shown in 4-28.


4-28a


4-28b


4-28c
${ }^{6}$ The first two structures of this type were discovered in 1986: H. Werner, A. Hohn, M. Dziallas Angew. Chem. Int. Ed. Engl. 25, 1090 (1986); M. D. Fryzuk, P. A. McNeil, R. G. Ball J. Am. Chem. Soc. 108, 6414 (1986).


4-29a

There is nothing surprising about this coordination mode for diamagnetic $d^{6} \mathrm{ML}_{5}$ systems. The d block of an SBP complex, with the metal either in the basal plane of the pyramid or close to it, contains three nonbonding orbitals that are derived from the $t_{2 g}$ block of the octahedron (see Chapter 2, § 2.3.1). This environment is therefore ideal to accommodate three pairs of electrons. Subsequently, however, far more surprising structures were observed. ${ }^{6}$ These complexes adopt the geometry of a distorted trigonal bipyramid, where one of the equatorial angles is only some $75-80^{\circ}$. Three examples are shown in 4-29.


4-29b


4-29c

Notice that the movement from an SBP structure to a distorted TBP geometry can be described by the change of one bond angle ( $\alpha$ ) from $\sim 180^{\circ}$ to $\sim 75^{\circ}$, the intermediate value of $120^{\circ}$ corresponding to a regular TBP geometry (4-30). Referring to the relative positions of the ligands $\mathrm{L}_{1}, \mathrm{~L}_{2}$, and $\mathrm{L}_{3}$, the SBP structure is said to be ' T -shaped', and the distorted TBP to be 'Y-shaped'.


Three questions arise:

1. Why do the diamagnetic $d^{6} \mathrm{ML}_{5}$ complexes not adopt a regular TBP geometry ( $\alpha=120^{\circ}$ )?
2. Why are two different structures observed, 'T-shaped' and ' Y shaped'?
3. What are the electronic factors that favour one or other of these structures?

Figure 4.8. d-block of a TBP complex (the strongly antibonding orbital $z^{2}$ is not shown). In the interests of clarity, the axial ligands located on the $z$-axis are omitted. The electronic occupation is given for a $d^{6}$ paramagnetic $\mathrm{ML}_{5}$ complex (a triplet state).


The answer to the first question is straightforward, if we analyse the structure of the d block of a TBP complex (Figure 4.8, where the equatorial plane $x y$ is the plane of the page); two orbitals are nonbonding $(x z$ and $y z)$ and two are weakly antibonding $\left(x y\right.$ and $\left.x^{2}-y^{2}\right)$.

Six electrons must be placed in these orbitals. It is clear that the splitting pattern of the $d$ block favours a paramagnetic ground state, with one electron in each of the antibonding orbitals (a triplet state, Figure 4.8). However, this orbital arrangement is not favourable for diamagnetic complexes (all electrons paired), and all those we have mentioned are of this type. That is why none of them adopts the regular TBP structure.

### 4.2.2.1. Correlation diagram $Y \leftarrow T B P \rightarrow T$ ( $\sigma$ interactions only)

The geometries that are observed experimentally suggest that we should study the correlation diagram for the d-block orbitals as the TBP structure $\left(\alpha=120^{\circ}\right)$ is changed either to the T-shaped $\left(\alpha=180^{\circ}\right)$ or to the Y-shaped geometry ( $\alpha$ close to $80^{\circ}$ ). We shall suppose initially that the ligands only have $\sigma$-type interactions with the metal.

During this deformation, the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{3}$ move in the equatorial plane of the initial TBP (4-30). They therefore remain in one of the nodal planes ( $x y$ ) of the nonbonding $x z$ and $y z$ orbitals (Figure 4.7), so the shapes and energies of these two stay constant. But substantial changes do occur for the two other orbitals on leaving the regular TBP structure. In the T -shaped structure, the ligands $\mathrm{L}_{2}$ and $\mathrm{L}_{3}$ are situated



4-31
in a nodal plane of the $x y$ orbital, which therefore becomes nonbonding (4-31, right-hand side). On the other hand, the antibonding interaction with $L_{2}$ and $L_{3}$ increases in the Y-shaped structure, leading to a destabilization of the $x y$ orbital (4-31, left-hand side).


4-32
Turning to the $x^{2}-y^{2}$ orbital, movement to the T-shaped structure increases its antibonding interactions with $L_{2}$ and $L_{3}$ located on the $x$-axis, producing a destabilization (4-32, right-hand side). In contrast, these interactions are decreased, if not eliminated, in the Y-shaped structure, as these ligands are placed approximately on the nodal surfaces of the $x^{2}-y^{2}$ orbital (4-32, left-hand side). This orbital is therefore stabilized, but not to the point of becoming nonbonding, as there is still an antibonding interaction with ligand $\mathrm{L}_{1}$.

The correlation diagram shown in Figure 4.9 groups all these results. It is striking that the deformation of the TBP to the T- or Y-shaped structures stabilizes one of the two antibonding orbitals ( $x y$ for T, $x^{2}-y^{2}$ for Y) but destabilizes the other. As a consequence of this lifting of the orbital degeneracy, associated with a deformation of the complex which lowers its symmetry ( $D_{3 h} \rightarrow C_{2 v}$ ), the diamagnetic state in which the two electrons are paired is stabilized. This is an example of what is known as a Jahn-Teller distortion.

We are now able to understand why diamagnetic $d^{6} \mathrm{ML}_{5}$ complexes can adopt either the T- or Y-shaped structure, as there are three


Figure 4.9. Schematic correlation diagram for the four lowest-energy $d$ orbitals, linking the central TBP structure to the Y- or T-shaped geometries (left- and right-hand sides, respectively).
low-energy $d$ orbitals in each of them. In the first, the ground-state electronic configuration is $\left((x z)^{2},(y z)^{2},(x y)^{2}\right)$, whereas in the second it is $\left((x z)^{2},(y z)^{2},\left(x^{2}-y^{2}\right)^{2}\right)$. The T-shaped structure does appear a little more favourable from the diagram, since the three occupied orbitals are nonbonding.

### 4.2.2.2. The role of $\pi$-type interactions

The crucial electronic factor turns out to be the $\pi$-donor or $\pi$-acceptor properties of the ligand $\mathrm{L}_{1}$ that is trans to the angle $\alpha$. We consider first a $\pi$-acceptor ligand such as CO. It is easy to show that the $x z$ and $y z$ orbitals, both of which are occupied, are influenced by the $\pi_{\mathrm{CO}}^{*}$ orbitals to the same extent in the T- and Y-shaped structures: $x z$ stays unchanged (the overlap is zero by symmetry), whereas $y z$ is stabilized by a bonding interaction with the $\pi_{\mathrm{CO}}^{*}$ orbital that is parallel to the $z$-axis. It is the third occupied orbital that makes the difference: $x y$ (T-shaped structure) is stabilized by a bonding interaction, but $x^{2}-y^{2}$ (Y-shaped structure) stays unchanged, since its overlap with the $\pi_{\mathrm{CO}}^{*}$ orbitals is zero by symmetry (4-33). The preference for the T-shaped structure due to the $\sigma$ interactions is therefore reinforced by the presence of a $\pi$-acceptor ligand in the position $\mathrm{L}_{1}$.


4-33 $\left(\mathrm{L}_{1}=\pi \operatorname{acceptor}(\mathrm{A})\right)$

We now consider a $\pi$ donor such as Cl . The same analysis is applied, but we now find that the $x y$ orbital of the T-shaped structure is destabilized by an antibonding interaction with a lone pair on chlorine (4-34).


Y


T

$$
4-34\left(\mathrm{~L}_{1}=\pi \operatorname{donor}(\mathrm{D})\right)
$$


#### Abstract

The presence of a $\pi$-donor therefore destabilizes the T-shaped structure with respect to the Y -shaped one. Looking at the structures shown in 4-29, we do indeed notice that they all have a $\pi$-donor ligand (chloride, alkoxy, amino) trans to the acute angle $\alpha$. In the case of a single-face $\pi$-donor ligand (amino, for example), the destabilizing effect for the T-shaped structure shown in 4-34, which therefore favours the Y-shaped structure, can only occur if the lone pair on the donor lies in the equatorial plane, so that it can interact with the $x y$ orbital. That is just what happens in the complex shown in 4-29c.


### 4.3. Carbene complexes

### 4.3.1. Ambiguity in the electron count for carbene complexes

Carbene complexes, whose general formula is $\left[\mathrm{L}_{n} \mathrm{M}=\mathrm{CR}_{2}\right]$ and which formally contain an $\mathrm{M}=\mathrm{C}$ double bond, create a real problem for the calculation of the metal's oxidation state. This problem arises because the bent $\mathrm{CR}_{2}$ ligand possesses two nonbonding orbitals, close in energy, in which two electrons must be placed (see Chapter 1, Figure 1.5). The lower of these is a hybrid orbital $n_{\sigma}(4-35 a)$, whereas the higher is a pure $p$ orbital (if we consider the simplest example, methylene, $\mathrm{CH}_{2}$ ), $n_{p}(4-35 b)$.

In the covalent model, the two electrons are placed in the lowerenergy orbital (configuration $\left(n_{\sigma}\right)^{2}$ ), so that the carbene is considered as an L-type ligand which does not oxidize the metal (4-36a). However, due to the small difference in energy between the $n_{\sigma}$ and $n_{p}$ nonbonding MO, the ground electronic state of several carbenes $\left(\mathrm{CH}_{2}\right.$, for example) is in fact the triplet ${ }^{3}\left[\left(n_{\sigma}\right)^{1}\left(n_{p}\right)^{1}\right]$. In this situation, with two unpaired electrons, it seems more logical to consider the carbene as an $X_{2}$ ligand $(4-36$ b) which oxidizes the metal by two units.


If we now consider the ionic model, the electronic octet of the carbon must be completed, which means that we must consider the dianionic form $\left(\mathrm{CR}_{2}\right)^{2-}(\mathbf{4 - 3 6 c})$. Since the two extra electrons were supplied by the metal, we obtain the same oxidation state as that yielded by the covalent model with an $\mathrm{X}_{2}$ ligand. In the ionic description, the dianion $\left(\mathrm{CR}_{2}\right)^{2-}$ is a strong $\pi$-donor, due to its doubly occupied $p$ orbital.

It is important to be aware of these different possible points of departure, since the ambiguity in the calculation of the oxidation
state appears very frequently in experimental articles about these complexes. An example is given in 4-37, where we see that the complex $\left[\mathrm{Cp}_{2} \mathrm{Ta}\left(\mathrm{CH}_{3}\right)\left(\mathrm{CH}_{2}\right)\right]$ can be described as a $d^{2} \mathrm{Ta}$ (III) complex, if $\mathrm{CH}_{2}$ is considered as an L-type ligand in the covalent model, or as a $d^{0} \mathrm{Ta}(\mathrm{V})$ complex if we use either the $X_{2}$ description of the covalent model or the ionic description. It may well seem disconcerting not to know if the electronic configuration of this complex is $d^{2}$ or $d^{0}$, since that appears to imply that one does not know how many electrons must be placed in the d block! Note, however, that the total number of electrons does not depend on the model chosen $\left(N_{t}=18\right)$.


$$
\begin{aligned}
& \left.\begin{array}{l}
\mathrm{CH}_{2}=\mathrm{L} \Rightarrow\left[\mathrm{Ta}\left(\mathrm{~L}_{2} \mathrm{X}\right)_{2}(\mathrm{~L})(\mathrm{X})\right] \Rightarrow \mathrm{Ta}(\mathrm{III}): d^{2} \\
\mathrm{CH}_{2}=\mathrm{X}_{2} \Rightarrow>\left[\mathrm{Ta}\left(\mathrm{~L}_{2} \mathrm{X}_{2}\left(\mathrm{X}_{2}\right)(\mathrm{X})\right] \Rightarrow \mathrm{Ta}(\mathrm{~V}): d^{0}\right.
\end{array}\right\} \text { covalent } \\
& \left(\mathrm{CH}_{2}\right)^{2-} ; 2 \mathrm{Cp}^{-} ;\left(\mathrm{CH}_{3}\right)^{-}=>\mathrm{Ta}^{5+}: d^{0} \text { ionic }
\end{aligned}
$$

This apparently complicated situation (which is at least fairly complicated in reality) arises from the different possible choices for the electron distribution within the $\mathrm{CR}_{2}$ ligand ( L or $\mathrm{X}_{2}$ ?) or between the ligand and the metallic fragment (a neutral or a dianionic ligand?). The initial distribution is in fact somewhat arbitrary no matter which choice is made, and the situation becomes clearer if one considers the electronic structure of the complex as a whole, rather than that of the separate fragments.

### 4.3.2. Two limiting cases: Fischer carbenes and Schrock carbenes

To clarify these points, we shall consider the carbene as an L-type ligand (4-36a). It therefore acts as a $\sigma$ donor, using its lone pair described by the $n_{\sigma}$ orbital, which interacts with an empty orbital on the metal (e.g. $\left.z^{2}, 4-38 a\right)$. In this model, the $n_{p}$ orbital is empty, so the carbene acquires a $\pi$-acceptor character (single face) (4-38b). The interaction scheme is similar to that in the Dewar-Chatt-Duncanson model (Chapter 3, § 3.4.1) used, for example, to describe ethylene complexes or molecular hydrogen complexes ( $\S$ 4.1.4).


4-38a


4-38b

However, an important difference arises, since the empty $p$ orbital on the carbene can be strictly nonbonding (as in $\mathrm{CH}_{2}$ ), and therefore at low energy, rather than antibonding $\left(\pi_{\mathrm{CC}}^{*}, \sigma_{\mathrm{H}_{2}}^{*}\right)$ as in the preceding examples. Depending on the nature of the metal, the ligands and the substituents on the carbene, this orbital can therefore be either higher or lower in energy than the $d$ orbitals on the metal. To prove this point, one need only consider the energy of the $d$ orbitals for the first transition series, which ranges from $-8.5(\mathrm{Sc})$ to $-14 \mathrm{eV}(\mathrm{Cu})$, whereas the energy of a $p$ orbital on carbon is -11.4 eV (parameters of the extended Hückel method).

### 4.3.2.1. Two interaction schemes for back-donation

When the $p$ orbital is higher in energy than the $d$ orbital, we obtain a typical back-donation scheme, with the formation of a bonding MO mainly located on the metal (4-39): formally, we may consider that the $\pi$ interaction does not oxidize the metal, since the two electrons stay mainly localized on this centre. In this case, the description of the carbene as an L-type ligand is therefore quite appropriate. However, if the $p$ orbital is lower in energy than the $d$ orbital, the occupied bonding MO is mainly concentrated on the carbene (4-40), so that in a formal

sense, two electrons have been transferred from the metal to this ligand. In other words, the $d$ orbital that was occupied before the interaction is
transformed to an MO of the complex that is mainly concentrated on the ligand: the d block therefore 'loses' two electrons. In this case, the description of the carbene as an $\mathrm{X}_{2}$-type ligand or as a dianionic ligand $\left(\mathrm{CR}_{2}\right)^{2-}$, which leads to an increase of the metal's oxidation state by two units, corresponds more closely to reality.

The ambiguity that we noticed previously for the electron count in carbene complexes in fact corresponds to chemical reality: the L and $\mathrm{X}_{2}$ (or dianion) formulations are appropriate for the situations described in 4-39 and 4-40, respectively.

We shall now examine how the nature of the metal, the ligands, and the substituents on the carbene can favour one or other of these two possibilities.

### 4.3.2.2. Fischer carbenes and Schrock carbenes

How can the situation described in 4-39 be favoured, with an acceptor orbital on the carbene that is clearly higher in energy than the metal $d$ orbital? We require that the $d$ orbital is low in energy and at the same time that the $p$ orbital is high. The $d$ orbital factor therefore leads us to consider metals towards the right-hand side of the periodic classification, or at the limit in the centre, since the energy of the metal orbitals is lowered on passing from left to right in a transition series (see Chapter 1, Table 1.4). In addition, the presence of $\pi$-acceptor ligands also leads to a lowering of the level of the $d$ orbitals. As far as the carbene is concerned, the energy of the empty $p$ orbital is raised if the substituents are $\pi$ donors, that is, if they have lone pairs (halogens, $\mathrm{O}-\mathrm{R}, \mathrm{NR}_{2}$, etc.). An example is shown in 4-41 for the carbene $\mathrm{CCl}_{2}$. The in-phase combination of the $p$ lone pairs on the chlorine atoms

destabilizes the empty orbital on carbon and therefore favours the situation depicted in 4-39. Carbene complexes that possess these characteristics are called Fischer carbenes. As examples, we may consider complexes such as $\left[(\mathrm{CO})_{5} \mathrm{~W}=\mathrm{C}(\mathrm{Ph})(\mathrm{OMe})\right],\left[(\mathrm{CO})_{4} \mathrm{Fe}=\mathrm{C}(\mathrm{Ph})(\mathrm{OMe})\right]$, $\left[(\mathrm{CO})_{5} \mathrm{Cr}=\mathrm{C}\left(\mathrm{N}^{i} \mathrm{Pr}_{2}\right)(\mathrm{OEt})\right]$, or $\left[\mathrm{Cp}(\mathrm{CO})\left(\mathrm{PPh}_{3}\right) \mathrm{Fe}=\mathrm{CF}_{2}\right]^{+}$. Since the carbene is usually considered to be an L-type ligand for the calculation of the oxidation state of the metal, these are described as complexes of $\mathrm{W}(0), \mathrm{Fe}(0), \mathrm{Cr}(0)$, and $\mathrm{Fe}(\mathrm{II})$, respectively.

In this group of compounds, the carbene acts as a $\sigma$ donor and a rather weak $\pi$ acceptor, due to the relative energies of the $d$ and $p$ orbitals. Overall, the electron density around the carbon of the carbene is decreased, leading to a metal-carbon bond that is polarized $\mathrm{M}\left(\delta^{-}\right)=$ $\mathrm{C}\left(\delta^{+}\right)$, with an electrophilic character for the carbon centre: it is therefore likely to be attacked by nucleophiles, and it is possible, for example, to interconvert two carbene complexes (4-42).


In order to obtain the situation shown in $\mathbf{4 - 4 0}$, where the $p$ orbital on carbon is at lower energy than the $d$ orbital, the conditions previously suggested for the Fischer carbenes must be reversed: the metal must come from the left of the periodic table, and must have $\pi$-donor ligands to destabilize the $d$ orbital, instead of $\pi$-acceptors. Moreover, to ensure that the $p$ orbital on the carbene is as low in energy as possible, its substituents cannot be $\pi$ donors: $\mathrm{CH}_{2}$ itself is a good candidate, and, more generally, alkyl substituents are suitable (the term 'alkylidene' is often used for a carbene substituted by alkyl groups). The complexes $\left[\mathrm{Cp}_{2}\left(\mathrm{CH}_{3}\right) \mathrm{Ta}=\mathrm{CH}_{2}\right]$ and $\left[\mathrm{CpCl}_{2} \mathrm{Ta}=\mathrm{C}(\mathrm{H})\left(\mathrm{CMe}_{3}\right)\right]$ are typical examples. As we have already noted, the metal in this group of complexes is usually considered to be oxidized by two units by the alkylidene ligand, so the two examples mentioned above contain $\mathrm{Ta}(\mathrm{V})$.

Carbene complexes that possess these properties are called Schrock carbenes. The $\pi$ interaction shown in 4-40 is accompanied by a substantial transfer of electron density from the metal to the carbon atom. This transfer more than outweighs the $\sigma$ donation (ligand $\rightarrow$ metal), so that the metal-carbon bond is polarized in the sense $\mathrm{M}\left(\delta^{+}\right)=\mathrm{C}\left(\delta^{-}\right)$. Schrock carbenes therefore possess a nucleophilic carbon, and give, for example, addition products with Lewis acids such as $\mathrm{AlMe}_{3}(4-43)$.


4-43

In summary, Fischer carbenes (electrophiles) and Schrock carbenes (nucleophiles) correspond to two different possibilities for the orbitals that participate in the $\pi$ interaction ( $\mathbf{4 - 3 9}$ and $\mathbf{4 - 4 0}$ ). We often distinguish between them for the calculation of the metal's oxidation state: the carbene ligand is considered to be of L-type for the first category, but as an $\mathrm{X}_{2}$ ligand (or $\mathrm{CR}_{2}^{2-}$ ) for the second. The characteristics that we have suggested for each category are, of course, only general indications. Some complexes may be found on the border between the two groups. For example, the complex $\left[\mathrm{Cp}_{2}\left(\mathrm{CH}_{3}\right) \mathrm{Ta}=\mathrm{CH}_{2}\right]$ is, as expected, a nucleophilic carbene (Schrock-type), but the tungsten complex $\left[\mathrm{Cp}_{2}\left(\mathrm{CH}_{3}\right) \mathrm{W}=\mathrm{CH}_{2}\right]^{+}$, which has exactly the same ligands, is electrophilic (Fischer-type). The presence of a positive charge on the metallic fragment, which stabilizes the metal's $d$ orbitals and thereby favours the situation shown in 4-39, provides a plausible explanation of this difference in behaviour, despite the low energy of the nonbonding $p$ orbital on the carbon atom.

### 4.4. Bimetallic complexes: from a single to a quadruple bond

In organic chemistry, we know that two carbon atoms can be linked by a single bond $(\sigma)$, a double $(\sigma+\pi)$, or a triple bond $(\sigma+2 \pi)$. The increase in bond order is accompanied by a decrease in interatomic distance, from $1.54 \AA(\mathrm{C}-\mathrm{C})$, through $1.34(\mathrm{C}=\mathrm{C})$ to $1.20 \AA(\mathrm{C} \equiv \mathrm{C})$. From the orbital point of view, these different bond orders correspond to the occupation of one, two, or three bonding MO between the two carbon atoms (4-44), it being clear that the corresponding antibonding MO are empty. The orbital that characterizes the $\sigma$ bond has cylindrical symmetry about the internuclear axis in each case, whereas the orbitals that characterize the $\pi$ bond(s) have a nodal plane (4-44).

It is natural to wonder how many bonds, of what type, can exist between the two metal centres in bimetallic complexes. The characteristics presented above will reappear in the description of metal-metal bonds, but the presence of $d$-type orbitals introduces a significant new feature.





$\mathrm{HC} \equiv \mathrm{CH}$




4-44

### 4.4.1. $\sigma, \pi$, and $\delta$ interactions

We shall consider the interactions between the $d$ orbitals on two metallic centres, initially ignoring any influence of the ligands on these orbitals. Each interaction leads to the formation of a bonding and an antibonding MO, but only the first of these, written ${ }^{(+)}$, is shown in 4-45. The interaction between the $z^{2}$ orbitals (4-45a) is of $\sigma$ type; the axial overlap between the two orbitals has cylindrical symmetry around the internuclear axis $(z)$. The MO $x z^{(+)}$and $y z^{(+)}(4-45 \mathbf{b})$ have the same characteristics as the $\pi$ orbitals of acetylene (4-44); they are bonding, and they possess a nodal plane that passes through the nuclei.
$\sigma$

$z^{2(+)}$

4-45a

$x z^{(+)}$


$x y^{(+)}$


4-45c

The novel feature compared to organic systems arises in the interactions between the pairs of orbitals $x y$ and $x^{2}-y^{2}$. In each case, the bonding MO that are formed have two nodal planes, instead of just one for the $\pi$ MO: these are the $x z$ and $y z$ planes for the $x y^{(+)}$orbital, and the planes that bisect the $x$-and $y$-axes for the $\left(x^{2}-y^{2}\right)^{(+)}$orbital (4-45c). These are described as $\delta$ interactions.


4-46


Figure 4.10. Interaction scheme for the $d$ orbitals on two $\mathrm{ML}_{5}$ fragments (SBP) to form the MO of a bimetallic complex $\mathrm{M}_{2} \mathrm{~L}_{10}$.

The overlaps involved depend on the type of interaction. The axial overlap $S_{\sigma}$ is larger than or comparable to the lateral overlaps $S_{\pi}$. However, the $S_{\delta}$ overlap is much smaller than these, only about a third or a quarter of $S_{\pi}$.

### 4.4.2. $\mathrm{M}_{2} \mathrm{~L}_{10}$ complexes

Consider a bimetallic complex of the type $\mathrm{M}_{2} \mathrm{~L}_{10}$ (4-46), whose MO may be constructed by the interaction of two monometallic fragments $\mathrm{ML}_{5}$ with an SBP geometry.

On each fragment, we shall consider the four lowest-energy $d$ orbitals, as in the preceding examples that concerned this fragment (§ 4.1.2, 4.1.4, and 4.2.1.2): three are nonbonding ( $x y, x z$, and $y z$ ) and one is weakly antibonding $\left(z^{2}\right)$. These orbitals are not involved in the description of the $M-L$ bonds in each fragment, and are therefore available to form bonding and antibonding MO between the two metallic centres.

There is a very strong $\sigma$ interaction between the $z^{2}$ orbitals, since the two fragment orbitals are polarized towards each other, two $\pi$-type interactions between the $x z$ and $y z$ orbitals, and a weaker $\delta$-type interaction between the $x y$ orbitals (4-47).


The interaction diagram for these fragment orbitals is shown in Figure 4.10, where it is clear that the $\sigma$ interaction is stronger than the $\pi$ interactions, which in turn are larger than the $\delta$ interaction.

Since the $x y, x z$, and $y z$ orbitals are initially at the same energy, there is no ambiguity in the energetic order of the $\pi$ and $\delta \mathrm{MO}: \varepsilon_{\pi}<$ $\varepsilon_{\delta}<\varepsilon_{\delta^{*}}<\varepsilon_{\pi^{*}}$. Moreover, the highest-energy MO (not shown) is $\sigma^{*}$. However, it is less easy to be sure of the position of the $\sigma$ bonding MO. Although the interaction is the strongest of those concerned, the orbitals involved are the highest in energy $\left(z^{2}\right)$. The energy separation between the nonbonding and antibonding levels in each fragment, which is variable from one complex to another, is thus an important parameter. The ordering shown in Figure 4.10, where the $\sigma$ orbital is placed between the $\pi$ and $\delta$ levels, corresponds to a fairly small energy separation in the fragments. The number and the nature of the metal-metal bonds clearly depend on the electronic occupation of these orbitals.

We shall examine first the complex $\left[\mathrm{Re}_{2}(\mathrm{CO})_{10}\right]$. It may be considered to be formed from two $\left[\operatorname{Re}(\mathrm{CO})_{5}\right]$ fragments, whose electronic configuration is $d^{7}(\operatorname{Re}(0))$. There are therefore 14 electrons in all to be placed in these MO, which leads to the electronic configuration $(\pi)^{4}(\sigma)^{2}(\delta)^{2}\left(\delta^{*}\right)^{2}\left(\pi^{*}\right)^{4}(4-48 \mathbf{a})$. The interactions of $\pi$ and $\delta$ type therefore do not contribute to bonding between the metallic centres, since both the bonding and antibonding MO are occupied. However, the $\sigma$ orbital is occupied, whereas its antibonding $\sigma^{*}$ counterpart is empty. We can therefore conclude that there is a single bond between the two metallic centres, of $\sigma$ type, in the complex $\left[\mathrm{Re}_{2}(\mathrm{CO})_{10}\right]$.

${ }^{7}$ F. A Cotton Inorg. Chem. 4, 334 (1965)

It is clear that to increase the bond order, the number of electrons must be reduced, so that fewer antibonding MO are occupied. In the complex $\left[\mathrm{Re}_{2}(\mathrm{Cl})_{8}\left(\mathrm{H}_{2} \mathrm{O}\right)_{2}\right]^{2-}$, each monometallic fragment $\left[\mathrm{Re}(\mathrm{Cl})_{4}\left(\mathrm{H}_{2} \mathrm{O}\right)\right]^{-}$ has a $d^{4}$ electronic configuration $(\operatorname{Re}(I I I))$. There are therefore eight electrons to be placed in the MO of this complex, which leads to the electronic configuration $(\pi)^{4}(\sigma)^{2}(\delta)^{2}(\mathbf{4 - 4 8 b})$. As only the bonding orbitals are occupied, there is a quadruple bond between the two metal centres. There are two $\pi$ bonds, one $\sigma$ bond, and a new entity that does not exist organic chemistry: $a \delta$ bond.

The change from a single to a quadruple bond is accompanied by a very substantial shortening of the $\mathrm{Re}-\mathrm{Re}$ bond, whose length decreases from $3.04 \AA$ in $\left[\operatorname{Re}_{2}(\mathrm{CO})_{10}\right]$ to $2.22 \AA$ in $\left[\operatorname{Re}_{2}(\mathrm{Cl})_{8}\left(\mathrm{H}_{2} \mathrm{O}\right)_{2}\right]^{2-}$.

### 4.4.3. The $\left[\mathrm{Re}_{2}(\mathrm{Cl})_{8}\right]^{2-}$ complex: a staggered or an eclipsed conformation?

The most famous example of a quadruply bound bimetallic complex is $\left[\mathrm{Re}_{2}(\mathrm{Cl})_{8}\right]^{2-}$, whose structure inspired Cotton to propose, for the first time, the existence of a $\delta$ bond, in addition to $\sigma$ and $\pi$ bonds, between two metallic centres. ${ }^{7}$ The Re-Re distance is very short ( $2.24 \AA$ ), a value close to that found in the complex $\left[\mathrm{Re}_{2}(\mathrm{Cl})_{8}\left(\mathrm{H}_{2} \mathrm{O}\right)_{2}\right]^{2-}$ already studied above; most strikingly, the complex is observed to have an eclipsed structure (4-49a) rather than the staggered conformation (4-49b) that would have been expected. It is also important to note that this complex is diamagnetic.



4-49a (exp)


4-49b

As a first approximation, we can construct the MO for this complex by considering the interaction of two monometallic $\left[\mathrm{ReCl}_{4}\right]^{-}$fragments with a square-planar geometry (in fact, the $\mathrm{Re}-\mathrm{Re}-\mathrm{Cl}$ angles are $103.7^{\circ}$ ).

We consider first the eclipsed conformation (4-49a). The interaction scheme for the fragment orbitals is similar to that shown in Figure 4.10 for two SBP $\mathrm{ML}_{5}$ fragments which also have an eclipsed conformation. The only change in the $d$ orbitals of the fragments is the lowering in energy of the $z^{2}$ orbital, which is almost nonbonding in a square-planar $\mathrm{ML}_{4}$ complex (see Scheme 4-11), whereas it is antibonding in an SBP $\mathrm{ML}_{5}$ complex (Figure 4.10). As a consequence, the strong interaction
between the $z^{2}$ orbitals leads to a $\sigma$ MO that is now lower in energy than the $\pi$ orbitals. The energetic ordering of the MO is therefore: $\varepsilon_{\sigma}<\varepsilon_{\pi}<\varepsilon_{\delta}<\varepsilon_{\delta^{*}}<\varepsilon_{\pi^{*}}<\varepsilon_{\sigma^{*}}$. Since each fragment has a $d^{4}$ electronic configuration $(\operatorname{Re}(\mathrm{III}))$, there are eight electrons to be placed in these MO. In this diamagnetic complex, the ground-state electronic configuration $(\sigma)^{2}(\pi)^{4}(\delta)^{2}$ is thus characteristic of a quadruple bond between the two metallic centres (4-50a).


We now turn to the staggered conformation. One monometallic unit has thus been rotated by $45^{\circ}$ around the $z$-axis with respect to the other (4-49b). This has no consequence for the $\sigma$ interaction, due to the cylindrical symmetry of the orbitals concerned $\left(z^{2}\right)$. The $\pi$ interactions involve two degenerate orbitals on each centre that are concentrated in two perpendicular planes, (as in the ML ${ }_{5}$ fragments, 4-47). The sum of the two $\pi$ interactions is not changed by the rotation of one fragment, as the reduction of the overlap with one of the degenerate orbitals is compensated by the increase in overlap with the other. But an important change does occur for the $\delta$ interaction: the overlap between the $x y$ orbitals in the eclipsed conformation disappears in the staggered form, since the two orbitals that are now concerned $\left(x y^{(1)}\right.$ and $\left.\left(x^{2}-y^{2}\right)^{(2)}\right)$ have different symmetries (with respect to the plane of the paper, for example) (4-51). There are therefore two nonbonding $d$ levels in this conformation (4-50b).

$S \neq 0$

$S=0$

In the diamagnetic ground state, movement from the eclipsed to the staggered confirmation therefore leads to the rupture of the $\delta$ bond, which explains why the first structure (quadruply bound) is more stable than the second ('only' a triple bond). However, if we consider the triplet state ${ }^{3} \delta \delta^{*}$, the staggered conformation with two nonbonding $d$ electrons is more stable than the eclipsed (one electron in a $\delta$ bonding orbital, one in an antibonding $\delta^{*}$ ). As a result, the movement from the diamagnetic ground state to the first excited triplet state is accompanied by a change in conformation for the complex (eclipsed $\rightarrow$ staggered).

### 4.5. The reductive elimination reaction

In this last example, we shall study a chemical reaction. In general, problems concerning reactivity are more difficult to treat than structural problems. A reaction is accompanied by a complete electronic reorganization, involving the breaking and making of bonds, and several vital factors, such as the relative stabilities of the reactants and products, are not easy to obtain from qualitative analyses of electronic structure. However, some worthwhile information can be obtained, usually from a correlation diagram that relates the MO in the reactants to those in the products. We thus obtain a description of the electronic reorganization that is associated with the reaction under study, in terms of molecular orbitals.

### 4.5.1. Definition

The elimination of the molecule $R-R^{\prime}\left(R\right.$ and $R^{\prime}$ are $X$-type ligands such as H , alkyl, halogen, etc.) is a decomposition mode that is frequently found for organometallic complexes $\left[L_{n} M R R^{\prime}\right](4-52)$.

$$
\left[\mathrm{L}_{n} \mathrm{MRR}^{\prime}\right] \underset{\text { elimination }}{\text { reductive }}\left[\mathrm{L}_{n} \mathrm{M}\right]+\mathrm{R}-\mathrm{R}^{\prime}
$$

4-52
The removal of two X-type ligands leads to a decrease of two units in the metal's oxidation state; the electronic configuration changes from $d^{n}$ to $d^{n+2}$. The metal is therefore reduced, which explains the name 'reductive elimination'. The opposite reaction is called 'oxidative addition'.

### 4.5.2. Simplified model for the reaction <br> $\left[\mathrm{L}_{n} \mathrm{MR}_{2}\right] \rightarrow\left[\mathrm{L}_{n} \mathrm{M}\right]+R-R$

Four electrons are intimately involved in the reorganization of the bonds. In the reactant, these are the electrons associated with the two $M-R$

Figure 4.11. Schematic correlation diagram for the reductive elimination reaction $\left[\mathrm{L}_{n} \mathrm{MR}_{2}\right] \rightarrow\left[\mathrm{L}_{n} \mathrm{M}\right]+\mathrm{R}-\mathrm{R}$. Only the four most strongly perturbed orbitals are shown.
bonds that will be broken. In the products, they are the two electrons in the new $\mathrm{R}-\mathrm{R}$ bond and the two electrons that remain on the reduced metal. In the simplest description of this reaction by MO theory, we shall consider only the orbitals associated with these electrons, that is:

1. In the reactant, two (occupied) bonding MO that describe the bonds that are to be broken, and the two corresponding (empty) antibonding MO.
2. In the products, the (occupied) bonding and (empty) antibonding MO associated with the new bond that is formed ( $R-R$ ), and the two nonbonding MO that remain on the metal (with two electrons in total) after the departure of the two ligands.

A typical correlation diagram for these MO is shown in Figure 4.11. The MO of the reactant are placed on the left. The bonding orbitals, $\sigma^{(+)}$ and $\sigma^{(-)}$, are mainly concentrated on the ligands, whereas the antibonding orbitals, $\sigma^{*(+)}$ and $\sigma^{*(-)}$, are mainly on the metal (Chapter 1, § 6.1). We shall not try to discover the exact shape of these MO, in particular the contribution from the different atomic orbitals (AO) on the metal ( $s$, $p, d$, or mixture?). The representation shown here is schematic, since in reality the MO spread over the whole complex, rather than being limited to the two bonds. The important point is that there is a symmetric bonding MO $\left(\sigma^{(+)}\right)$and an antisymmetric bonding MO $\left(\sigma^{(-)}\right)$, just as there is a 'symmetric' $\left(\mathrm{a}_{1}\right)$ and an 'antisymmetric' ( $\mathrm{b}_{2}$ ) MO that describe the two A-H bonds in a bent $\mathrm{AH}_{2}$ molecule (Chapter 1, Figure 1.5). The same remark may be made about the antibonding $\sigma^{*(+)}$ and $\sigma^{*(-)}$. In the products, the MO that are $\mathrm{R}-\mathrm{R}$ bonding ( $\sigma_{\mathrm{R}-\mathrm{R}}$ ) and antibonding $\left(\sigma_{R-R}^{*}\right)$ are the lowest and highest in energy of the four orbitals, respectively. The relative energies of $\sigma_{\mathrm{R}-\mathrm{R}}$ and $\sigma^{(+)}$indicate that the


[^5]$R-R$ bond is assumed to be stronger than the $M-R$ bond. Among the two nonbonding orbitals ( $n_{\mathrm{A}}$ and $n_{\mathrm{S}}$ ), the antisymmetric orbital $n_{A}$ has been placed lower in energy than the symmetric $n_{S}$. This is the case when the first is a polarized $d$ orbital and the second a hybrid orbital of $s-p$ type. ${ }^{8}$

The diagram in Figure 4.11 shows us that the antisymmetric orbital $\sigma^{(-)}$is strongly destabilized, as the two metal-ligand bonding interactions are removed. The activation energy for the reaction is linked to the change in energy for this orbital. In contrast, the symmetrical orbital $\sigma^{*(+)}$ is strongly stabilized, as the two antibonding interactions disappear. The final point to note is the reduction of the metal centre. The $\sigma^{(-)}$orbital, which is mainly concentrated on the ligands, correlates with the $n_{\mathrm{A}}$ orbital that is localized on the metal: in a formal sense, the metal has gained two electrons.

### 4.5.3. An example: $d^{8}-\left[\mathrm{L}_{2} \mathrm{MR}_{2}\right] \rightarrow d^{10}-\left[\mathrm{L}_{2} \mathrm{M}\right]+\mathrm{R}-\mathrm{R}$.

We shall consider a square-planar $\mathrm{ML}_{4}$ complex with a $d^{8}$ electronic configuration which possesses two alkyl ligands ( R ) in cis positions. The reductive elimination of $\mathrm{R}_{2}$ leads to the formation of a $d^{10}\left[\mathrm{ML}_{2}\right]$ complex. An example is shown in 4-53, that involves the elimination of ethane from a complex that contains a chelating diphosphine ligand.


4-53

### 4.5.3.1. Correlation diagram for the highest-symmetry mechanism ( $C_{2 v}$ )

Several different mechanisms can be considered for this reaction, depending on the way in which $\mathrm{R}_{2}$ is eliminated. We shall study here the pathway with the highest symmetry possible; the orientation of the $\mathrm{R}_{2}$ unit with respect to the rest of the complex stays unchanged (4-54). $C_{2 v}$ symmetry is maintained along this reaction pathway, and the symmetry labels used for the orbitals will be appropriate.

In the correlation diagram, we shall represent not only the four MO described in the preceding section, but also the nonbonding or weakly antibonding orbitals in the d block and the other nonbonding orbitals on the metal. In this way, we shall obtain a more complete description

Figure 4.12. Correlation diagram for the MO involved in the reductive elimination reaction $d^{8}-\left[\mathrm{L}_{2} \mathrm{MR}_{2}\right] \rightarrow d^{10}-\left[\mathrm{L}_{2} \mathrm{M}\right]$ (bent) $+\mathrm{R}_{2}$ in the mechanism where $C_{2 v}$ symmetry is conserved.


4-54

of the electronic reorganization that is associated with this elimination reaction.

For the reactant (left-hand side of Figure 4.12), there are therefore the $\sigma_{M-R}$ bonding ( $1 \mathrm{a}_{1}$ and $1 \mathrm{~b}_{2}$ ) and $\sigma_{M-\mathrm{R}}^{*}$ antibonding ( $4 \mathrm{a}_{1}$ and $2 \mathrm{~b}_{2}$ ) MO and, at an intermediate energy, the three nonbonding $d$ orbitals $\left(x^{2}-y^{2}\right.$, $x z$, and $y z$ ), together with the slightly antibonding $z^{2}$ orbital and the nonbonding $p_{z}$ orbital (Chapter 2, § 2.2). The bonding MO are doubly occupied, as are the four $d$ orbitals $\left(d^{8}\right)$. For the products (right-hand side of Figure 4.12), we consider the $\sigma_{\mathrm{R}-\mathrm{R}}$ and $\sigma_{\mathrm{R}-\mathrm{R}}^{*} \mathrm{MO}\left(1 \mathrm{a}_{1}\right.$ and $2 \mathrm{~b}_{2}$ ), the five nonbonding ( $x^{2}-y^{2}, x z$, and $y z$ ), or weakly antibonding ( $z^{2}$ and, at slightly higher energy, $x y$ ) $d$ orbitals as well as the two nonbonding orbitals $p_{z}$ and $s-p_{y}\left(2 \mathrm{~b}_{1}\right.$ and $\left.4 \mathrm{a}_{1}\right)$ (Chapter 2, § 2.8.4). In the ground-state electronic configuration, the $\sigma_{\mathrm{R}-\mathrm{R}} \mathrm{MO}$ and the five $d$ orbitals are doubly occupied ( $d^{10}$ ). $C_{2 v}$ labels are used for the symmetries of these orbitals.

The first conclusion that we may draw from this correlation diagram is that the elimination reaction is symmetry-allowed according to this mechanism: the ground-state electronic configuration of the reactant correlates with the ground state of the products. Moreover, the energy changes of the four 'principal' MO (solid lines) are the same as those given in the simplified diagram of the previous section (Figure 4.11). In particular, notice the destabilization of the $1 \mathrm{~b}_{2}$ orbital and its transformation from a ligand orbital in the initial complex into a d-block orbital of the $\mathrm{ML}_{2}$ complex.

### 4.5.3.2. Use of the correlation diagram: influence of the nature of the metal

Experimentally, it has been shown that the reductive elimination of $R_{2}$ is easier for a nickel complex than for the isoelectronic palladiumcontaining species. This result may be readily understood from the correlation diagram in Figure 4.12.


As we have already noted, the activation energy of this reaction is mainly caused by the destabilization of the $1 \mathrm{~b}_{2}$ orbital. In the reactant, this MO is mainly concentrated on the ligands, but in the products it is a metal-based orbital, since it is part of the $d$ block of the $\left[\mathrm{ML}_{2}\right]$ complex. On moving from palladium to nickel, the energy of the $d$ orbitals is lowered appreciably, from -9.58 to -12.92 eV (see Table 1.4, Chapter 1). This lowering in energy stabilizes the $1 \mathrm{~b}_{2}$ orbital, only weakly in the reactant (small coefficient on the metal) but strongly in the product (a $d$ orbital). This analysis shows us that the destabilization of the $1 \mathrm{~b}_{2}$ orbital is smaller for a nickel complex than for one containing palladium (4-55). The activation energies follow the same trend, and the greater tendency for a nickel complex to undergo reductive elimination is thereby explained.

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## Exercises

Conformation of $d^{10}-\left[\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right) \mathrm{ML}_{2}\right]$ complexes

## 4.1

1. Which is the more stable conformation for an $\left[\left(\eta^{2}-\mathrm{C}_{2} \mathrm{H}_{4}\right) \mathrm{ML}_{2}\right]$ complex with a $d^{10}$ electronic configuration, 1 or 2 ?


1


2
2. How does the barrier to olefin rotation change when the substituents R are varied: $\mathrm{R}=\mathrm{CN}$ (tetracyanoethylene), $\mathrm{R}=\mathrm{H}$ (ethylene), $\mathrm{R}=\mathrm{Cl}$ (tetrachloroethylene)?
3. Repeat question 2, replacing the two non-ethylene ligands on the metal by a bidentate ligand that constrains the $L-M-L$ angle to be much larger than $120^{\circ}$.
$\mathrm{ML}_{3}$ complexes

## 4.2

1. Starting from the $d$ orbitals of an octahedron, construct the d orbitals (shapes and relative energies) for a 'T-shaped' $\mathrm{ML}_{3}$ fragment $\left(\mathrm{L}_{2}-\mathrm{M}-\mathrm{L}_{3}=180^{\circ}, \mathrm{L}_{1}-\mathrm{M}-\mathrm{L}_{2}=\mathrm{L}_{1}-\mathrm{M}-\mathrm{L}_{3}=90^{\circ}\right)$.

2. How do the orbital energies of this fragment change on moving from the T-shape to a trigonal geometry ( $120^{\circ}$ bond angles)?

3. From the results you obtain in 2 above, predict the geometries adopted by the complexes $\left[\mathrm{Rh}\left(\mathrm{PPh}_{3}\right)_{3}\right]^{+}$and $\left[\mathrm{Pt}\left(\mathrm{PPh}_{3}\right)_{3}\right]$.

## Octahedral dioxo complexes

## 4.3

We wish to compare the relative stabilities of the cis and trans isomers of $d^{0}$ and $d^{2}$ octahedral dioxo complexes; remember that the oxo group is an $\mathrm{X}_{2}$-type ligand.


Trans


Cis


1. Give the shapes and relative energies of the $d$ orbitals for an octahedral complex in the absence of any $\pi$-type interactions, in the two orientations shown above. In what follows, we shall only consider the nonbonding d orbitals.
2. We now turn to dioxo complexes. If the four non-oxo ligands are pure $\sigma$ donors, give the interaction schemes $d$ (metal) $\leftrightarrow p$ (oxo) for each isomer.
3. Which isomer is the more stable for a $d^{0}$ complex? For a $d^{2}$ complex?
4. Deduce the structures of the complexes $\left[\mathrm{OsO}_{2} \mathrm{~F}_{4}\right]^{2-}$ and $\left[\mathrm{MoO}_{2} \mathrm{Cl}_{4}\right]^{2-}$.

## Carbene-dihydrogen coupling in an octahedral complex

## 4.4

In the molecular hydrogen complex $\left[\mathrm{Os}\left(\mathrm{NH}_{3}\right)_{4}\left(\mathrm{CR}_{2}\right)\left(\mathrm{H}_{2}\right)\right]^{2+}$, a Fischer carbene and a dihydrogen molecule are in cis positions. We shall consider four limiting structures $(0,0),(0,90),(90,0)$, and ( 90 , 90 ) that are characterized by the orientations of the carbene (first angle) and of the dihydrogen molecule (second angle).


1. What is the electronic configuration $d^{n}$ for this complex?
2. For each of the four conformations above, analyse the backdonation interactions towards the carbene and towards the dihydrogen.
3. From the results of this analysis, suggest the favoured structure(s) for this complex.
4. Which limiting conformations are possible for the trans isomer? Which should be the most stable a priori?

## Coupling of two trans carbenes in an octahedral complex

## 4.5

Consider the conformations 1 and 2 of octahedral complexes in which two Fischer carbenes are in trans positions. The non-carbene ligands are pure $\sigma$ donors.

1. Give the shapes and relative energies of the three lowest-energy d-block MOs for the two conformations.


1


2
2. Deduce the conformations for $a d^{2}, d^{4}$, or $d^{6}$ diamagnetic complex.
3. How is the barrier to rotation modified if the non-carbene ligands are $\pi$ acceptors?

## The reductive elimination reaction

## 4.6

We wish to study the reductive elimination of $\mathrm{H}_{2}$ from an octahedral complex with a $d^{6}$ electronic configuration, following the mechanism shown in $\mathbf{4 - 5 6}$, which leads to a square-planar $d^{8} \mathrm{ML}_{4}$ complex.


4-56

1. Sketch the correlation diagram for the MO associated with this reaction. For the reactant, consider the nonbonding $d \mathrm{MO}$ and the MO that characterize the two $\mathrm{M}-\mathrm{H}$ bonds which will be broken; for the products, the MO of $\mathrm{H}_{2}$, the occupied $d$ orbitals, and the nonbonding $p$ orbital of the $\mathrm{ML}_{4}$ complex. Two planes should be used to characterize the symmetries of the MO.
2. Is this reaction allowed by symmetry?
3. How does the correlation diagram show the reduction of the metal?

## The isolobal analogy

${ }^{1}$ R. Hoffmann 'Building bridges between inorganic and organic chemistry (Nobel lecture)', Angew. Chem. Int. Ed. 21, 711 (1982).
${ }^{2}$ J. Halpern in Advances in Chemistry; Homogeneous Catalysis, No 70, American Chemical Society, Washington, DC (1968) pp. 1-24.

In Chapter 2, the $d$ orbitals (and, in some cases, the $s, p$, or $s-p$ hybrid orbitals on the metal) were constructed for many types of $\mathrm{ML}_{n}$ complexes in which the ligands have $\sigma$-type interactions with the metal. The shapes and relative energies of these orbitals depend on the number of ligands and their geometrical arrangement around the metal. Naturally, any $\pi$-type interactions that are present also play a role. Thus, the $d$ orbitals in the octahedral complex $\left[\mathrm{WCl}_{6}\right]$ (six $\pi$-donor ligands) are different from those in the octahedral complex $\left[\mathrm{W}(\mathrm{CO})_{6}\right]$ (six $\pi$-acceptor ligands). But 'there is a time for detail and there is a time for generality': in both cases, the $d$ orbitals are split $3+2$, with three degenerate orbitals (the $\mathrm{t}_{2 g}$ block) lower in energy than the two degenerate orbitals of the $\mathrm{e}_{g}$ block, and this result is characteristic of all octahedral $\mathrm{ML}_{6}$ complexes. In the same way, the orbital structure of $\mathrm{AH}_{n}$ molecules or fragments (or more generally $\mathrm{AR}_{n}$ ), in which A is a main-group element, depends essentially on the number and geometrical arrangement of the substituents on the central atom. But beyond the obvious difference created by the presence of $d$ orbitals in organometallic complexes, there are resemblances between the molecular orbitals (MO) of $\mathrm{ML}_{n}$ complexes and those of $\mathrm{AH}_{n}$ molecules. These similarities can lead to related properties: thus J. Halpern had noted in 1968 the similarity in behaviour of organic radicals and of $\mathrm{d}^{7} \mathrm{ML}_{5}$ complexes. ${ }^{2}$ But it is chiefly R. Hoffmann who developed this concept and showed its remarkable fruitfulness. ${ }^{1}$

### 5.1. The analogy between fragments of octahedral $\mathrm{ML}_{6}$ and of tetrahedral $\mathrm{CH}_{4}$

In many cases, we established the orbitals of $\mathrm{ML}_{n}$ fragments earlier in this book by starting from a complex in which one or several ligands were removed. For example, the orbitals of $\mathrm{ML}_{5}$ complexes (SBP, $C_{4 v}$ ), $\mathrm{ML}_{4}$ ('butterfly', $C_{2 v}$ ), and $\mathrm{ML}_{3}$ (pyramidal, $C_{3 v}$ ) were obtained from those of an octahedral $\mathrm{ML}_{6}$ complex, by removing one, two, or three ligands, respectively (Chapter 2, § 2.3.1, 2.8.1, and 2.8.3). In certain respects, each of these complexes was therefore considered as a fragment of an octahedron (5-1).

$\mathrm{ML}_{6}$



5-1

The same procedure can be applied to molecules of the type $\mathrm{AH}_{n}$, where A is an element from the second or third row of the periodic table. In particular, starting from tetrahedral methane, $\mathrm{CH}_{4}$, the archetype of organic molecules, one can obtain the organic fragments pyramidal $\mathrm{CH}_{3}\left(C_{3 v}\right)$, bent $\mathrm{CH}_{2}\left(C_{2 v}\right)$, and $\mathrm{CH}\left(C_{\infty v}\right)$, by removing one, two, or three atoms of hydrogen (5-2).



## 5-2

Even though these comments are purely geometrical in nature, it turns out that they have profound chemical consequences. Thus, an ML, complex (SBP) possesses a vacant site, and it is likely to bind an additional ligand to form an octahedral complex. In the same way, the $\mathrm{ML}_{4}\left(C_{2 v}\right)$ and $\mathrm{ML}_{3}\left(C_{3 v}\right)$ fragments have two and three available sites, respectively, with respect to the initial octahedron. In the organic series,

[^6]the $\mathrm{CH}_{3}, \mathrm{CH}_{2}$, and CH fragments possess the same characteristics, with the potential to form one, two, or three new bonds.

This similarity in the capacities of inorganic and organic fragments to form new bonds can be expressed in orbital terms, particularly if we consider the number, the shape, and the electronic occupation of the orbitals that are available to form these new bonds. The resemblances between these orbitals are at the heart of the theory of the isolobal analogy, developed by R. Hoffmann, which enables us to build bridges between the electronic structures of inorganic and organic molecules.

### 5.1.1. Fragment orbitals by the valence-bond method

The orbitals of methane, $\mathrm{CH}_{4}$, and those of the related fragments $\mathrm{CH}_{3}$, $\mathrm{CH}_{2}$, and CH can be described using the molecular orbital method, as we have done for all the systems studied so far in this book. But the valence-bond approach, introduced by L. Pauling, ${ }^{3}$ can also be used; this is perhaps the simplest way to establish an initial relationship between the electronic structures of organic and inorganic fragments.

In $\mathrm{CH}_{4}$, the carbon atom is at the centre of a tetrahedron defined by the four hydrogen atoms. The carbon atom's four valence orbitals (the $2 s$ orbital and the three $2 p$ orbitals) can be combined to give four equivalent hybrid orbitals (written $s p^{3}$ ) that point towards the vertices of the tetrahedron. Since carbon has four valence electrons, one electron can be placed in each hybrid orbital. A $1 s_{\mathrm{H}}$ orbital, which is also singly occupied, can interact with the $s p^{3}$ hybrid that points towards it, and in this way we can describe the four equivalent bonds in a tetrahedral arrangement.

Using this model, how can we describe the orbital structure of the $\mathrm{CH}_{3}, \mathrm{CH}_{2}$, and CH fragments that are obtained by homolytic rupture of one, two, or three $\mathrm{C}-\mathrm{H}$ bonds? In the case of $\mathrm{CH}_{3}$, three of the hybrid orbitals point towards three hydrogen atoms and interact strongly with the $1 s_{\mathrm{H}}$ orbitals. But the fourth orbital points in the direction where there is no hydrogen atom (5-3a). It is therefore a nonbonding hybrid orbital, which contains a single electron (from the homolytic rupture of the $\mathrm{C}-\mathrm{H}$ bond). In the same way, there are two nonbonding singly occupied $s p^{3}$ hybrid orbitals in $\mathrm{CH}_{2}(5-3 \mathbf{b})$, and three in $\mathrm{CH}(5-3 \mathbf{c})$.


5-3a


5-3b


5-3c

We now turn to an octahedral $\mathrm{ML}_{6}$ complex with a $d^{6}$ electronic configuration, which we shall consider as the starting point for inorganic fragments. This complex obeys the 18 -electron rule, just as methane obeys the octet rule. To be a little more concrete, though this choice is in no way unique, we shall consider a chromium complex [ $\mathrm{CrL}_{6}$ ], with six neutral L-type ligands ( $\mathrm{PR}_{3}, \mathrm{CO}$, etc.), each of which supplies a pair of electrons to the metal. In this complex, chromium is in the oxidation state zero, and the electronic configuration is indeed $d^{6}$.

In the spirit of valence-bond theory, we must construct six hybrid orbitals that point towards the vertices of the octahedron. We must therefore combine just six of the nine valence orbitals of the metal: it turns out that these are the $s$ orbital, the three $p$ orbitals, and two of the $d$ orbitals (this hybridisation pattern will be written $d^{2} s p^{3}$, see Chapter 2, § 2.1.2.3). Three $d$ orbitals are therefore not involved in the hybridisation, and they stay unchanged (5-4). The link with MO theory is plain: these are the three nonbonding orbitals of the octahedral $t_{2 g}$ block. In a $d^{6}$ complex, these three non-hybridized orbitals are doubly occupied.



5-5

With the six metal electrons placed in this way, the hybrid orbitals are empty. Each of these interacts with the doubly occupied orbital of the ligand L towards which it points, thereby forming six bonds in an octahedral arrangement (5-5).

We now break one $\mathrm{Cr}-\mathrm{L}$ bond homolytically, as we did above for methane. If the ligand is neutral, and of L type $\left(\mathrm{CO}, \mathrm{PR}_{3}, \ldots\right)$, this homolytic rupture leaves one of the bonding electrons on the metal, forming the radicals ${ }^{\prime}{ }^{+}$and $\left[{ }^{\cdot} \mathrm{CrL}_{5}\right]^{-}$. In the latter (5-6a, right-hand side), there are still five $d^{2} s p^{3}$ hybrid orbitals that interact strongly with the orbitals on the five ligands. The sixth orbital points towards the vacant site of the octahedron: it is therefore nonbonding, and it contains the electron that came from the homolytic rupture of the $\mathrm{Cr}-\mathrm{L}$ bond. From this viewpoint, the resemblance with $\mathrm{CH}_{3}$ is striking (5-6a): the
$\mathrm{CH}_{3}$ and $\left[\mathrm{CrL}_{5}\right]^{-}$fragments are both characterized by the presence of a singly occupied nonbonding orbital which points towards a vacant site, of a tetrahedron for the former, of an octahedron for the latter.


5-6a
5-6b


5-6c

In the same way, the homolytic rupture of a second $\mathrm{Cr}-\mathrm{L}$ bond that leads to the complex $\left[\mathrm{CrL}_{4}\right]^{2-}$ with a 'butterfly' geometry leaves two singly occupied nonbonding hybrid orbitals on the metal, as in the organic fragment $\mathrm{CH}_{2}(5-6 \mathbf{b})$. And the rupture of a third bond leads to the fragment $\left[\mathrm{CrL}_{3}\right]^{3-}(5-6 \mathrm{c})$, whose electronic structure is similar to that of the fragment CH with three nonbonding hybrid orbitals each containing one electron.

These organic and inorganic fragments, which have the same number of nonbonding hybrid orbitals occupied by the same number of electrons, are said to be isolobal analogues. ${ }^{4}$ The isolobal analogy is represented by a double-headed arrow, underneath which is a symbol that can be described either as a small hybrid orbital or as a tear.s

$$
\left.\begin{array}{l}
\mathrm{CH}_{3} \longleftrightarrow 0 \\
\mathrm{CH}_{2} \longleftrightarrow 0
\end{array} \mathrm{CrL}_{5}\right]^{-}
$$

${ }^{6}$ In the isolobal analogy, fragments (organic or inorganic) are considered with geometries which are not necessarily the most stable. For example, the most stable structure for $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$ is trigonal-bipyramidal (TBP) rather than square-pyramidal. Similarly, the $\mathrm{CH}_{3}^{+}$cation is planar.

Our choice of a chromium complex was quite arbitrary. In each of the inorganic fragments above, we may obviously replace chromium by molybdenum or tungsten, as they have the same number of valence electrons. More generally, we may also consider a metal from another group in the periodic table, so long as the number of electrons is not changed. Thus, all $d^{7} \mathrm{ML}_{5}$ complexes with a square-based pyramidal (SBP) geometry possess a singly occupied nonbonding hybrid orbital; they are therefore isolobal analogues of the pyramidal fragment $\mathrm{CH}_{3}$. As examples, we may consider the neutral complexes $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]$, $\left[\mathrm{Tc}(\mathrm{CO})_{5}\right]$, and $\left[\operatorname{Re}(\mathrm{CO})_{5}\right]$, or the cationic complex $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]^{+}$. In the same way, the complexes $\left[\mathrm{Mn}(\mathrm{CO})_{4}\right]^{-}$or $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$, with a 'butterfly' geometry, are, like $\left[\mathrm{CrL}_{4}\right]^{2-}$, complexes with a $d^{8}$ electronic configuration and therefore isolobal analogues of the organic fragment $\mathrm{CH}_{2}$. And pyramidal $d^{9} \mathrm{ML}_{3}$ complexes $\left(\left[\mathrm{Cr}(\mathrm{CO})_{3}\right]^{3-},\left[\mathrm{Mn}(\mathrm{CO})_{3}\right]^{2-}\right.$, $\left[\mathrm{Fe}(\mathrm{CO})_{3}\right]^{-}$, and $\left[\mathrm{Co}(\mathrm{CO})_{3}\right]$, for example) are all isolobal analogues of CH .


The same type of change may also be made on the organic fragment, by replacing $\mathrm{CH}_{3}$ with $\mathrm{BH}_{3}^{-}$or $\mathrm{NH}_{3}^{+}$(still with a pyramidal geometry). All these fragments are isolobal to $d^{7} M L_{5}$ complexes such as $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]$. One can also modify the charge on the organic or inorganic fragment. Since $\mathrm{CH}_{3}$ is isolobal to $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right], \mathrm{CH}_{3}^{+}$(pyramidal) is isolobal to $\left[\mathrm{Cr}(\mathrm{CO})_{5}\right]$, and $\mathrm{CH}_{3}^{-}$to $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$, if all these complexes are considered with a SBP geometry. ${ }^{6}$ In these three pairs of analogous compounds, the single nonbonding hybrid orbital contains zero $\left(\mathrm{CH}_{3}^{+}\right.$ and $\left.\left[\mathrm{Cr}(\mathrm{CO})_{5}\right]\right)$, one $\left(\mathrm{CH}_{3}\right.$ and $\left.\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]\right)$ or two electrons $\left(\mathrm{CH}_{3}^{-}\right.$ and $\left.\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]\right)$.

### 5.1.2. Fragment molecular orbitals

The analogy between the electronic structures of two isolobal fragments is, of course, still present if we consider their molecular orbitals.

We start with the organic series $\mathrm{CH}_{3}, \mathrm{CH}_{2}$, and CH . Just as there were one, two, or three nonbonding hybrid orbitals, respectively, in the valence-bond model, the electronic structure in terms of
molecular orbitals shows the presence of one, two or three nonbonding or essentially nonbonding MO, respectively (Chapter 1, Figures 1.4-1.6). Moreover, the number of electrons to be placed in these orbitals is the same as in the nonbonding hybrid orbitals: one for $\mathrm{CH}_{3}(5-7 \mathrm{a})$, two for $\mathrm{CH}_{2}(\mathbf{5 - 7 b})$, and three for $\mathrm{CH}(5-7 \mathbf{c})$. Notice that the electronic occupation of these MO can sometimes be problematical, as the orbital energies are quite close. For $\mathrm{CH}_{2}$ and CH , we have chosen to place the electrons in the MO so as to create as many pairs as possible. This electronic occupation is indeed the ground state of CH , but it is not for methylene, $\mathrm{CH}_{2}$, for which the ground state is a triplet, with one electron in each of the two nonbonding orbitals (parallel spins). In fact, this is not important when one is using the isolobal analogy. The capacity of a fragment to form new bonds depends more on the number of available nonbonding orbitals and the total number of electrons they contain, than on the more-or-less arbitrary initial distribution of the electrons in the fragment.


In contrast to hybrid orbitals, MO are adapted to molecular symmetry. Several features therefore arise that are new, compared to the previous description. First, even though the MO that are considered are always concentrated in the same region of space as the vacant site(s) of the tetrahedron, they do not always point directly towards these sites, except, of course, in the case of $\mathrm{CH}_{3}$ where there is only a single vacant site. For example, the shapes of the $\mathrm{a}_{1}$ and $\mathrm{b}_{2} \mathrm{MO}$ in the $\mathrm{CH}_{2}$ fragment mean that the two sites of the initial tetrahedron are equivalent for these two MO (5-7b), but neither orbital points along the direction of the
broken $\mathrm{C}-\mathrm{H}$ bonds. Moreover, in contrast to hybrid orbitals, the MO do not have the same shape or energy. Still using $\mathrm{CH}_{2}$ as an example, the $\mathrm{a}_{1}$ (a hybrid orbital) and $\mathrm{b}_{2}$ (a pure $p$ orbital) MO have different shapes, and the $a_{1}$ orbital is lower in energy than the $b_{2}$ orbital ( $\mathbf{5 - 7 b}$ ). In $C H$, there is one hybrid orbital with $\sigma$ symmetry and two pure $p$ orbitals with $\pi$ symmetry, the pair of degenerate orbitals ( $\pi$ ) being higher in energy than the $\sigma$ orbital (5-7c).

We shall now consider the inorganic fragments that are the isolobal analogues of $\mathrm{CH}_{3}, \mathrm{CH}_{2}$, and CH , that is, the SBP $d^{7} \mathrm{ML}_{5}$, 'butterfly' $d^{8} \mathrm{ML}_{4}$ and pyramidal $d^{9} \mathrm{ML}_{3}$ fragments (§ 5.1.1), and compare the MO that are concentrated around the vacant site(s) following the homolytic rupture of one, two, or three bonds in the initial structure (tetrahedral $\mathrm{CH}_{4}$ or octahedral $\mathrm{ML}_{6}$ ).

In the $\mathrm{CH}_{3}$ and $d^{7} \mathrm{ML}_{5}$ fragments, there is only one orbital of this type. It is totally symmetric ( $a_{1}$ symmetry) in both cases ( $C_{3 v}$ for $\mathrm{CH}_{3}$ and $C_{4 v}$ for $\mathrm{ML}_{5}$ ) (5-8), it points towards the vacant site (of the tetrahedron or of the octahedron) and of course it has axial symmetry with respect to the main symmetry axis ( $C_{3}$ and $C_{4}$, respectively). These two orbitals are not identical-the first is an $s p^{3}$ hybrid orbital, the second a polarized $z^{2}$ orbital-but they are very similar. Counting electrons, the nonbonding orbital of the methyl radical is singly occupied (5-7a). In a $d^{7} \mathrm{ML}_{5}$ fragment, such as $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]$, six electrons occupy the nonbonding orbitals derived from the $\mathrm{t}_{2 g}$ block of the octahedron, and the seventh occupies the polarized $z^{2}$ orbital (Chapter 2, Figure 2.7). In both fragments, the orbital pointing towards the vacant site is therefore singly occupied.


## 5-8

The set of all these characteristics-the same number of orbitals to be considered, similar symmetry properties, and the same number of electrons-allows us to establish that two fragments are isolobal in the framework of MO theory.
${ }^{7}$ This orbital $\left(p_{y}\right)$ has $\mathrm{b}_{2}$ symmetry (see the character table for the $C_{2 v}$ point group, Chapter 6, Table 6.5). If the $x$ - and $y$-axes are interchanged (the $z$-axis stays as the $C_{2}$ axis), this orbital becomes $p_{x}$, which has $\mathrm{b}_{1}$ symmetry (as was noted in Chapter 1, Figure 1.5). The choice is arbitrary. The important point is that this orbital is antisymmetric with respect to the molecular plane. This comment also applies to the polarized $d$ orbital on the metallic fragment, whose symmetry is either $\mathrm{b}_{1}$ or $\mathrm{b}_{2}$, depending on the orientation chosen for the $x$ - and $y$-axes.

## Comment

The hydrogen atom, with its atomic orbital $1_{s_{H}}$ singly occupied, is also an isolobal analogue of the $d^{7} \mathrm{ML}_{5}$ fragment.

The $\mathrm{CH}_{2}$ and $d^{8} \mathrm{ML}_{4}$ ('butterfly') fragments, both with $C_{2} v$ symmetry, provide a second illustration. Now there are two orbitals to consider for each fragment. For $\mathrm{CH}_{2}$, these are the two orbitals shown in $\mathbf{5 - 7 b}$ : the first, totally symmetric $\left(a_{1}\right)$, has of course axial symmetry with respect to the $C_{2}$ axis; the second (a pure $p$ orbital with $\mathrm{b}_{2}$ symmetry) ${ }^{7}$ is antisymmetric with respect to the molecular plane. For the metallic fragment, whose orbital structure was established in Chapter 2 (Figure 2.15), there is a polarized $d$ orbital, with $b_{2}$ symmetry, and the $s-p$ hybrid orbital on the metal, with $a_{1}$ symmetry. The orbitals involved in the two fragments have the same symmetries (5-9). Moreover, the same number of electrons must be placed in these orbitals: two in $\mathrm{CH}_{2}$ (5-7b) and also in a $d^{8} \mathrm{ML}_{4}$ fragment, since there are six electrons in the three nonbonding $d$ orbitals derived from the $t_{2 g}$ block of the octahedron (Chapter 2, Figure 2.15). Notice, however, that the energetic order of the orbitals is inverted (5-9), so that with the chosen electronic occupation (the lowest-energy MO doubly occupied) the electronic configuration is not the same for the two fragments $\left(\mathrm{a}_{1}^{2}\right.$ in $\mathrm{CH}_{2}, \mathrm{~b}_{2}^{2}$ in $\left.d^{8} \mathrm{ML}_{4}\right)$. As long as one is mainly concerned by the fragments' capacity to form new bonds, this difference is of no real consequence (see $\oint$ 5.3.1).


$\mathrm{CH}_{2} \longleftrightarrow \underset{0}{\longleftrightarrow} d^{8}-\mathrm{ML}_{4}$ ('butterfly')
5-9

For the CH and pyramidal $d^{9} \mathrm{ML}_{3}$ fragments, three orbitals must be considered. For CH , these are the orbitals shown in $\mathbf{5 - 7} \mathbf{c}$ : one orbital with $\sigma$ symmetry and two degenerate orbitals with $\pi$ symmetry, with three electrons in total. In a $d^{9} \mathrm{ML}_{3}$ complex, such as $\mathrm{Co}(\mathrm{CO})_{3}$, six electrons occupy the three nonbonding orbitals derived from the $\mathrm{t}_{2 g}$ block
of the octahedron (Chapter 2, Figure 2.13). At slightly higher energy, we find the three orbitals that are concentrated in the region of the vacant sites of the octahedron - a pair of degenerate orbitals (e symmetry in the $C_{3 v}$ point group) and one totally symmetric orbital ( $\mathrm{a}_{1}$ ) -in which we must place the three remaining electrons. If we compare the degenerate MO in the two fragments, we notice that in each case one orbital is antisymmetric with respect to the plane of the paper ( $p_{y}$ and $e_{y}$ ) but the other is symmetric with respect to this plane ( $p_{x}$ and $e_{x}$ ). Even though the symmetry labels in the $C_{\infty v}$ and $C_{3 v}$ point groups are different, the similarity between the orbitals in the two fragments is plain: the $\sigma$ orbital in CH corresponds to the $\mathrm{a}_{1}$ orbital in the metallic fragment, and the degenerate $\pi$ pair to the degenerate e pair (5-10). It is true that the energetic ordering of the degenerate and the totally symmetric orbitals is inverted, but this detail has no more importance here than in the preceding example concerning the $\mathrm{CH}_{2}$ and $d^{8} \mathrm{ML}_{4}$ fragments.


### 5.2. Other analogous fragments

An octahedral $d^{6} \mathrm{ML}_{6}$ complex, such as $\left[\mathrm{Cr}(\mathrm{CO})_{6}\right]$, is not the only initial structure that can be used to construct inorganic fragments which are analogous to the organic fragments $\mathrm{CH}_{3}, \mathrm{CH}_{2}$, and $\mathrm{CH}(5-11 \mathbf{a}$ and $\mathbf{b}$ ). Two other structures are often used as starting points.

The cyclopentadienyl ligand ( Cp ) is very widespread in organometallic chemistry. When it is coordinated in the $\eta^{5}$ mode, the most common situation, it behaves as an $\mathrm{L}_{2} \mathrm{X}$ ligand (Chapter 1, § 1.1.1.3). [CpML ${ }_{3}$ ] complexes ( $\mathrm{L}=\mathrm{CO}, \mathrm{PR}_{3}, \ldots$ ) are therefore pseudo-octahedral complexes of the $\mathrm{ML}_{5} \mathrm{X}$ type, which have 18 electrons when their electronic configuration is $d^{6}$. A $d^{6}\left[\mathrm{CpML}_{3}\right]$ complex may therefore be used as a new initial structure, very similar to that used previously. Since the presence of the cyclopentadienyl ligand oxidizes the metal by one unit, a metal with seven valence electrons ( $\mathrm{M}=\mathrm{Mn}, \mathrm{Tc}, \mathrm{Re}$ ) must be used to obtain a neutral complex. The complex $\left[\mathrm{CpMn}(\mathrm{CO})_{3}\right]$ is therefore a
possible starting point. Subsequently, the fragments $d^{7}\left[\mathrm{CpMn}(\mathrm{CO})_{2}\right]^{-}$, $d^{8}[\mathrm{CpMn}(\mathrm{CO})]^{2-}$, and $d^{9}[\mathrm{CpMn}]^{3-}$ are obtained by homolytic rupture of one, two, or three $\mathrm{Mn}-\mathrm{CO}$ bonds. Neutral metallic fragments can be obtained, if the metal is taken from one group further to the right in the periodic table each time. Starting from the complex $d^{6}\left[\mathrm{CpMn}(\mathrm{CO})_{3}\right]$, we can thus deduce that the fragments $d^{7}\left[\mathrm{CpFe}(\mathrm{CO})_{2}\right], d^{8}[\mathrm{CpCo}(\mathrm{CO})]$, and $d^{9}[\mathrm{CpNi}]$ are isolobal to $\mathrm{CH}_{3}, \mathrm{CH}_{2}$, and CH , respectively $(5-11 \mathrm{a}$ and $\mathbf{c}$ ).


Another starting point is sometimes used, taking a complex whose 'normal' number of electrons is 16 instead of 18 . The most common example involves a square-planar complex with a $d^{8}$ electronic configuration (Chapter 2, § 2.2). Starting, for example, from $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$, and using the same method as above, one can show that the fragments $\left[\mathrm{Co}(\mathrm{CO})_{3}\right]$, with a 'T-shaped' geometry, bent $\left[\mathrm{Ni}(\mathrm{CO})_{2}\right]$ and $[\mathrm{Cu}(\mathrm{CO})]$ are isolobal analogues of $\mathrm{CH}_{3}, \mathrm{CH}_{2}$, and CH , respectively (5-11a and d).

### 5.3. Applications

### 5.3.1. Metal-metal bonds

Since the $\mathrm{CH}_{3}$ and $d^{7} M L_{5}$ fragments have analogous electronic structures, their chemical behaviour is similar in some respects. For example,
dimerization to form ethane $\left(\mathrm{C}_{2} \mathrm{H}_{6}\right)$ is an important property of the methyl radical, as is its initiation of radical-type chain reactions. Similarly, species such as $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]$ or $\left[\mathrm{Co}(\mathrm{CN})_{5}\right]^{3-}$, both of which are $d^{7} \mathrm{ML}_{5}$ systems, have a very rich radical-type chemistry; in particular, they dimerize to form $\left[\mathrm{Mn}_{2}(\mathrm{CO})_{10}\right]$ and $\left[\mathrm{Co}_{2}(\mathrm{CN})_{10}\right]^{6-}$, respectively, (despite the presence, in the latter, of three negative charges on each monomer!). These bimetallic complexes are therefore isolobal analogues of the ethane molecule.


The main orbital interaction that occurs during the dimerization of the fragments $\left(\mathrm{CH}_{3},\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]\right.$ or $\left.\left[\mathrm{Co}(\mathrm{CN})_{5}\right]^{3-}\right)$ involves in all cases the singly occupied nonbonding orbital on each monomer. The interaction schemes for the orbitals on the $\left[\mathrm{Mn}(\mathrm{CO})_{5}\right]\left(\right.$ or $\left.\left[\mathrm{Co}(\mathrm{CN})_{5}\right]^{3-}\right)$ and $\mathrm{CH}_{3}$ fragments produce a doubly occupied bonding MO in the dimer, with axial symmetry around the internuclear axis, and the corresponding antibonding orbital which is empty ( $\mathbf{5 - 1 2 a}$ and $\mathbf{b}$ ). There is therefore a single $\sigma$ bond between the two metallic centres in the $\left[\mathrm{Mn}_{2}(\mathrm{CO})_{10}\right]$ or $\left[\mathrm{Co}_{2}(\mathrm{CN})_{10}\right]^{6-}$ dimers, which implies nearly free rotation around the metal-metal bond, as there is around the $\mathrm{C}-\mathrm{C}$ single bond in ethane. The analogy can be extended to mixed inorganic-organic dimers such as $\left[\mathrm{Mn}(\mathrm{CO})_{5} \mathrm{CH}_{3}\right]$, formed by the combination of the metallic fragment [ $\mathrm{Mn}(\mathrm{CO})_{5}$ ] and its organic analogue $\mathrm{CH}_{3}$. In this last example, there is once again a single $\sigma$ bond between the metal and the carbon, resulting from the combination of the singly occupied orbitals on the two fragments.


Another example involves the dimerization of the $\mathrm{CH}_{2}$ and 'butterfly' $d^{8}-\mathrm{ML}_{4}$ fragments. The dimer of $\mathrm{CH}_{2}$ is just simply ethylene, a planar molecule characterized by a double bond between the two carbon atoms (a $\sigma$ bond and a $\pi$ bond). The isolobal analogy linking the $\mathrm{CH}_{2}$ and 'butterfly' $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$ fragments $(5-9)$ shows us that the
bimetallic complex $\left[\mathrm{Fe}_{2}(\mathrm{CO})_{8}\right]$, in the conformation shown in $\mathbf{5 - 1 3}$, is an isolobal analogue of ethylene, which implies the existence of a double bond between the two metallic centres.


The similarity in the electronic structures of the two molecules is confirmed by the interaction diagrams shown in $\mathbf{5 - 1 4 a}$ for $\mathrm{C}_{2} \mathrm{H}_{4}$ and $\mathbf{5 - 1 4 b}$ for $\left[\mathrm{Fe}_{2}(\mathrm{CO})_{8}\right]$.



In ethylene, the interaction between the fragment orbitals leads to the formation of two bonding ( $\sigma_{\mathrm{CC}}$ and $\pi_{\mathrm{CC}}$ ) and two antibonding $\mathrm{MO}\left(\sigma_{\mathrm{CC}}^{*}\right.$ and $\left.\pi_{\mathrm{CC}}^{*}\right)$. In the ground electronic state, the former are doubly occupied and characterize the $\sigma$ and $\pi$ bonds ( $\mathbf{5 - 1 4 a}$ ). In the $\left[\mathrm{Fe}_{2}(\mathrm{CO})_{8}\right]$ dimer, there are also two bonding MO, one of $\sigma$ type with axial symmetry around the internuclear axis, the other of $\pi$ type with a nodal plane in the plane of the paper, and the two corresponding antibonding MO (5-14b). There are still, in total, four electrons to be placed in these MO, so in the ground electronic state the two bonding $\mathrm{MO}\left(\sigma_{\mathrm{MM}}\right.$ and $\left.\pi_{\mathrm{MM}}\right)$ are doubly occupied, just as in ethylene.

We noted above that the different orbital occupations in the two isolated fragments $\left(\mathrm{a}_{1}^{2}\right.$ for $\mathrm{CH}_{2}, \mathrm{~b}_{2}^{2}$ for $\left.\left[\mathrm{Fe}(\mathrm{CO})_{4}\right], 5-9\right)$ has no influence on their capacity to form new bonds and is therefore of little importance for the isolobal analogy. The example above illustrates this point nicely: independently of the energetic order of the fragment orbitals and the initial electronic occupation, the occupied and empty MO in the two dimers have analogous symmetry properties (5-15).


The analogy is not limited to molecules that are formed by the combination of two fragments. For example, the complex $\left[\mathrm{Os}_{3}(\mathrm{CO})_{12}\right]$ can be considered as the combination of three $\left[\mathrm{Os}(\mathrm{CO})_{4}\right]$ fragments with $C_{2 v}$ geometry. Each of these is of course isoelectronic to $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$, and therefore isolobal to the organic fragment $\mathrm{CH}_{2}$. The $\left[\mathrm{Os}_{3}(\mathrm{CO})_{12}\right]$ complex is therefore an isolobal analogue of cyclopropane (5-16). Each metallic centre is bound to two neighbours by a single bond, as in
the organic analogue. Other complexes, which appear to be very different, also have this property. The fragment $[\mathrm{CpRhL}]$ is isoelectronic to $[\mathrm{CpCoL}]$, and therefore isolobal to the $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$ and $\mathrm{CH}_{2}$ fragments (5-11). As a result, the complex $\left[\mathrm{Cp}_{2} \mathrm{Rh}_{2}(\mathrm{CO})_{2}\left(\mu-\mathrm{CH}_{2}\right)\right]$ can be described as a 'two-thirds inorganic' cyclopropane (5-16).


5-16

### 5.3.2. Conformational problems

Several conformational properties that are well known in organic molecules reappear in inorganic complexes or in mixed organic/ inorganic systems. The ethylene molecule provides a good example; it adopts a planar geometry that allows the formation of two bonds between the carbon atoms (a $\sigma$ bond and a $\pi$ bond). In the 'orthogonal' structure, where the two methylene groups lie in perpendicular planes, only the $\sigma$ bond would be preserved. We now turn to the carbene complex $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\mathrm{CH}_{2}\right)\right]$ in a TBP geometry, with the carbene in an equatorial position. Since the $\mathrm{CH}_{2}$ and $\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]$ fragments are isolobal, this complex is an isolobal analogue of ethylene, $\mathrm{C}_{2} \mathrm{H}_{4}$. Two limiting conformations can be imagined, with the $\mathrm{CH}_{2}$ ligand either perpendicular to the equatorial plane of the complex, or in this plane. The analogue of the first of these conformations is indeed a planar molecule of ethylene ( $\mathbf{5 - 1 7} \mathbf{a}$ ), but the second corresponds to a strongly destabilized state of ethylene in the orthogonal geometry ( $\mathbf{5 - 1 7 b}$ ). However,

the similarity in electronic structures does not necessarily imply that the barriers to rotation ( $\mathbf{5 - 1 7 a} \rightarrow \mathbf{b}$ ) will be comparable in the organic molecule and in the organometallic complex. About $64 \mathrm{kcal} / \mathrm{mol}$ are needed to twist ethylene from its planar to its orthogonal structure, but the rotation of the carbene ligand in a complex of $d^{8}\left[\mathrm{ML}_{4}\left(\mathrm{CR}_{2}\right)\right]$ type needs only about $15-20 \mathrm{kcal} / \mathrm{mol}$ (see Exercise 5.8).

In the same way, we can predict (or rediscover) the conformational preferences of several ethylene complexes. If we replace the carbene in the preceding example by ethylene, still in an equatorial position, we obtain the complex $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\mathrm{C}_{2} \mathrm{H}_{4}\right)\right]$ which is isolobal to cyclopropane. The conformation shown in $\mathbf{5 - 1 8 a}$ is isolobal to cyclopropane in its most stable geometry, with three carbon atoms on a tetrahedral environment, whereas that shown in $\mathbf{5 - 1 8 b}$ is isolobal to a 'deformed' cyclopropane, in which the coordination around one of the carbon atoms is squareplanar. The first of these conformations is therefore more stable: in contrast to the carbene, ethylene 'prefers' to lie in the equatorial plane of the complex. In a simple way, we come to the same conclusion concerning the conformational preference as we did in the preceding chapter (§ 4.1.1), when we examined the problem starting from the interaction diagrams between fragment orbitals.



### 5.4. Limitations

The isolobal analogy draws out similarities in the electronic structures of organic and inorganic molecules that can easily escape our attention. As a result of these similarities, we can discover resemblances in the number and nature of bonds, molecular geometries, and sometimes even reactivities. But analogies must not be pushed too far, and there can be significant differences between analogous molecules.

One of the limitations concerns the kinetic stabilities of analogous species. For example, in the analogous series of compounds ethylene, carbene complexes of iron tetracarbonyl and its dimer, $\left[\mathrm{Fe}_{2}(\mathrm{CO})_{8}\right]$ (5-19), the first two are known, but the last is unstable and to date it
has only been observed in an inert matrix at very low temperatures. In other words, the analogy with a very stable organic molecule is not a guarantee that the inorganic species will be stable.


Another limitation concerns the capacity of some ligands, such as the carbonyl group, to coordinate in either terminal or bridging positions ( $\mu$-type coordination) in polynuclear complexes. Consider the bimetallic complex $\left[\mathrm{CpFe}(\mathrm{CO})_{2}\right]_{2}$, for example, $(\mathbf{5 - 2 0 b})$. Each $\left[\mathrm{CpFe}(\mathrm{CO})_{2}\right]$ fragment is analogous to the methyl radical (5-11), so the complex is isolobal to ethane (5-20a). However, it is known in two different forms: $\left[\mathrm{CpFe}(\mathrm{CO})_{2}\right]_{2}$, where all the carbonyl ligands are in terminal positions, and $[\mathrm{CpFe}(\mathrm{CO})(\mu-\mathrm{CO})]_{2}$, in which two carbonyl ligands are bridging $(5-20 c)$, and this latter form is the more stable. A structure of this type, with bridging hydrogen atoms, is of course unknown for ethane, the organic analogue of this complex.


In metallic complexes, the bridged and non-bridged structures are often very close in energy, and small changes can favour one geometry or the other. Thus, in the trimetallic complex $\left[\mathrm{Os}_{3}(\mathrm{CO})_{12}\right]$, which is analogous to cyclopropane (5-21a), all the carbonyl ligands are terminal. However, in the isoelectronic iron complex, two carbonyl


5-21a


5-21b
ligands bridge the two metallic centres, giving the isomer $\left[\mathrm{Fe}(\mathrm{CO})_{4}\left(\mathrm{Fe}(\mathrm{CO})_{3}(\mu-\mathrm{CO})\right)_{2}\right](5-21 \mathbf{b})$. As in the previous example, a
bridged structure of this type cannot be seriously imagined for the organic analogue, but that does not prevent the complex $\mathbf{5 - 2 1 b}$, like the complex 5-21a, from being an isolobal analogue of cyclopropane.

## Exercises

## 5.1

1. Give an organic and a bimetallic isolobal analogue of the complex $\left[\mathrm{Mn}(\mathrm{CO})_{5} \mathrm{CH}_{3}\right]$.
2. Suggest an orbital interaction diagram that is suitable for the description of the bond between the metallic centre and the methyl group in this mixed complex.

## 5.2

1. By using the isolobal analogy with an organic molecule, predict the number and nature ( $\sigma$ or $\pi$ ) of the metal-metal bond(s) in the bimetallic complex $[\mathrm{CpRh}(\mathrm{CO})]_{2}$.
2. Suggest a structure for this compound.

## 5.3

1. How many carbonyl ligands ( $n$ ) are necessary on each metal centre for the bimetallic complex of osmium (2) to be an isolobal analogue of cyclobutane (1)?


1
2
2. Indicate the geometrical arrangement of the ligands around the metal centres.

## 5.4

1. Show that the trimetallic complex $\left[(\mathrm{CO})_{4} \mathrm{Fe}(\mathrm{Pt}(\mathrm{CO}) \mathrm{L})_{2}\right]$ $\left(\mathrm{L}=\mathrm{PR}_{3}\right)$ is an isolobal analogue of cyclopropane.

2. Deduce the positions of the ligands around each metallic centre.

## 5.5

In the complex below, $\left[\mathrm{Cr}(\mathrm{CO})_{4}\left(\mathrm{BH}_{4}\right)\right]^{-}$, the borohydride ligand is bound to the metal centre in the $\eta^{2}$ mode, that is, with two bridging hydrogen atoms.


1. Give an organic isolobal analogue of the fragment $\left[\mathrm{Cr}(\mathrm{CO})_{4}\right]^{-}$.
2. Give a boron-based analogue of this fragment.
3. Which boron-based compound is the isolobal analogue of the full complex?
Deduce (or rediscover) its geometrical structure.

## 5.6

Use the isolobal analogy to predict the more stable conformations for the following ethylene complexes ( $\mathrm{Cp}=\eta^{5}$-cyclopentadienyl).

(a)

(a)

or
or

(b)

(b)
5.7

1. Show that the complex $\left[\operatorname{Ir}_{4}(\mathrm{CO})_{12}\right](\mathbf{1})$ is an isolobal analogue of tetrahedrane (2).


1


2
2. What can be said about the complex $\left[\left(\eta^{3}\right.\right.$-cyclopropenyl) $\left.\mathrm{Co}(\mathrm{CO})_{3}\right]$ 3 and the bimetallic complex $\left[\mu-\left(\eta^{2}\right.\right.$-alcyne $\left.)\left(\mathrm{Co}(\mathrm{CO})_{3}\right)_{2}\right]$ 4 , in which the acetylenic bond is perpendicular to the metalmetal bond?


3


4
3. Give another limiting representation for the structures of complexes 3 and 4.

## 5.8

Examine the analogies described in Scheme 5-17. Does the doublebond really disappear on passing from structure $\mathbf{a}$ to structure $\mathbf{b}$ for (i) the organic molecule and (ii) the organometallic complex?

## Elements of group theory and applications

Examination of the geometries of isolated molecules shows that there are four types of symmetry elements: reflection planes, axes of rotation, inversion centres, and improper axes of rotation. A symmetry operation moves a molecule from an initial configuration to another, equivalent configuration, either by leaving the position of the atoms unchanged, or by exchanging equivalent atoms. The ideas of symmetry element and symmetry operation are of course very closely linked, since a symmetry operation is defined with respect to a given element of symmetry, and, conversely, the presence of a symmetry element is established by the presence of one or more symmetry operations that are associated with that element.


6-1

### 6.1. Symmetry elements and symmetry operations

### 6.1.1. Reflection planes

Consider the dimethylether molecule $\left(\mathrm{O}_{\left.\left(\mathrm{CH}_{3}\right)_{2}\right) \text { in the conformation }}\right.$ shown in 6-1: the hydrogen atoms $\mathrm{H}_{1}$ and $\mathrm{H}_{4}$ lie in the plane $\mathrm{C}_{1}-\mathrm{O}-\mathrm{C}_{2}$ (the plane of the paper), but the atoms $\mathrm{H}_{2}, \mathrm{H}_{5}$ and $\mathrm{H}_{3}, \mathrm{H}_{6}$ are 'above' and 'below' this plane.

The $x z$ plane, which is perpendicular to the plane of the paper and which bisects the angle $\mathrm{C}_{1}-\mathrm{O}-\mathrm{C}_{2}$, is a symmetry element of the molecule, written $\sigma_{v}$. Reflection in this plane leads to a configuration that is equivalent to the initial one: the position of the oxygen atom (located in the $\sigma_{v}$ plane) stays unchanged, and the positions of the pairs of equivalent atoms $\left(\mathrm{C}_{1}, \mathrm{C}_{2}\right),\left(\mathrm{H}_{1}, \mathrm{H}_{4}\right),\left(\mathrm{H}_{2}, \mathrm{H}_{5}\right)$, and $\left(\mathrm{H}_{3}, \mathrm{H}_{6}\right)$ are interchanged (6-2).


The $y z$ plane (written $\sigma_{v}{ }^{\prime}$ ), which contains the three heavy atoms, is also a symmetry elements of the molecule, since reflection in this plane maintains the positions of the atoms $\mathrm{O}, \mathrm{C}_{1}, \mathrm{C}_{2}, \mathrm{H}_{1}$, and $\mathrm{H}_{4}$, and exchanges the pairs of hydrogen atoms $\left(\mathrm{H}_{2}, \mathrm{H}_{3}\right)$ and $\left(\mathrm{H}_{5}, \mathrm{H}_{6}\right)(\mathbf{6 - 3})$.


6-3

If the operation of reflection in either of these planes is performed twice (an operation which is written $\sigma^{2}$ ), a configuration is obtained that is identical to the initial one. This idea can be expressed as follows: $\sigma^{2}=$ $E$, where $E$ represents the identity operation. A reflection plane is therefore associated with a single symmetry operation that leads to a configuration that is equivalent to but not identical with the initial configuration.

### 6.1.2. Inversion centre

If the change of the coordinates $(x, y, z)$ for every atom to $(-x,-y,-z)$ leads to a molecular configuration that is equivalent to the initial one, the point at the origin of the frame of reference is an inversion centre (or centre of symmetry) for the molecule. The symbol $i$ is used both for this symmetry element and for the associated operation. Two examples will illustrate the fact that the inversion centre may be located on one of the atoms of the molecule, but this is not necessary. In the octahedral complex $\left[\mathrm{FeCl}_{6}\right]^{4-}$, the iron atom is situated at the inversion centre: the inversion operation leaves its position unchanged, whereas the chlorine atoms are exchanged in pairs: $\mathrm{Cl}_{1,2}, \mathrm{Cl}_{3,4}$, and $\mathrm{Cl}_{5,6}(6-4)$. But in the ethylene molecule, the inversion centre is situated in the middle of the carbon-carbon bond, and all the atoms are exchanged by the inversion operation (6-5).



6-5

Notice that a molecule may not contain more than one inversion centre, and that there is only one symmetry operation associated with this element. If the inversion operation is performed twice $\left(i^{2}\right)$, the resulting molecular configuration is identical to the initial one $\left(i^{2}=E\right)$.

### 6.1.3. Rotation axes

We return to the example of dimethylether $(\mathbf{6}-\mathbf{1})$. A rotation of $180^{\circ}(\pi)$ around the $z$-axis leads to a molecular configuration that is equivalent to the initial one (6-6), with the following pairs of atoms being exchanged: $\left(\mathrm{C}_{1}, \mathrm{C}_{2}\right),\left(\mathrm{H}_{1}, \mathrm{H}_{4}\right),\left(\mathrm{H}_{2}, \mathrm{H}_{6}\right)$, and $\left(\mathrm{H}_{3}, \mathrm{H}_{5}\right)$. The $z$-axis is therefore a symmetry element of the molecule. The ratio between $360^{\circ}$ and the rotation angle associated with the symmetry operation $\left(180^{\circ}\right.$ in this example) defines the order of the axis. In this case it is therefore a twofold axis (or of order two), written $C_{2}$. A single symmetry operation is associated with this axis, shown in 6-6 and written $C_{2}^{1}$ or simply $C_{2}$. If this operation is applied twice (a rotation of $2 \times 180^{\circ}=360^{\circ}$, written $C_{2}^{2}$ ), we return to the initial molecular configuration $\left(C_{2}^{2}=E\right)$. Note that by convention, the rotation is performed in a clockwise sense.


6-6


6-7

Ammonia $\left(\mathrm{NH}_{3}\right)$ is an example of a molecule that possesses a threefold axis $\left(C_{3}\right)$. This is the $z$-axis that passes through the N atom and is perpendicular to the equilateral triangle defined by the three hydrogen atoms (6-7).

A rotation of $120^{\circ}($ or $1 \times 2 \pi / 3)$ around this axis leads to an equivalent configuration in which $\mathrm{H}_{1}, \mathrm{H}_{2}$, and $\mathrm{H}_{3}$ have been replaced by $\mathrm{H}_{2}, \mathrm{H}_{3}$, and $\mathrm{H}_{1}$, respectively ( $\mathbf{6 - 8}$ ). This operation is written $C_{3}^{1}$ or $C_{3}$. A rotation of $240^{\circ}$ (or $2 \times 2 \pi / 3$ ) also leads to a configuration that is equivalent to the initial one $\left(\left(\mathrm{H}_{1}, \mathrm{H}_{2}, \mathrm{H}_{3}\right)\right.$, replaced by $\left(\mathrm{H}_{3}, \mathrm{H}_{1}\right.$, and $\left.\left.\mathrm{H}_{2}\right)\right)$, (6-9). This new operation is written $C_{3}^{2}$. A rotation of $360^{\circ}$ (or $3 \times 2 \pi / 3$ ) brings us
back to the initial configuration, that is, $C_{3}^{3}=E$. There are therefore two symmetry operations, different from the identity operation, that are associated with a $C_{3}$-axis.




6-8






6-9

More generally, an axis of order $n\left(C_{n}\right)$ is present if a rotation of $360^{\circ} / n$ (or $2 \pi / n$ ) leads to an equivalent configuration for the molecule. This axis generates $(n-1)$ symmetry operations, $C_{n}^{1}, C_{n}^{2}, \ldots, C_{n}^{n-1}$, the operation $C_{n}^{n}$ being equivalent to the identity operation ( $C_{n}^{n}=E$ ).

Some molecules have several rotation axes. The axis of highest order is called the principal axis. The complex $\left[\mathrm{PtCl}_{4}\right]^{2-}$, in which the platinum atom is located in the centre of the square defined by the four chlorine atoms $(\mathbf{6}-\mathbf{1 0})$, possesses a $C_{4}$-axis, perpendicular to the plane of the square, and four $C_{2}$-axes in the plane of the square: two of these are co-linear with the bonds $\mathrm{Cl}_{1}-\mathrm{Pt}-\mathrm{Cl}_{3}$ and $\mathrm{Cl}_{2}-\mathrm{Pt}-\mathrm{Cl}_{4}\left(C_{2}^{\prime}\right)$, and the two others bisect the angles $\mathrm{Cl}-\mathrm{Pt}-\mathrm{Cl}\left(C_{2}^{\prime \prime}\right)$. The principal axis is therefore four-fold (of order four). The existence of a four-fold rotation axis implies the presence of a co-linear two-fold axis, as the operation $C_{4}^{2}$ (a rotation of $2 \times 2 \pi / 4$ ) is identical to the operation $C_{2}^{1}$ (a rotation of $1 \times 2 \pi / 2$ ).

The set of operations associated with a rotation around the axis that is perpendicular to the molecular plane can thus be written: $C_{4}^{1}$, $C_{4}^{2}\left(=C_{2}^{1}\right), C_{4}^{3}$ and $C_{4}^{4}(=E)$. Usually, the notation used for these is $C_{4}^{1}$, $C_{2}, C_{4}^{3}$, and $E$, or sometimes $2 C_{4}, C_{2}$, and $E$, where the two operations $C_{4}^{1}$ and $C_{4}^{3}$ are grouped together. This type of situation occurs every time that there is an axis whose order is even and greater than 2 (in practice, 4,6 , or 8 ). For example, if a molecule possesses a $C_{6}$-axis, a $C_{3}$-axis and a $C_{2}$-axis are necessarily co-linear with this axis. The symmetry operations are thus: $C_{6}^{1}, C_{6}^{2}\left(=C_{3}^{1}\right), C_{6}^{3}\left(=C_{2}^{1}\right)$, $C_{6}^{4}\left(=C_{3}^{2}\right), C_{6}^{5}$ and $C_{6}^{6}(=E)$, which can also be written $2 C_{6}, 2 C_{3}, C_{2}$, and $E$.

## Comment

The $\left[\mathrm{PtCl}_{4}\right]^{2-}$ example ( $\mathbf{6 - 1 0}$ ) allows us to make several ideas more precise about planes of symmetry. The $z$-axis defines the 'vertical' direction. The molecular plane ( $x y$ ), which is perpendicular to the principal ( $z$ ) axis, is written $\sigma_{h}$ ('h' for 'horizontal'). The two planes which contain the principal axis and two $\mathrm{Pt}-\mathrm{Cl}$ bonds are written $\sigma_{v}$ ('v' for 'vertical'). And the two planes $x z$ and $y z$, which contain the principal axis and which bisect the angles $\mathrm{Cl}-\mathrm{Pt}-\mathrm{Cl}$, are written $\sigma_{d}$ ('d' for 'dihedral').

### 6.1.4. Improper rotation axes

The symmetry operation associated with an improper axis of order $n$ (written $S_{n}$ ) is performed in two steps: a rotation of $2 \pi / n$ around this axis, followed by a reflection in a plane that is perpendicular to this axis. It is important to realize that the axis used is not necessarily a $C_{n}$-axis, nor is the plane that is involved necessarily a symmetry element of the molecule.



6-12

As a first example, we consider the methane molecule $\left(\mathrm{CH}_{4}\right)$, whose geometry is tetrahedral. If we perform a rotation of $90^{\circ}(2 \pi / 4)$ around the axis bisecting the angles $\mathrm{H}_{1}-\mathrm{C}-\mathrm{H}_{2}$ and $\mathrm{H}_{3}-\mathrm{C}-\mathrm{H}_{4}$, followed by a reflection in the plane perpendicular to this axis passing through the carbon atom, we obtain a configuration that is equivalent (but not identical) to the initial one ( $\mathbf{6 - 1 1}$ ). The axis which we used is therefore an improper four-fold axis $\left(S_{4}\right)$, and the symmetry operation described in 6-11 is written $S_{4}^{1}$, or simply $S_{4}$. Notice that this axis is not a $C_{4}$-axis in the molecule (but a $C_{2}$-axis), and that the plane used is not a symmetry element in the molecule. As in the cases of $C_{n}$ axes, the $S_{4}$ operation can be performed $m$ times in succession, and these are written $S_{n}^{m}$. In the example above, it is easy to verify that $S_{4}^{2}=C_{2}^{1}$ (a two-fold axis, co-linear with $S_{4}$ ) and that $S_{4}^{4}=E$ (see Exercise 6.4). The only two operations that are unique to the $S_{4}$-axis are therefore $S_{4}^{1}$ and $S_{4}^{3}$.

The hypervalent molecule $\mathrm{PF}_{5}$ adopts a trigonal-bipyramidal (TBP) geometry (6-12). The fluorine atoms $\mathrm{F}_{1}, \mathrm{~F}_{2}$, and $\mathrm{F}_{3}$ define the base of the bipyramid (the $x y$ plane, written $\sigma_{h}$ ) and are situated at the vertices of an equilateral triangle (the angles between bonds are $120^{\circ}$ ). The atoms
$\mathrm{F}_{4}$ and $\mathrm{F}_{5}$ define the vertices of the bipyramid, the $\mathrm{F}_{4}-\mathrm{P}-\mathrm{F}_{5}$ axis (the $z$-axis) being perpendicular to the plane $\sigma_{h}$.

A rotation of $120^{\circ}(2 \pi / 3)$ around the $z$-axis, followed by a reflection in the $x y$ plane, leads to an equivalent configuration for the molecule $(6-13)$. The $z$-axis is therefore an improper three-fold axis $\left(S_{3}\right)$. Notice that in this example both the axis and the plane that were used to perform the $S_{3}$ operation are themselves symmetry elements in the molecule: the $z$-axis is also a $C_{3}$-axis, and the $\sigma_{h}$ plane is a reflection plane. A characteristic property of the $S_{3}$-axis (and, more generally, of improper axes whose order is odd) is that the $S_{3}^{3}$ operation does not lead to a configuration that is identical to the initial one ( $S_{3}^{3}$ is different from $E)$. Although a rotation of $3 \times 2 \pi / 3\left(=360^{\circ}\right)$ around the $z$-axis does leave the positions of the five fluorine atoms unchanged, an odd number of reflections (three) in the $x y$ plane interchanges the positions of the atoms $\mathrm{F}_{4}$ and $\mathrm{F}_{5}\left(\sigma_{h}^{3}=\sigma_{h}\right)$. We therefore obtain $S_{3}^{3}=\sigma_{h}$. To obtain the identity operation, the operation $S_{3}^{6}$ must be performed (a rotation of twice $360^{\circ}$ and an even number (six) of reflections in the plane). Of the six $S_{3}^{m}(m=1-6)$ that seem possible, only the $S_{3}^{1}$ and $S_{3}^{5}$ operation need to be considered, as the others are equivalent to symmetry operations associated with the $C_{3}$-axis the $\sigma_{h}$ plane or the identity operation in the case of $S_{3}^{6}$.


### 6.2. Symmetry groups

### 6.2.1. Definitions

The set of symmetry operations associated with a molecule makes up a group, in the mathematical sense of the term. One can indeed verify $(i)$ If $A, B$, and $C$ are three symmetry operations, the relation $A(B C)=(A B) C$ holds, that is, the combination of symmetry operations is associative, (ii) there is a null element (the identity operation $E$ ): $A E=E A=\mathrm{A}$, (iii) for each symmetry operation $A$, there is an inverse operation $A^{-1}$ such that $A^{-1} A=A A^{-1}=E$ (see Exercise 6.5).

Moreover, all the symmetry elements intersect at a single point (e.g. the oxygen atom in dimethylether (6-1), the nitrogen atom in ammonia (6-7), the middle of the carbon-carbon bond in ethylene $(6-5)$ ). The expression point-group symmetry is therefore used. For each molecule, there is a corresponding point group that completely characterizes its symmetry properties.

The different symmetry operations can be grouped into classes. From a mathematical point of view, two symmetry operations $A$ and $B$ belong to the same class if there is a symmetry operation $X$ in the group, such that $B=X^{-1} A X$. Verification of this relationship can be tedious, and it will be sufficient for us to note that a class gathers together symmetry operations of the same 'nature', even if this term is rather vague: as examples, we may give the rotations $C_{3}^{1}$ and $C_{3}^{2}$ associated with a $C_{3}$-axis, the rotations $C_{4}^{1}$ and $C_{4}^{3}$ associated with a $C_{4}$-axis, and the reflections in two $\sigma_{v}$ planes or in two $\sigma_{d}$ planes in $\left[\mathrm{PtCl}_{4}\right]^{2-}$ (see the comment in § 6.1.3). As far as notation is concerned, one often groups together the operations in a given class, for example, replacing $C_{3}^{1}$ and $C_{3}^{2}$ by $2 C_{3}, C_{4}^{1}$ and $C_{4}^{3}$ by $2 C_{4}$, etc.

### 6.2.2. Determination of the symmetry point group

A set of symmetry elements (or of symmetry operations) enables us to define a symmetry point group which is represented by a symbol. There is a simple method to determine the point group, and it is not necessary to establish a complete list of all the symmetry operations in the molecule under consideration.

1. Initially, check whether the molecule adopts one of the four most easily identifiable geometries: octahedral (e.g. $\left[\mathrm{FeCl}_{6}\right]^{4-}(6-4)$ ), tetrahedral (e.g. $\mathrm{CH}_{4}$ ), linear with an inversion centre (e.g. $\mathrm{CO}_{2}$ ), or linear without an inversion centre (e.g. HCN ). If this is the case, the problem is solved; the symbols associated with these four point groups are $O_{h}, T_{d}, D_{\infty h}$, and $C_{\infty v}$ respectively.
2. If none of these groups is appropriate, and if there is no rotation axis, there are three possibilities: the molecule (i) does not possess and symmetry element (the point group $C_{1}$, in which the identity operation is the only symmetry operation); (ii) possesses a plane of symmetry (the point group $C_{s}$ ); (iii) possesses an inversion centre (the point group $C_{i}$ ).
3. If the molecule possesses $a$ single rotation axis (of order $n$ ), there are four possibilities: (i) there is no other symmetry element besides this axis, with the possible exception of a co-linear axis of lower order. The point-group symbol is then $C_{n}$; (ii) the existence of a reflection plan $\sigma_{h}$ (perpendicular to the $C_{n}$-axis) indicates that the point group is of the type $C_{n h}$; (iii) if there are $n \sigma_{v}$ planes (that contain the $C_{n}$
axis), the point-group symbol is $C_{n v} ;$ (iv) the presence of an improper axis $S_{2 n}$, co-linear with the $C_{n}$-axis, indicates that the point group is of the type $S_{2 n}$.
4. If the molecule possesses several rotation axes (the principal axis being of order $n$ ), there are three possibilities: $(i)$ there are no other symmetry elements besides these axes ( $D_{n}$ point-groups); (ii) there are $n \sigma_{d}$ planes that bisect the $C_{2}$ axes ( $D_{n d}$ point groups); (iii) there is also a $\sigma_{h}$ plane, perpendicular to the principal axis ( $D_{n h}$ point groups).

### 6.2.3. Basis of an irreducible representation

The application of group theory to the construction of molecular orbitals leads us to study the way in which atomic orbitals (or linear combinations of these orbitals) transform when the operations of the point group are applied. Several definitions, which are also appropriate for other types of function besides orbitals, will be presented and illustrated by some simple examples.

### 6.2.3.1. Basis for a representation

If a set of functions $f=\left\{f_{1}, f_{2}, \ldots, f_{i}, \ldots, f_{n}\right\}$ is such that any symmetry operation, $R_{k}$, of the group $G$ transform one of the functions, $f_{i}$, into a linear combination of the various functions of the set $f$, the set is said to be globally stable and to constitute a basis for the representation of the group $G$. As the symmetry operations maintain the positions of the atoms or interchange the positions of equivalent atoms, it can be shown that the set of atomic orbitals (AO) of a molecule constitute a basis for the representation of the point-group symmetry of the molecule. In what follows, we shall adopt the usual notation in group theory, and indicate a basis for a representation by $\Gamma$.

### 6.2.3.2. Basis for an irreducible representation

Suppose that a basis $\Gamma$, of dimension $n$, can be decomposed into several bases $\Gamma_{i}$, whose dimensions are smaller $\left(n_{i}\right)$, each of which is globally stable with respect to all the symmetry operations of the group. Suppose also that it is not possible to decompose any of the representations $\Gamma_{i}$ into representations whose dimensions are smaller than $n_{i}$. The reducible representation $\Gamma$ is said to have been decomposed into a sum of irreducible representations $\Gamma_{i}$, which is written:

$$
\begin{equation*}
\Gamma=a_{1} \Gamma_{1} \oplus a_{2} \Gamma_{2} \oplus \cdots \oplus a_{m} \Gamma_{m} \tag{6.1}
\end{equation*}
$$

### 6.2.3.3. $\mathrm{H}_{2} \mathrm{O}$ as an example



6-14

To illustrate this point, we shall consider the set of valence $A O$ of the atoms in the water molecule: $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ on the hydrogen atoms and $2 s, 2 p_{x}, 2 p_{y}$, and $2 p_{z}$ on the oxygen atom. The water molecule possesses a two-fold rotation axis $(z)$, and two planes of symmetry, $x z$ and $y z(6-14)$; its point group is therefore $C_{2 v}$. The set of atomic orbitals constitutes a basis $(\Gamma)$ for the representation of this point group.

We apply one of the symmetry operations of the group $\left(E, C_{2}{ }^{z}, \sigma_{x z}\right.$, and $\sigma_{y z}$ ) to these orbitals. A rotation by $180^{\circ}$ around the $z$-axis $\left(C_{2}{ }^{z}\right)$ interchanges the orbitals $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ (6-15 and 6-16). On the oxygen atom, the $2 s$ and $2 p_{z}$ are transformed into themselves (the axis of revolution for $2 p_{z}$ is identical to the rotation axis) (6-17 and 6-18), whereas $2 p_{x}$ and $2 p_{y}$ are transformed into their opposites (6-19 and 6-20).


6-15



6-17


6-18

If the same exercise is undertaken for the three other symmetry operations, we obtain the results that are presented in Table 6.1.

6-19

6-20

Table 6.1. Transformation of the atomic orbitals of the hydrogen and oxygen atoms of the $\mathrm{H}_{2} \mathrm{O}$ molecule by the action of the symmetry operations of the $C_{2 v}$ point group

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | :--- | :---: | :---: | :---: |
| $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ |
| $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{2}}$ |
| $2 s$ | $2 s$ | $2 s$ | $2 s$ | $2 s$ |
| $2 p_{x}$ | $2 p_{x}$ | $-2 p_{x}$ | $2 p_{x}$ | $-2 p_{x}$ |
| $2 p_{y}$ | $2 p_{y}$ | $-2 p_{y}$ | $-2 p_{y}$ | $2 p_{y}$ |
| $2 p_{z}$ | $2 p_{z}$ | $2 p_{z}$ | $2 p_{z}$ | $2 p_{z}$ |

No matter which symmetry operation is applied, it is clear that the $2 s$ orbital on the oxygen atom is transformed into itself. It therefore constitutes by itself a set that is globally stable, and so it is a basis for the representation of the point group. As the dimension of this basis is 1 , it is impossible to reduce it further: the $2 s$ orbital is a basis of a one-dimensional irreducible representation in the $C_{2 v}$ point group. The same applies for the other AO of the oxygen atom ( $2 p_{x}, 2 p_{y}$, and $2 p_{z}$ ): the action of the symmetry operations transforms each of them either into itself, or into its opposite (Table 6.1), so each of them constitutes a globally stable set.

We now examine the action of the symmetry operations on the $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ orbitals on the hydrogen atoms (the representation $\Gamma_{\mathrm{H}}$ ). It is immediately clear that neither $1 s_{\mathrm{H}_{1}}$ nor $1 s_{\mathrm{H}_{2}}$ constitutes by itselfa set that is globally stable, since some symmetry operations transform $1 s_{\mathrm{H}_{1}}$ into $1 s_{\mathrm{H}_{2}}$, and vice-versa (Table 6.1). However, the set $\left\{1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}\right\}$ is stable, and so it is a basis for a representation of the group. Moreover, if the symmetry operations are applied to the two linear combinations $\left(1 s_{\mathrm{H}_{1}}+\right.$ $\left.1 s_{\mathrm{H}_{2}}\right)$ and ( $1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}$ ), we see that they are globally stable towards the action of the symmetry operations, since each is transformed either into itself or into its opposite (Table 6.2). Both these linear combinations are
${ }^{1}$ These two functions are also called the symmetry-adapted linear combinations of the atomic orbitals $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ (we shall return to this point in greater detail in $\S 6.4$ ).


6-21

Table 6.2. Transformation of the linear combinations $\left(1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}\right)$ and $\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}\right)$ of the atomic orbitals of the hydrogen atoms in the $\mathrm{H}_{2} \mathrm{O}$ molecule by the action of the symmetry operations of the $C_{2 v}$ point group

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | :--- | :--- | :--- | :--- |
| $1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{2}}+1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{2}}+1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}$ |
| $1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}$ | $-\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}\right)$ | $-\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}\right)$ | $1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}$ |

therefore bases for an irreducible representation whose dimension is 1 in the $C_{2 v}$ point group. ${ }^{1}$

### 6.2.4. Characters

### 6.2.4.1. Representation of the group

As we have already remarked, the action of a symmetry operation $R_{k}$ on a basis for a representation $f=\left\{f_{1}, f_{2}, \ldots, f_{i}, \ldots, f_{n}\right\}$ of the group transforms each of the functions $f_{i}$ into a linear combination $f_{i}^{\prime}$ of the different functions of the set $f$. This action can be represented by a matrix $\mathbf{M}_{k}$ such that:

$$
\begin{equation*}
\mathbf{M}_{k} f_{i}=f_{i}^{\prime} \tag{6.2}
\end{equation*}
$$

This is an $(n \times n)$ matrix, where $n$ is the dimension of the basis. The collection of matrices associated with the different symmetry operations constitutes a representation of the point group.

### 6.2.4.2. Characters: the $\mathrm{C}_{3 v}$ point group

The following symmetry elements are present in the $\mathrm{NH}_{3}$ molecule: a $C_{3}(z)$ axis and three planes $\sigma_{v}(i)$ that contain the $\mathrm{N}-\mathrm{H}_{i}$ bonds and bisect the opposite $\mathrm{H}-\mathrm{N}-\mathrm{H}$ angle (6-21). These symmetry elements are characteristic of the $C_{3 v}$ point group.

Consider the orbitals $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ which constitute a basis written $\Gamma_{\mathrm{H}}$. The action of the symmetry operations ( $E, C_{3}^{1}, C_{3}^{2}, \sigma_{v}(1)$, $\left.\sigma_{v}(2), \sigma_{v}(3)\right)$ on these orbitals gives the results presented in Table 6.3.

Each $(3 \times 3)$ matrix associated with a symmetry operation (Figure 6.1) can easily be constructed from this table. Notice in particular that the matrix associated with the identity operation $(E)$ is diagonal; all its non-zero terms lie on the diagonal and moreover, they are all equal to 1 .

Each of these matrices can be characterized by its trace, that is the sum of its diagonal elements. In group theory, this trace is called the character of a matrix, written $\chi$. Notice that the character in Figure 6.1

Table 6.3. Transformation of the atomic orbitals $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ of the hydrogen atoms of the $\mathrm{NH}_{3}$ molecule by the action of the symmetry operations of the $C_{3 v}$ point group

| $C_{3 v}$ | $E$ | $C_{3}^{1}$ | $C_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ | $\sigma_{v}(3)$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ |
| $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ |
| $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ |

$$
\begin{aligned}
& E
\end{aligned} \begin{array}{ll}
{\left[\begin{array}{lll}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{l}
1 s_{H_{1}} \\
1 s_{H_{2}} \\
1 s_{H_{3}}
\end{array}\right]=\left[\begin{array}{l}
1 s_{H_{1}} \\
1 s_{H_{2}} \\
1 s_{H_{3}}
\end{array}\right]} & \chi(E)=3 \\
C_{3}^{1} & {\left[\begin{array}{lll}
0 & 0 & 1 \\
1 & 0 & 0 \\
0 & 1 & 0
\end{array}\right]\left[\begin{array}{l}
1 s_{H_{1}} \\
1 s_{H_{2}} \\
1 s_{H_{3}}
\end{array}\right]=\left[\begin{array}{l}
1 s_{H_{3}} \\
1 s_{H_{1}} \\
1 s_{H_{2}}
\end{array}\right]}
\end{array} \begin{aligned}
& \chi\left(C_{3}^{1}\right)=0 \\
& C_{3}^{2} \\
& \sigma_{v}(1) \quad\left[\begin{array}{lll}
0 & 1 & 0 \\
0 & 0 & 1 \\
1 & 0 & 0
\end{array}\right]\left[\begin{array}{l}
1 s_{H_{1}} \\
1 s_{H_{2}} \\
1 s_{H_{3}}
\end{array}\right]=\left[\begin{array}{l}
1 s_{H_{2}} \\
1 s_{H_{3}} \\
1 s_{H_{1}}
\end{array}\right]
\end{aligned} \begin{aligned}
& \chi\left(C_{3}^{2}\right)=0 \\
& 0 \\
& 0
\end{aligned} 1
$$

Figure 6.1. Matrix representation of the action of the symmetry operations of the $C_{3 v}$ point group on the $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ orbitals of the hydrogen atoms of the $\mathrm{NH}_{3}$ molecule.
associated with the identity operation (three in this example) is equal to the dimension of the basis; this result is general. Notice also that the characters associated with the symmetry operations that belong to a given class (see $\S$ 6.2.1) are equal (to 0 for the operations $C_{3}^{1}$ and $C_{3}^{2}$, and to 1 for the operations $\sigma_{v}(1), \sigma_{v}(2)$, and $\left.\sigma_{v}(3)\right)$.

From now on, the action of the symmetry operations on a basis will be represented by the set of characters associated with each symmetry operation; the operations may be grouped into classes where appropriate. The symmetry properties of the basis $\left(\Gamma_{\mathrm{H}}\right)$ formed by the $1 s_{\mathrm{H}_{1}}$, $1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ orbitals of the hydrogen atoms of the $\mathrm{NH}_{3}$ molecule can be summarized as follows (Table 6.4):

Table 6.4. Characters associated with the basis $\Gamma_{\mathrm{H}}$ formed by the orbitals $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ of the hydrogen atoms in the $\mathrm{NH}_{3}$ molecule ( $C_{3 v}$ point group)

| $C_{3 v}$ | $E$ | $C_{3}^{1}$ | $C_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ | $\sigma_{v}(3)$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\Gamma_{\mathrm{H}}$ | 3 | 0 | 0 | 1 | 1 | 1 |

Or alternatively

| $C_{3 v}$ | $E$ | $2 C_{3}$ | $3 \sigma_{v}$ |
| :--- | :--- | :--- | :--- |
| $\Gamma_{\mathrm{H}}$ | 3 | 0 | 1 |

### 6.2.5. Character tables

A character table provides a complete list of the irreducible representations associated with a given point group, as well as other useful information. The appearance of these tables, which is identical for all point groups, is illustrated below for the $C_{2 v}$ (e.g. $\mathrm{H}_{2} \mathrm{O}$ ) and $C_{3 v}$ (e.g. $\mathrm{NH}_{3}$ ) point groups.

### 6.2.5.1. Character table for the $C_{2 v}$ point group

This table can be described in the following ways (Table 6.5):

1. The symbol of the point group appears in the top left-hand corner.
2. The symmetry operations, grouped in classes, appear on the first line. However, in this point group, there is only one operation per class.
3. The symbols given to the different irreducible representations of the point group appear underneath the point-group symbol, in the first column. These are known as Mulliken symbols.
4. The characters associated with each irreducible representation are found in the central part of the table. For example, the characters for the $A_{1}$ representation are $(1,1,1,1)$, for the $B_{2}$ representation

Table 6.5. Character table for the $C_{2 v}$ point group

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |  |  |
| :--- | ---: | ---: | ---: | ---: | :---: | :---: |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 | 1 | $z$ | $x^{2}, y^{2}, z^{2}$ |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 | -1 |  | $x y$ |
| $\mathrm{~B}_{1}$ | 1 | -1 | 1 | -1 | $x$ | $x z$ |
| $\mathrm{~B}_{2}$ | 1 | -1 | -1 | 1 | $y$ | $y z$ |

$(1,-1,-1,1)$, etc. Every irreducible representation of this group is one-dimensional, since $\chi(E)=1$ for each of them.
5. The last two columns on the right contain algebraic functions which are bases for the irreducible representations on the same line. For example, $z$ is a basis for the $A_{1}$ irreducible representation, and $x z$ a basis for the $B_{1}$ irreducible representation.

These last two points need further comment. The $\mathrm{A}_{1}$ representation, for which every character is 1 , is also known as the totally symmetric representation, since any of the possible symmetry operations of the group transforms a basis function into itself. There is a totally symmetric representation in each point group. The characters for the other representations are either 1 or -1 , depending on whether the function is transformed into itself or into its opposite. The algebraic functions given in the last two columns on the right allow us to find immediately the symmetry of the AO on the 'central' atom, that is, the one whose position is unchanged by any of the symmetry operations, or the one located at the intersection of all the symmetry elements (the oxygen atom in the case of the water molecule). So the $p_{x}, p_{y}$, and $p_{z}$ orbitals transform in the same way as $x, y$, and $z$, respectively, and the $s$ orbital (with spherical symmetry) like $\left(x^{2}+y^{2}+z^{2}\right)$. The symmetries of these orbitals are therefore $\mathrm{A}_{1}\left(s, p_{z}\right), \mathrm{B}_{1}\left(p_{x}\right)$, and $\mathrm{B}_{2}\left(p_{y}\right)$. The character table for the $C_{2 v}$ point group therefore allows us to reduce the basis $\Gamma_{\mathrm{O}}$ formed by the valence orbitals of oxygen:

$$
\begin{equation*}
\Gamma_{\mathrm{O}}=2 \mathrm{~A}_{1} \oplus \mathrm{~B}_{1} \oplus \mathrm{~B}_{2} \tag{6.3}
\end{equation*}
$$

Notice that the last column of the character table contains functions that are squares or second-order products of $x, y$, and $z$. The symmetries of these functions indicate the symmetries of the $d$ orbitals on the central atom, which is clearly very useful for the study of transition metal complexes. Thus, in a complex with $C_{2 v}$ symmetry, the $d_{x^{2}-y^{2}}$ and $d_{z^{2}}$ orbitals have $A_{1}$ symmetry, whereas the $d_{x y}, d_{x z}$, and $d_{y z}$ orbitals have $A_{2}$, $B_{1}$, and $B_{2}$ symmetries, respectively (Table 6.5). The basis of dimension 5 formed by the $d$ orbitals of the central atom $\left(\Gamma_{d}\right)$ can therefore be
reduced as follows:

$$
\begin{equation*}
\Gamma_{d}=2 \mathrm{~A}_{1} \oplus \mathrm{~A}_{2} \oplus \mathrm{~B}_{1} \oplus \mathrm{~B}_{2} \tag{6.4}
\end{equation*}
$$

### 6.2.5.2. Character table for the $C_{3 v}$ point group

We shall now examine a second character table, that associated with the $C_{3 v}$ point group (Table 6.6), the point group for the $\mathrm{NH}_{3}$ molecule. It is presented in the same format as the previous one, but it illustrates several characteristics of the point groups in which a rotation axis is present whose order is higher than 2 (here, a $C_{3}$-axis).

Notice first that on the first line, the symmetry operations are grouped in classes. Thus, the two rotation operations around the $C_{3}$-axis ( $C_{3}^{1}$ and $C_{3}^{2}$, 6-8 and 6-9) are written $2 C_{3}$, and the reflections in the three planes of symmetry, $\sigma_{v}(1), \sigma_{v}(2)$, and $\sigma_{v}(3)(6-21)$ are written $3 \sigma_{v}$. In general, the notation that is used indicates the number of symmetry operations in the class, followed by the symmetry element that is concerned. One of the irreducible representations in this point group has a character of two associated with the identity operation. This is therefore a two-dimensional irreducible representation, written E (be careful not to confuse this with the identity operation, $E$ ).

If we consult the last two columns of the character table, we can establish the symmetry properties of the orbitals on the central atom. The $s$ orbital is a basis for the $\mathrm{A}_{1}$ representation (or, more simply, it has $\mathrm{A}_{1}$ symmetry), like the $p_{z}$ orbital. The $p_{x}$ and $p_{y}$ orbitals form a basis for the two-dimensional representation $(\mathrm{E})$. This indicates that from the symmetry point of view, these orbitals cannot be separated. Neither of them taken separately, nor any linear combination of them, forms a set that is stable to the action of the symmetry operations $C_{3}^{1}$ and $C_{3}^{2}$. If we consider the $d$ orbitals on the central atom, the final column in the character table shows us that $d_{z^{2}}$ has $\mathrm{A}_{1}$ symmetry, whereas the orbital pairs $\left(d_{x^{2}-y^{2}}, d_{x y}\right)$ and $\left(d_{x y}, d_{y z}\right)$ have E symmetry.

To finish this section, we provide a few comments about the Mulliken symbols that are used for irreducible representations, though without full details. One-dimensional representations are indicated by the letters A or B, whereas the letters E and T are used for two- and

Table 6.6. Character table for the $C_{3 v}$ point group

| $C_{3 v}$ | $E$ | $2 C_{3}$ | $3 \sigma_{v}$ |  |  |
| :--- | ---: | ---: | ---: | :---: | :---: |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 | $z$ | $x^{2}+y^{2}, z^{2}$ |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 |  |  |
| E | 2 | -1 | 0 | $(x, y)$ | $\left(x^{2}-y^{2}, x y\right),(x z, y z)$ |

three-dimensional representations, respectively. In the groups which contain an inversion centre, the subscript $g$ (from the German gerade, which means even) is added to the representations that are symmetric with respect to inversion (a positive character in the column headed by $i$ ), and the subscript $u$ (from the German ungerade, for odd) to the representations that are antisymmetric with respect to inversion (a negative character). Thus in the $C_{i}$ point group, in which the inversion centre is the only symmetry element, the two (one-dimensional) irreducible representations are written $\mathrm{A}_{g}$ and $\mathrm{A}_{u}$, and the functions that are bases of these representations transform into themselves or their opposites, respectively.

### 6.3. The reduction formula

As we have already seen ( $\S 6.2 .5$ ), the character table gives us information on orbital symmetry properties. If the molecule contains a central atom, the symmetries of the orbitals of this atom are indicated in the last two columns of the table. However, the orbitals on non-central atoms, for example the $1 s_{\mathrm{H}}$ orbitals in $\mathrm{H}_{2} \mathrm{O}$ or $\mathrm{NH}_{3}$, are not individually bases for an irreducible representation (Tables 6.1 and 6.3). These AO form a basis for a reducible representation that can be decomposed into a sum of irreducible representations of the point group. Although the character table does not give the result immediately, it does enable us to find it by using the reduction formula.

### 6.3.1. The reduction formula

If the characters $\chi_{\Gamma}$ associated with a reducible representation $\Gamma$ are known, it can be decomposed into a sum of irreducible representations ( $\Gamma=\sum_{i} a_{i} \Gamma_{i}$ ) of the point group by using the reduction formula:

$$
\begin{equation*}
a_{i}=\frac{1}{h} \sum_{k} \chi_{\Gamma}\left(R_{k}\right) \times \chi_{i}\left(R_{k}\right) \times n\left(R_{k}\right) \tag{6.5}
\end{equation*}
$$

where the summation $(k)$ is carried out over the classes, $h$ is the number of symmetry operations in the point group, also known as its order, $\chi_{\Gamma}\left(R_{k}\right)$ is the character of the reducible representation for a given class, $\chi_{i}\left(R_{k}\right)$ is the character of the irreducible representation for that class, and $n\left(R_{k}\right)$ is the number of symmetry operations in that class.

Before being able to apply the reduction formula, it is therefore necessary to determine the characters of the reducible representation being studied.

### 6.3.2. Characters of a reducible representation

We need intially to establish the characters of a reducible representation $\Gamma$. For a given symmetry operation $R$, only the diagonal terms of the matrix associated with this operation contribute to the character $\chi_{R}$ (§ 6.2.4.2). If the symmetry operation transforms the orbital under consideration into itself, the contribution to the character is +1 . If, however, it transforms the orbital into its opposite, the contribution is -1 . If the orbital is transformed into another one, the contribution is zero. We have already given an example of the calculation of these characters for the $1 s_{\mathrm{H}}$ orbitals on the hydrogen atoms in $\mathrm{NH}_{3}$ (Table 6.4). We shall now consider tow other examples, for the molecules $\mathrm{H}_{2} \mathrm{O}$ and $\mathrm{C}_{2} \mathrm{H}_{4}$.

Consider the basis $\Gamma_{\mathrm{H}}$ constituted by the two orbitals $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ on the hydrogen atoms in the water molecule. From Table 6.1, we notice that the two orbitals are transformed into themselves by the operations $E$ and $\sigma_{y z}$ (characters equal to $1+1=2$ ), whereas they are interchanged by the operations $C_{2}^{z}$ and $\sigma_{x z}$ (characters equal to $0+0=0$ ). The characters associated with the basis $\Gamma_{\mathrm{H}}$, listed in Table 6.7, do not correspond to any of the irreducible representations of the $C_{2 v}$ point group, which are all one-dimensional (Table 6.5). This is therefore a basis for a reducible representation.

The examples of $\mathrm{H}_{2} \mathrm{O}$ and $\mathrm{NH}_{3}$ (Table 6.4) illustrate a characteristic property of the $\Gamma_{\mathrm{H}}$ representations constituted by the set of $1 s_{\mathrm{H}_{1}}$ orbitals on the hydrogen atoms of a molecule. An orbital of this type can only be transformed either into itself (contribution to the character equal to +1 ) or into an orbital located on another hydrogen atom (contribution to the character of 0 ). To determine the character associated with a particular symmetry operation, it is therefore sufficient to count the number of hydrogen atoms that are left unchanged by this operation. This comment is also applicable to other orbitals that possess this property (e.g. the $s$ orbitals on heavy atoms that are equivalent by symmetry).

In other cases, it is necessary to examine carefully the ways in which the orbitals are transformed by the symmetry operations. Consider, for example, the ethylene molecule (6-22). It belongs to the $D_{2 h}$ point group and contains the following symmetry elements: the $C_{2}^{x}, C_{2}^{y}$, and $C_{2}^{z}$ axes,

Table 6.7. Characters associated with the representation $\Gamma_{\mathrm{H}}$ constituted by the $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ orbitals of the hydrogen atoms in the $\mathrm{H}_{2} \mathrm{O}$ molecule ( $C_{2 v}$ point group)

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | :--- | :--- | :--- | :--- |
| $\Gamma_{\mathrm{H}}$ | 2 | 0 | 0 | 2 |

an inversion center $i$ located in the middle of the carbon-carbon bond and the three planes $\sigma_{x y}, \sigma_{x z}$, and $\sigma_{y z}(\S 6.2 .2)$.

The results of applying the symmetry operations to the $p_{z 1}$ and $p_{z 2}$ orbitals are given in Table 6.8. It is clear that an orbital can be transformed into its opposite even if the symmetry operations does not move the atom on which it is found (e.g. the action of $C_{2}^{x}$ transform $p_{z 1}$ into $-p_{z 1}$, even though the position of the $\mathrm{C}_{1}$ atom is unaltered). The characters associated with the basis $\Gamma_{p z}$ cannot therefore be obtained simply by counting the number of carbon atoms whose position is left unchanged by each symmetry operation, in contrast to those of the $\Gamma_{\mathrm{H}}$ representation (Table 6.8).

Table 6.8. Transformation of the atomic orbitals $p_{z 1}$ and $p_{z 2}$ on the carbon atoms of the ethylene molecule (6-22) by the action of the symmetry operations of the $D_{2 h}$ point group, and the characters of the $\Gamma_{p z}$ and $\Gamma_{\mathrm{H}}$ representations

| $D_{2 h}$ | $E$ | $C_{2}^{z}$ | $C_{2}^{y}$ | $C_{2}^{x}$ | $i$ | $\sigma_{x y}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $p_{z 1}$ | $p_{z 1}$ | $p_{z 2}$ | $-p_{z 2}$ | $-p_{z 1}$ | $-p_{z 2}$ | $-p_{z 1}$ | $p_{z 1}$ | $p_{z 2}$ |
| $p_{z 2}$ | $p_{z 2}$ | $p_{z 1}$ | $-p_{z 1}$ | $-p_{z 2}$ | $-p_{z 1}$ | $-p_{z 2}$ | $p_{z 2}$ | $p_{z 1}$ |
| $\Gamma_{p z}$ | 2 | 0 | 0 | -2 | 0 | -2 | 2 | 0 |
| $\Gamma_{\mathrm{H}}$ | 4 | 0 | 0 | 0 | 0 | 4 | 0 | 0 |

### 6.3.3. Applications

### 6.3.3.1. $\mathrm{H}_{2} \mathrm{O}$ as an example

We shall now apply the reduction formula (6.5) to some of the reducible representations that we have already studied. Table 6.9 contains the character table for the $C_{2 v}$ point group and the characters that we have obtained for the $\Gamma_{\mathrm{H}}$ basis that is formed by the two $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ orbitals in the $\mathrm{H}_{2} \mathrm{O}$ molecule (Table 6.7).

Table 6.9. Characters of the irreducible representations of the $C_{2 v}$ point group and of the representation $\Gamma_{\mathrm{H}}$ in the $\mathrm{H}_{2} \mathrm{O}$ molecule

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | :--- | ---: | ---: | ---: |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 | 1 |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 | -1 |
| $\mathrm{~B}_{1}$ | 1 | -1 | 1 | -1 |
| $\mathrm{~B}_{2}$ | 1 | -1 | -1 | 1 |
| $\Gamma_{\mathrm{H}}$ | 2 | 0 | 0 | 2 |

As the order of the group ( $h$, the number of symmetry operations) is 4 , we obtain:

$$
\begin{aligned}
& a_{\mathrm{A}_{1}}=\frac{1}{4}[(2 \times 1 \times 1)+(0 \times 1 \times 1)+(0 \times 1 \times 1)+(2 \times 1 \times 1)]=1 \\
& a_{\mathrm{A}_{2}}=\frac{1}{4}[(2 \times 1 \times 1)+(0 \times 1 \times 1)-(0 \times 1 \times 1)-(2 \times 1 \times 1)]=0 \\
& a_{\mathrm{B}_{1}}=\frac{1}{4}[(2 \times 1 \times 1)+(0 \times 1 \times 1)+(0 \times 1 \times 1)-(2 \times 1 \times 1)]=0 \\
& a_{\mathrm{B}_{2}}=\frac{1}{4}[(2 \times 1 \times 1)-(0 \times 1 \times 1)+(0 \times 1 \times 1)+(2 \times 1 \times 1)]=1
\end{aligned}
$$

which leads to the expression

$$
\begin{equation*}
\Gamma_{\mathrm{H}}=\mathrm{A}_{1} \oplus \mathrm{~B}_{2} \tag{6.6}
\end{equation*}
$$

### 6.3.3.2. $\mathrm{NH}_{3}$ as an example

In the same way, we can decompose the $\Gamma_{\mathrm{H}}$ basis in the $\mathrm{NH}_{3}$ molecule (Table 6.10): its characters have already been obtained (Table 6.4):
Table 6.10. Characters of the irreducible
representations of the $C_{3 v}$ point group
and of the $\Gamma_{\mathrm{H}}$ representation in the

| $\mathrm{NH}_{3}$ | molecule |  |  |
| :--- | ---: | ---: | ---: |
| $C_{3 v}$ | $E$ | $2 C_{3}$ | $3 \sigma_{v}$ |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 |
| $\mathrm{~A}_{2}$ | 1 | 1 | -1 |
| E | 2 | -1 | 0 |
| $\Gamma_{\mathrm{H}}$ | 3 | 0 | 1 |

We obtain

$$
\begin{aligned}
& a_{\mathrm{A}_{1}}=\frac{1}{6}[(3 \times 1 \times 1)+(0 \times 1 \times 2)+(1 \times 1 \times 3)]=1 \\
& a_{\mathrm{A}_{2}}=\frac{1}{6}[(3 \times 1 \times 1)+(0 \times 1 \times 2)+(1 \times(-1) \times 3)]=0 \\
& a_{\mathrm{E}}=\frac{1}{6}[(3 \times 2 \times 1)+(0 \times(-1) \times 2)+(1 \times 0 \times 3)]=1
\end{aligned}
$$

so

$$
\begin{equation*}
\Gamma_{\mathrm{H}}=\mathrm{A}_{1} \oplus \mathrm{E} \tag{6.7}
\end{equation*}
$$

The three-dimensional representation $\Gamma_{\mathrm{H}}$ is therefore decomposed into a one-dimensional $\left(\mathrm{A}_{1}\right)$ and a two-dimensional (E) irreducible representation.

### 6.3.4. Direct products

It is often useful to be able to determine the symmetry of a function that is a product of two (or more) functions whose symmetry is already known. This need arises, for example, when we consider a polyelectronic wave function that is written as a product of monoelectronic functions (atomic or molecular orbitals), or when we are interested in the overlap between two orbitals (see § 6.5.1).

Consider two functions $\phi_{1}$ and $\phi_{2}$, whose symmetries are $\mathrm{B}_{1}$ and $\mathrm{B}_{2}$, respectively, in the $C_{2 v}$ point group. The characters associated with the product function ( $\phi_{1} \times \phi_{2}$ ) are obtained, for each symmetry operation $R$, from the product of the characters $\chi_{R}\left(\mathrm{~B}_{1}\right)$ and $\chi_{R}\left(\mathrm{~B}_{2}\right)$ of the $\mathrm{B}_{1}$ and $B_{2}$ representations. Inspection of the character table for the $C_{2 v}$ point group (Table 6.5) shows that the characters that are obtained for the product function are identical to those of the $\mathrm{A}_{2}$ irreducible representation (Table 6.11).

Table 6.11. Direct product of the $B_{1}$ and $B_{2}$ irreducible representations in the $C_{2 v}$ point group

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | ---: | ---: | ---: | ---: |
| $\chi_{R}\left(\mathrm{~B}_{1}\right)$ | 1 | -1 | 1 | -1 |
| $\chi_{R}\left(\mathrm{~B}_{2}\right)$ | 1 | -1 | -1 | 1 |
| $\chi_{R}\left(\mathrm{~B}_{1}\right) \times \chi_{R}\left(\mathrm{~B}_{2}\right)$ | 1 | 1 | -1 | -1 |
| $\chi_{R}\left(\mathrm{~A}_{2}\right)$ | 1 | 1 | -1 | -1 |

This result can be written as follows:

$$
\begin{equation*}
\mathrm{B}_{1} \times \mathrm{B}_{2}=\mathrm{A}_{2} \tag{6.8}
\end{equation*}
$$

Note that if the two functions $\phi_{1}$ and $\phi_{2}$ have the same symmetry in a group whose irreducible representations are one-dimensional, the product of the characters is necessarily equal to 1 for all the symmetry operations. The product function is therefore a basis for the totally symmetric representation ( $\mathrm{A}_{1}$ in the case of the $C_{2 v}$ point group). But the situation is more complicated when the point group contains a two- or three-dimensional (degenerate) representation. Consider, for example, the product $\mathrm{E} \times \mathrm{E}$ in the $C_{3 v}$ point group (Table 6.12). The characters that are obtained do not correspond to any of the irreducible representations of this group (see Table 6.6). The representation $\mathrm{E} \times \mathrm{E}$ in fact forms a basis for a reducible four-dimensional representation; when this is reduced, the totally symmetric representation $\left(A_{1}\right)$ of the $C_{3 v}$ group appears as one component.

$$
\begin{equation*}
\mathrm{E} \times \mathrm{E}=\mathrm{A}_{1} \oplus \mathrm{~A}_{2} \oplus \mathrm{E} \tag{6.9}
\end{equation*}
$$

Table 6.12. The $\mathrm{E} \times \mathrm{E}$ direct product in the $C_{3 v}$ point group

| $C_{3 v}$ | $E$ | $2 C_{3}$ | $3 \sigma_{v}$ |
| :--- | ---: | ---: | ---: |
| $\chi_{R}(\mathrm{E})$ | 2 | -1 | 0 |
| $\chi_{R}(\mathrm{E}) \times \chi_{R}(\mathrm{E})$ | 4 | 1 | 0 |

In general, the product of two functions of the same symmetry is either a basis for the totally symmetric representation of the group, or is a basis for a reducible representation that contains it. In contrast, the product of two functions whose symmetries are different never contains the totally symmetric representation.

### 6.4. Symmetry-adapted orbitals

Once the decomposition of a reducible representation into a sum of irreducible representations has been achieved, the following step consists of finding the linear combination of orbitals that are bases for these irreducible representations. These are often referred to as symmetryadapted linear combinations of orbitals (SALCO).

### 6.4.1. Projection operator

To find these combinations, we use a 'projection operator' $P$, whose action on a function $\phi$ is defined as follows (without considering normalization):

$$
\begin{equation*}
P \phi=\left[\sum_{k} \chi_{i}\left(R_{k}\right) R_{k}\right] \phi \tag{6.10}
\end{equation*}
$$

where $\chi_{i}\left(R_{k}\right)$ is the character associated with the operation $R_{k}$ for the irreducible representation being considered, $\Gamma_{i},{ }^{2} R_{k}$ is the symmetry operation whose character is $\chi_{i}\left(R_{k}\right)$ and $\phi$ is one of the orbitals of the reducible representation or a linear combination of these orbitals. It is called a generating function.

Note that the summation $(k)$ in this formula is carried out over all the symmetry operations. If there are several operations in a given class of the group, we must take account of each of them.

### 6.4.2. Application

### 6.4.2.1. The $\Gamma_{\mathrm{H}}$ basis in $\mathrm{H}_{2} \mathrm{O}$

A simple example will show how this formula is used. Consider the $\Gamma_{\mathrm{H}}$ basis that is constituted by the two $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ orbitals in the $\mathrm{H}_{2} \mathrm{O}$
molecule. We have already show that $\Gamma_{\mathrm{H}}$ can be reduced to the sum of two irreducible representations, $\mathrm{A}_{1}$ and $\mathrm{B}_{2}$ (formula (6.6)). We shall now continue by finding the combination of these orbitals that forms a basis for the $\mathrm{A}_{1}$ representation. We choose one of these orbitals, $1 s_{\mathrm{H}_{1}}$, for example, which will act as the generating function $\phi$ in equation (6.10). We then construct a table which contains (i) the result of the action of symmetry operation $R_{k}$ on the generating function (the term $R_{k} \phi$ in equation (6.10)); (ii) the characters of the irreducible representation under consideration (the term $\chi_{i}\left(R_{k}\right)$ ) and last, the product of these two terms (Table 6.13). The sum of these products over all the symmetry operations (equation (6.10)) gives function $\phi_{\mathrm{A}_{1}}$ that we seek, that is, the linear combination of the $1 s_{\mathrm{H}_{1}}$, and $1 s_{\mathrm{H}_{2}}$, orbitals that is a basis for the $\mathrm{A}_{1}$ representation.

Table 6.13. Quantities necessary for the determination of the linear combination of the $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ orbitals that has $\mathrm{A}_{1}$ symmetry in the $\mathrm{H}_{2} \mathrm{O}$ molecule

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | :--- | :--- | :--- | :--- |
| $R_{k}\left(1 s_{\mathrm{H}_{1}}\right)$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 | 1 |
| Product | $1 \times 1 s_{\mathrm{H}_{1}}$ | $1 \times 1 s_{\mathrm{H}_{2}}$ | $1 \times 1 s_{\mathrm{H}_{2}}$ | $1 \times 1 s_{\mathrm{H}_{1}}$ |

We therefore obtain: $\phi_{\mathrm{A}_{1}}=\left(1 \times 1 s_{\mathrm{H}_{1}}\right)+\left(1 \times 1 s_{\mathrm{H}_{2}}\right)+\left(1 \times 1 s_{\mathrm{H}_{2}}\right)+$ $\left(1 \times 1 s_{\mathrm{H}_{1}}\right)=2 \times\left(1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}\right)$. It is clear that we would have obtained the same result if we had taken the $1 s \mathrm{H}_{2}$ orbital as the generating function, as this would merely have led to an interchange of the subscripts 1 and 2 in the second line of Table 6.13.

The linear combination that has $B_{2}$ symmetry is obtained in the same way, from the data in Table 6.14. It is: $\phi_{\mathrm{B}_{2}}=\left(1 \times 1 s_{\mathrm{H}_{1}}\right)-(1 \times$ $\left.1 s_{\mathrm{H}_{2}}\right)-\left(1 \times 1 s_{\mathrm{H}_{2}}\right)+\left(1 \times 1 s_{\mathrm{H}_{1}}\right)=2 \times\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}\right)$.

In summary, the linear combinations of the $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ orbitals that are adapted to the molecular symmetry of $\mathrm{H}_{2} \mathrm{O}$ are the sum and difference of these two. If we include the normalization factor that is

Table 6.14. Quantities necessary for the determination of the linear combination of the $1 s_{\mathrm{H}_{1}}$ and $1 s_{\mathrm{H}_{2}}$ orbitals that has $\mathrm{B}_{2}$ symmetry in the $\mathrm{H}_{2} \mathrm{O}$ molecule

| $C_{2 v}$ | $E$ | $C_{2}^{z}$ | $\sigma_{x z}$ | $\sigma_{y z}$ |
| :--- | :--- | :--- | :--- | :--- |
| $R_{k}\left(1 s_{\mathrm{H}_{1}}\right)$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ |
| $\mathrm{~B}_{2}$ | 1 | -1 | -1 | 1 |
| Product | $1 \times 1 s_{\mathrm{H}_{1}}$ | $(-1) \times 1 s_{\mathrm{H}_{2}}$ | $(-1) \times 1 s_{\mathrm{H}_{2}}$ | $1 \times 1 s_{\mathrm{H}_{1}}$ |

${ }^{3}$ We have already shown (see Table 6.2) that these linear combinations are stable with respect to all the symmetry operations of the $C_{2 v}$ point group.



Figure 6.2. Symmetry-adapted orbitals for $\mathrm{H}_{2} \mathrm{O}$ ( $C_{2 v}$ point group).
given by Hückel theory, the expressions are:

$$
\begin{align*}
\phi_{\mathrm{A}_{1}} & =\frac{1}{\sqrt{2}}\left(1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}\right)  \tag{6.11a}\\
\phi_{\mathrm{B}_{2}} & =\frac{1}{\sqrt{2}}\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}\right) \tag{6.11b}
\end{align*}
$$

This result could have been anticipated, since we have only two equivalent functions to combine. ${ }^{3}$ These symmetry-adapted orbital are shown in Figure 6.2.

## Comment

In the standard (non-extended) Hückel method, the overlap between orbitals $(S)$ is neglected in the calculation of the normalization factor. As a result, the sum of the squares of the coefficients is 1 . In what follows, when we describe symmetry-adapted orbitals as normalized, we shall always adopt that approximation.

### 6.4.2.2. The $\Gamma_{\mathrm{H}}$ basis in $\mathrm{NH}_{3}$

We now consider a second example, one where the result is more difficult to predict: the three $1 s_{\mathrm{H}}$ orbitals in the $\mathrm{NH}_{3}$ molecule. These orbitals form a basis $\Gamma_{H}=A_{1} \oplus E$ (formula (6.7)). We shall proceed in the same way, taking the $1 s_{\mathrm{H}_{1}}$ orbital as the generating function. The second line of Table 6.3 shows how this orbital is transformed by the symmetry operation of the $C_{3 v}$ point group. We may use this result to construct Table 6.15, and hence find the linear combination that has $\mathrm{A}_{1}$ symmetry.

By adding the terms in the last line, we obtain: $\phi_{\mathrm{A}_{1}}=2 \times\left(1 s_{\mathrm{H}_{1}}+\right.$ $1 s_{\mathrm{H}_{2}}+1 s_{\mathrm{H}_{3}}$ ). The (non-normalized) totally symmetric function is therefore simply the sum of the initial functions, the same result as that obtained for $\mathrm{H}_{2} \mathrm{O}$.

For the representation with E symmetry, we must find two independent linear combinations of the $1 s_{\mathrm{H}}$ orbitals (a two-dimensional

Table 6.15. Quantities necessary for the determination of the linear combination of the $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ orbitals that has $\mathrm{A}_{1}$ symmetry in the $\mathrm{NH}_{3}$ molecule

| $C_{3 v}$ | $E$ | $C_{3}^{1}$ | $C_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ | $\sigma_{v}(3)$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $R_{k}\left(1 s_{\mathrm{H}_{1}}\right)$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 | 1 | 1 | 1 |
| Product | $1 \times 1 s_{\mathrm{H}_{1}}$ | $1 \times 1 s_{\mathrm{H}_{3}}$ | $1 \times 1 s_{\mathrm{H}_{2}}$ | $1 \times 1 s_{\mathrm{H}_{1}}$ | $1 \times 1 s_{\mathrm{H}_{3}}$ | $1 \times 1 s_{\mathrm{H}_{2}}$ |

Table 6.16. Quantities necessary for the determination of the first linear combination of the $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ orbitals with E symmetry in the $\mathrm{NH}_{3}$ molecule

| $C_{3 v}$ | $E$ | $C_{3}^{1}$ | $C_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ | $\sigma_{v}(3)$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $R_{k}\left(1 s_{\mathrm{H}_{1}}\right)$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{2}}$ |
| E | 2 | -1 | -1 | 0 | 0 | 0 |
| Product | $2 \times 1 s_{\mathrm{H}_{1}}$ | $(-1) \times 1 s_{\mathrm{H}_{3}}$ | $(-1) \times 1 s_{\mathrm{H}_{2}}$ | $0 \times 1 s_{\mathrm{H}_{1}}$ | $0 \times 1 s_{\mathrm{H}_{3}}$ | $0 \times 1 s_{\mathrm{H}_{2}}$ |

Table 6.17. Quantities necessary for the determination of the second linear combination of the $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ orbitals with E symmetry in the $\mathrm{NH}_{3}$ molecule

| $C_{3 v}$ | $E$ | $C_{3}^{1}$ | $C_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ | $\sigma_{v}(3)$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $R_{k}\left(1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}\right)$ | $1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}$ | $1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{3}}-1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{3}}-1 s_{\mathrm{H}_{2}}$ | $1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{1}}$ | $1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{3}}$ |
| E | 2 | -1 | -1 | 0 | 0 |  |
| Product | $2 \times\left(1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}\right)$ | $(-1) \times\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}\right)$ | $(-1) \times\left(1 s_{\mathrm{H}_{3}}-1 s_{\mathrm{H}_{1}}\right)$ | $0 \times\left(1 s_{\mathrm{H}_{3}}-1 s_{\mathrm{H}_{2}}\right)$ | $0 \times\left(1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{1}}\right)$ | $0 \times\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{3}}\right)$ |

${ }^{4}$ The second function can also be found by another method, from the requirement that it be orthogonal to the symmetry-adapted orbitals already found ( $\phi_{\mathrm{A} 1}$ and $\phi_{\mathrm{E}}(1)$ ) (see Exercise 6.7).
representation). If we use the $1 s_{\mathrm{H}_{1}}$ orbital once again as the generating function, and the characters of the E representation (Table 6.16), we obtain $\phi_{\mathrm{E}}(1)=2 \times\left(1 s_{\mathrm{H}_{1}}\right)-1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}$.

In this way, we have found one of the two linear combinations that we need. If we now change the generating function to $1 s_{\mathrm{H}_{2}}$ or $1 s_{\mathrm{H}_{3}}$, we shall clearly obtain the same type of function, but where the subscripts 1,2 , and 3 have been permuted. This function is acceptable, but it is not orthogonal to the previous one, and in general we prefer to use linear combinations of the orbitals that are mutually orthogonal. ${ }^{4}$ In this example, the second linear combination can be obtained if the combination $\left(1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}\right)$ is used as the generating function. Action of the symmetry operations on this function gives the results in Table 6.17.

Taking the characters of the E representation into account, we obtain the function: $\phi_{\mathrm{E}}(2)=2 \times\left(1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}\right)+(-1) \times\left(1 s_{\mathrm{H}_{1}}-\right.$ $\left.1 s_{\mathrm{H}_{2}}\right)+(-1) \times\left(1 s_{\mathrm{H}_{3}}-1 s_{\mathrm{H}_{1}}\right)=3 \times\left(1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}\right)$, that is, the generating function itself (without considering normalization).

In conclusion, the normalized linear combinations of the $1 s_{\mathrm{H}_{1}}, 1 s_{\mathrm{H}_{2}}$, and $1 s_{\mathrm{H}_{3}}$ orbitals that are adapted to the molecular symmetry of $\mathrm{NH}_{3}$ are:

$$
\begin{align*}
\phi_{\mathrm{A}_{1}} & =\frac{1}{\sqrt{3}}\left(1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}+1 s_{\mathrm{H}_{3}}\right)  \tag{6.12a}\\
\phi_{\mathrm{E}}(1) & =\frac{1}{\sqrt{6}}\left(2 \times\left(1 s_{\mathrm{H}_{1}}\right)-1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}\right)  \tag{6.12b}\\
\phi_{\mathrm{E}}(2) & =\frac{1}{\sqrt{2}}\left(1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}\right) \tag{6.12c}
\end{align*}
$$




$2 s$

$2 p_{z}$


Figure 6.3. Symmetry-adapted orbitals for $\mathrm{NH}_{3}\left(C_{3 v}\right.$ point group) and orbitals of the same symmetry on the central atom.

These orbitals are shown in Figure 6.3, together with the orbitals on the central atom whose symmetry is given in the character table for the $C_{3 v}$ point group (Table 6.6).

The set of orbitals that we have found for the E representations, $\phi_{\mathrm{E}}(1)$ and $\phi_{\mathrm{E}}(2)$, is not unique. Any pair of independent linear combinations of these two orbitals also constitutes a basis for this representation, and the same is true for the $2 p_{x}$ and $2 p_{y}$ orbitals on the central atom. The $\phi_{\mathrm{E}}(1)$ and $\phi_{\mathrm{E}}(2)$ orbitals shown here are, however, the set that is used most frequently. The first is transformed in the same way as $2 p_{x}$ by all the symmetry operations of the $C_{3 v}$ point group and the second like $2 p_{y}$. They are therefore often written $e_{x}$ and $e_{y}$, respectively.

### 6.5. Construction of $\mathrm{MO}: \mathrm{H}_{2} \mathrm{O}$ as an example

The problem of allowing AO, or SALCO of these orbitals, to interact to form molecular orbitals (MO), is simplified considerably if the symmetry properties of the system are taken into account. The use of symmetry allows us to identify rapidly those interactions which are exactly zero, and therefore to consider only those which really do contribute to the formation of the MO.

### 6.5.1. Symmetry and overlap

Two orbitals $\phi_{1}$ and $\phi_{2}$ interact if their overlap is non-zero (Chapter $1, \S$ 1.3). This overlap is equal to the integral over all space of the product of the functions $\phi_{1}$ and $\phi_{2}$ :

$$
\begin{equation*}
S_{12}=\int_{\text {space }} \phi_{1}^{*} \phi_{2} \mathrm{~d} \tau \tag{6.13}
\end{equation*}
$$

It can be shown that this integral is non-zero if the product function is a basis for the totally symmetric representation of the group (or of a
reducible representation that contains the totally symmetric representation). It is therefore necessary for the two orbitals to have the same symmetry (see $\oint$ 6.3.4) for the integral to be non-zero. If their symmetries are different, the integral is rigorously equal to zero, and their overlap is said to be zero by symmetry.

## Comment

These rules (that are not proved here) can be illustrated by two simple examples. Consider the function with a single variable $y=x^{2}$. The integral of this symmetric (or even) function between $-a$ and $+a$ is non-zero $\left(\left[x^{3} / 3\right]_{-a}^{+a}=\left(a^{3} / 3\right)-\left(-a^{3} / 3\right)=2 a^{3} / 3\right)$. However, the integral of the antisymmetric (or odd) function $y=x^{3}$, which changes sign when $x$ is replaced by $-x$, is equal to zero $\left(\left[x^{4} / 4\right]_{-a}^{+a}=\left(a^{4} / 4\right)-\left(a^{4} / 4\right)=0\right)$.

### 6.5.2. Molecular orbitals for $\mathrm{H}_{2} \mathrm{O}$

To construct the MO of the $\mathrm{H}_{2} \mathrm{O}$ molecule, we shall allow the atomic orbitals of the central oxygen atom to interact with the symmetryadapted orbitals on the hydrogen atoms (the fragment method). We have already established the symmetry properties of the various orbitals ( $C_{2 v}$ point group), and they are repeated in Figure 6.4. It is easy to check visually that the orbitals with the same symmetry on two different fragments have a non-zero overlap, whereas those with different symmetries have zero overlap ('by symmetry').

The interaction diagram is shown in Figure 6.5, where the orbitals on the fragments and in the full molecule are labelled by symmetry $\left(a_{1}, b_{2}, \ldots\right)$.

There is an interaction between the two orbitals with $\mathrm{B}_{2}$ symmetry that leads to the formation of a bonding MO, written $1 b_{2}$, and an antibonding MO, written $2 \mathrm{~b}_{2}$. As there are three orbitals with $\mathrm{A}_{1}$
$\mathrm{A}_{1}$

$2 s$

$2 p_{z}$

$\phi_{\mathrm{A}_{1}}$
$\mathrm{B}_{2}$

$2 p_{y}$

$\phi_{\mathrm{B}_{2}}$
$B_{1}$


Figure 6.4. Symmetry-adapted orbitals for $\mathrm{H}_{2} \mathrm{O}$ ( $C_{2 v}$ point group) and orbitals with the same symmetry on the central atom.

Figure 6.5. Construction of the MOs for $\mathrm{H}_{2} \mathrm{O}$ from AO on oxygen and the symmetry-adapted orbitals on the hydrogen atoms.


6-23


6-24

symmetry, three molecular orbitals are formed: $1 \mathrm{a}_{1}$ (bonding), $2 \mathrm{a}_{1}$ (nonbonding), and $3 \mathrm{a}_{1}$ (antibonding). The atomic orbital with $\mathrm{B}_{1}$ symmetry $\left(2 p_{x}\right)$ cannot, by symmetry, interact with any other, and so it stays unchanged in shape and in energy ( $1 \mathrm{~b}_{1}$, nonbonding).

### 6.6. Symmetry-adapted orbitals in several $\mathbf{M L}_{n}$ complexes

The aim of this section is to construct the symmetry-adapted orbitals for the principal ligand fields, making use of the reduction (6.5) and projection (6.10) formulae. In what follows, with the exception of § 6.6.6, we shall only consider a single orbital on each ligand, the one that is used to create the $\sigma$ bond with the metal (Chapter 1, § 1.5.1). The orbital on ligand $L_{i}$ will be written $\sigma_{i}$. We shall suppose that all the ligands, and thus all the orbitals $\sigma_{i}$, are identical.

### 6.6.1. Square-planar $\mathrm{ML}_{4}$ complexes

Consider a complex in which the metallic atom is surrounded by four ligands that are placed at the corners of a square ( $\mathbf{6 - 2 3}$ ). The symmetry elements of this system are characteristic of the $D_{4 h}$ point group. The axes are shown in $\mathbf{6 - 2 4}$. The planes of symmetry are $x y\left(\sigma_{h}\right), x z$ $\left(\sigma_{d a}\right)$, and $y z\left(\sigma_{d b}\right)$, respectively, together with the planes that bisect $x z$ and $y z$ and each contain two $\mathrm{M}-\mathrm{L}$ bonds ( $\sigma_{\mathrm{v} a}$ and $\sigma_{\mathrm{v} b}$, respectively). The inversion centre is of course at the origin, coincident with the central atom.

### 6.6.1.1. Reduction of the $\Gamma_{\sigma}$ representation

The action of a symmetry operation on one of the orbitals $\sigma_{i}$ can transform it only into itself or into an orbital $\sigma_{j}$ on another ligand. Following the procedure that we have already established for bases $\Gamma_{\mathrm{H}}$ constituted by $1 s_{\mathrm{H}}$ orbitals located on hydrogen atoms (§ 6.3.2), the character of the representation $\Gamma_{\sigma}$ is obtained simply by counting the number of ligands whose position stays unchanged. It is easy to see that the identity operation $(E)$ and reflection in the molecular plane $\left(\sigma_{h}\right)$ leave the positions of the four ligands unchanged $(\chi=4)$. Rotations around the $C_{2}^{\prime}$ axes and reflections in the $\sigma_{v}$ planes maintain the positions of the two ligands situated on the symmetry element concerned $(\chi=2)$. The other symmetry operations move all the ligands $(\chi=0)$. The characters obtained for the representation are listed in Table 6.18, together with the characters of the irreducible representations of the $D_{4 h}$ point group.

The order of the group ( $h$, the number of symmetry operations) is 16 (first line of Table 6.18). The reduction formula (6.5) enables us to decompose the four-dimensional representation $\Gamma_{\sigma}$ into a sum of irreducible representations, $\sum_{i} a_{i} \Gamma_{i}$. The only non-zero values of $a_{i}$ are:

$$
\begin{aligned}
& a_{\mathrm{A}_{1 g}}=\frac{1}{16}[(4 \times 1 \times 1)+(2 \times 1 \times 2)+(4 \times 1 \times 1)+(2 \times 1 \times 2)]=1 \\
& a_{\mathrm{B}_{1 g}}=\frac{1}{16}[(4 \times 1 \times 1)+(2 \times 1 \times 2)+(4 \times 1 \times 1)+(2 \times 1 \times 2)]=1 \\
& a_{\mathrm{E}_{u}}=\frac{1}{16}[(4 \times 2 \times 1)+(2 \times 0 \times 2)+(4 \times 2 \times 1)+(2 \times 0 \times 2)]=1
\end{aligned}
$$

hence

$$
\begin{equation*}
\Gamma_{\sigma}=\mathrm{A}_{1 g} \oplus \mathrm{~B}_{1 g} \oplus \mathrm{E}_{u} \tag{6.14}
\end{equation*}
$$

Table 6.18. Character table for the $D_{4 h}$ point group and characters of the reducible representation $\Gamma_{\sigma}$ of a square-planar complex $\mathrm{ML}_{4}$

| $D_{4 h}$ | $E$ | $2 C_{4}$ | $C_{2}$ | $2 C_{2}^{\prime}$ | $2 C_{2}^{\prime \prime}$ | $i$ | $2 S_{4}$ | $\sigma_{h}$ | $2 \sigma_{v}$ | $2 \sigma_{d}$ |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $\mathrm{~A}_{1 g}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | $x^{2}+y^{2}, z^{2}$ |
| $\mathrm{~A}_{2 g}$ | 1 | 1 | 1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 |  |
| $\mathrm{~B}_{1 g}$ | 1 | -1 | 1 | 1 | -1 | 1 | -1 | 1 | 1 | -1 | $x^{2}-y^{2}$ |
| $\mathrm{~B}_{2 g}$ | 1 | -1 | 1 | -1 | 1 | 1 | -1 | 1 | -1 | 1 | $x y$ |
| $\mathrm{E}_{g}$ | 2 | 0 | -2 | 0 | 0 | 2 | 0 | -2 | 0 | 0 |  |
| $\mathrm{~A}_{1 u}$ | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 |  |
| $\mathrm{~A}_{2 u}$ | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 | 1 | 1 | $z$ |
| $\mathrm{~B}_{1 u}$ | 1 | -1 | 1 | 1 | -1 | -1 | 1 | -1 | -1 | 1 |  |
| $B_{2 u}$ | 1 | -1 | 1 | -1 | 1 | -1 | 1 | -1 | 1 | -1 |  |
| $\mathrm{E}_{u}$ | 2 | 0 | -2 | 0 | 0 | -2 | 0 | 2 | 0 | 0 | $(x, y)$ |
| $\Gamma_{\sigma}$ | 4 | 0 | 0 | 2 | 0 | 0 | 0 | 4 | 2 | 0 |  |

The representation $\Gamma_{\sigma}$ is thus decomposed into two one-dimensional representations ( $\mathrm{A}_{1 g}$ and $\mathrm{B}_{1 g}$ ) and one degenerate two-dimensional representation $\left(\mathrm{E}_{u}\right)$.

### 6.6.1.2. Symmetry-adapted orbitals

Application of the projection formula (6.10) requires that we know how one of the ligand orbitals (the generating function) is transformed by all the symmetry operations. We shall consider two generating functions in turn, $\sigma_{1}$ and $\sigma_{2}$, where the numbering of the orbitals is shown in $\mathbf{6 - 2 3}$. The results of the action of the symmetry operations on these two orbitals, $R_{k}\left(\sigma_{1}\right)$ and $R_{k}\left(\sigma_{2}\right)$, are given in Table 6.19 , together with the characters of the irreducible representations $\mathrm{A}_{1 g}, \mathrm{~B}_{1 g}$, and $\mathrm{E}_{u}$.

The symmetry-adapted functions are obtained by multiplying the function $R_{k}\left(\sigma_{1}\right)$ (or $R_{k}\left(\sigma_{2}\right)$ ) for each symmetry operation by the character of the irreducible representation considered, and adding the sum of these products for all the symmetry operations (formula (6.10)).

By using $\sigma_{1}$ as a generating function for the symmetries $A_{1 g}$ and $B_{1 g}$, then $\sigma_{1}$ and $\sigma_{2}$ successively for the symmetry $\mathrm{E}_{u}$ (a two-dimensional, degenerate representation), we obtain:

$$
\begin{aligned}
\phi_{\mathrm{A}_{1 g}} & =4 \times\left(\sigma_{1}+\sigma_{2}+\sigma_{3}+\sigma_{4}\right) \\
\phi_{\mathrm{B}_{1 g}} & =4 \times\left(\sigma_{1}-\sigma_{2}+\sigma_{3}-\sigma_{4}\right) \\
\phi_{\mathrm{E}_{u}}(1) & =4 \times\left(\sigma_{1}-\sigma_{3}\right) \\
\phi_{\mathrm{E}_{u}}(2) & =4 \times\left(\sigma_{2}-\sigma_{4}\right)
\end{aligned}
$$

Table 6.19. Action of the symmetry operations of the $D_{4 h}$ point group on the orbitals $\sigma_{1}$ and $\sigma_{2}$ (see 6-23 and 6-24 for the numbering of the orbitals and the symmetry elements) and the characters of the irreducible representations $\mathrm{A}_{1 g}, \mathrm{~B}_{1 g}$, and $\mathrm{E}_{u}$

| $D_{4 h}$ | $E$ | $C_{4}^{1}$ | $C_{4}^{3}$ | $C_{2}$ | $C_{2 a}^{\prime}$ | $C_{2 b}^{\prime}$ | $C_{2 a}^{\prime \prime}$ | $C_{2 b}^{\prime \prime}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $R_{k}\left(\sigma_{1}\right)$ | $\sigma_{1}$ | $\sigma_{4}$ | $\sigma_{2}$ | $\sigma_{3}$ | $\sigma_{1}$ | $\sigma_{3}$ | $\sigma_{4}$ | $\sigma_{2}$ |
| $R_{k}\left(\sigma_{2}\right)$ | $\sigma_{2}$ | $\sigma_{1}$ | $\sigma_{3}$ | $\sigma_{4}$ | $\sigma_{4}$ | $\sigma_{2}$ | $\sigma_{3}$ | $\sigma_{1}$ |
| $\mathrm{~A}_{1 g}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |
| $\mathrm{~B}_{1 g}$ | 1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 |
| $\mathrm{E}_{u}$ | 2 | 0 | 0 | -2 | 0 | 0 | 0 | 0 |
|  |  |  |  |  |  |  |  |  |
| $D_{4 h}$ | $i$ | $S_{4}^{1}$ | $S_{4}^{3}$ | $\sigma_{h}$ | $\sigma_{v a}$ | $\sigma_{v b}$ | $\sigma_{d a}$ | $\sigma_{d b}$ |
| $R_{k}\left(\sigma_{1}\right)$ | $\sigma_{3}$ | $\sigma_{4}$ | $\sigma_{2}$ | $\sigma_{1}$ | $\sigma_{1}$ | $\sigma_{3}$ | $\sigma_{4}$ | $\sigma_{2}$ |
| $R_{k}\left(\sigma_{2}\right)$ | $\sigma_{4}$ | $\sigma_{1}$ | $\sigma_{3}$ | $\sigma_{2}$ | $\sigma_{4}$ | $\sigma_{2}$ | $\sigma_{3}$ | $\sigma_{1}$ |
| $\mathrm{~A}_{1 g}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |
| $\mathrm{~B}_{1 g}$ | 1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 |
| $\mathrm{E}_{u}$ | -2 | 0 | 0 | 2 | 0 | 0 | 0 | 0 |

$\phi_{\mathrm{A}_{1 g}}$





Figure 6.6. Symmetry-adapted $\sigma$ orbitals for a square-planar $\mathrm{ML}_{4}$ complex. Two different representations are shown for the degenerate $\mathrm{E}_{\mathcal{U}}$ orbitals.
which gives the following normalized expressions for the orbitals:

$$
\begin{align*}
\phi_{\mathrm{A}_{1 g}} & =\frac{1}{2}\left(\sigma_{1}+\sigma_{2}+\sigma_{3}+\sigma_{4}\right)  \tag{6.15a}\\
\phi_{\mathrm{B}_{1 g}} & =\frac{1}{2}\left(\sigma_{1}-\sigma_{2}+\sigma_{3}-\sigma_{4}\right)  \tag{6.15b}\\
\phi_{\mathrm{E}_{u}}(1) & =\frac{1}{\sqrt{2}}\left(\sigma_{1}-\sigma_{3}\right)  \tag{6.15c}\\
\phi_{\mathrm{E}_{u}}(2) & =\frac{1}{\sqrt{2}}\left(\sigma_{2}-\sigma_{4}\right) \tag{6.15~d}
\end{align*}
$$

These symmetry-adapted orbitals are shown in Figure 6.6. The solution found for the degenerate $\mathrm{E}_{u}$ orbitals is not unique, since any pair of independent linear combinations of these orbitals is also a basis for this representation. Another choice is often made, the (normalized) sum and difference of the two functions $\phi_{\mathrm{E}_{u}}(1)$ and $\phi_{\mathrm{E}_{u}}(2)$ (right-hand side of Figure 6.6).

### 6.6.2. Tetrahedral $\mathrm{ML}_{4}$ complexes

Consider a complex $\mathrm{ML}_{4}$ in which the four ligands are situated at the apices of a tetrahedron. Each ligand has a $\sigma$ orbital which points towards the metallic centre ( $\mathbf{6 - 2 5}$ ). The symmetry elements, which are characteristic of the $T_{d}$ point group, are:

- four $C_{3}$ axes, each of which is co-linear with one of the $\mathrm{M}-\mathrm{L}$ bonds;
- three $C_{2}$ axes that bisect the $\mathrm{L}-\mathrm{M}-\mathrm{L}$ angles;
- three $S_{4}$ axes that are co-linear with the $C_{2}$ axes;
- six $\sigma_{d}$ planes, each of which contains two $\mathrm{M}-\mathrm{L}$ bonds.

An example of each of these elements is shown in 6-26.

### 6.6.2.1. Reduction of the $\Gamma_{\sigma}$ representation

To obtain the characters that are associated with the reducible representation $\Gamma_{\sigma}$, we just need to count the number of ligands that are left unmoved by the various symmetry operations. For the identity operation, we clearly have $\chi=4$. Rotation around a $C_{3}$-axis leaves one ligand in its original position $(\chi=1)$, the one located on that axis. Rotation around the $C_{2}$ and $S_{4}$ axes moves all the ligands ( $\chi=0$ ). Reflection in a $\sigma_{d}$ plane leaves the two ligands that are in this plane in their original positions ( $\chi=2$ ). The characters that are obtained are given in Table 6.20, together with those for the irreducible representations of the $T_{d}$ point group.

The order of the group (the number of symmetry operations, $h$ ) is 24 (first line of Table 6.20). Application of the reduction formula (6.5)

Table 6.20. Character table for the $T_{d}$ point group and the characters of the reducible representation $\Gamma_{\sigma}$ in a tetrahedral $\mathrm{ML}_{4}$ complex

| $T_{d}$ | $E$ | $8 C_{3}$ | $3 C_{2}$ | $6 S_{4}$ | $6 \sigma_{d}$ |  |  |
| :--- | ---: | ---: | ---: | ---: | ---: | :--- | :--- |
| $\mathrm{~A}_{1}$ | 1 | 1 | 1 | 1 | 1 | $x^{2}+y^{2}+z^{2}$ |  |
| $\mathrm{~A}_{2}$ | 1 | 1 | 1 | -1 | -1 |  |  |
| E | 2 | -1 | 2 | 0 | 0 |  | $\left(2 z^{2}-x^{2}-y^{2}, x^{2}-y^{2}\right)$ |
| $\mathrm{T}_{1}$ | 3 | 0 | -1 | 1 | -1 |  |  |
| $\mathrm{~T}_{2}$ | 3 | 0 | -1 | -1 | 1 | $(x, y, z)$ | $(x y, x z, y z)$ |
| $\Gamma_{\sigma}$ | 4 | 1 | 0 | 0 | 2 |  |  |

gives:

$$
\begin{aligned}
& a_{\mathrm{A}_{1}}=\frac{1}{24}[(4 \times 1 \times 1)+(1 \times 1 \times 8)+(0 \times 1 \times 3)+(0 \times 1 \times 6) \\
&+(2 \times 1 \times 6)]=1 \\
& a_{\mathrm{A}_{2}}=\frac{1}{24}[(4 \times 1 \times 1)+(1 \times 1 \times 8)+(0 \times 1 \times 3)-(0 \times 1 \times 6) \\
&\quad(2 \times 1 \times 6)]=0 \\
& \\
& a_{\mathrm{E}}=\frac{1}{24}[(4 \times 2 \times 1)-(1 \times 1 \times 8)+(0 \times 2 \times 3)-(0 \times 0 \times 6) \\
&+(2 \times 0 \times 6)]=0 \\
& a_{\mathrm{T}_{1}}=\frac{1}{24}[(4 \times 3 \times 1)+(1 \times 0 \times 8)-(0 \times 1 \times 3)+(0 \times 1 \times 6) \\
&-(2 \times 1 \times 6)]=0 \\
& a_{\mathrm{T}_{2}}=\frac{1}{24}[(4 \times 3 \times 1)+(1 \times 0 \times 8)-(0 \times 1 \times 3)-(0 \times 1 \times 6) \\
&+(2 \times 1 \times 6)]=1
\end{aligned}
$$

hence

$$
\begin{equation*}
\Gamma_{\sigma}=\mathrm{A}_{1} \oplus \mathrm{~T}_{2} \tag{6.16}
\end{equation*}
$$

The representation $\Gamma_{\sigma}$ is thus decomposed into a one-dimensional representation $\left(\mathrm{A}_{1}\right.$, the totally symmetric representation) and a degenerate three-dimensional representation $\left(\mathrm{T}_{2}\right)$.

### 6.6.2.2. Symmetry-adapted orbitals

The linear combinations of the $\sigma_{i}$ orbitals that are bases for the $\mathrm{A}_{1}$ and $\mathrm{T}_{2}$ irreducible representations can be determined from the projection formula (6.10). Since there are many symmetry operations in this point group (24), the exercise is rather tedious, and we shall be content here


Figure 6.7. Symmetry-adapted $\sigma$ orbitals in a tetrahedral ML4 complex. Two different representations are given for the degenerate $\mathrm{T}_{2}$ orbitals.


6-27


6-28
just to give the result. The symmetry-adapted orbitals are:

$$
\begin{align*}
\phi_{\mathrm{A}_{1}} & =\frac{1}{2}\left(\sigma_{1}+\sigma_{2}+\sigma_{3}+\sigma_{4}\right)  \tag{6.17a}\\
\phi_{\mathrm{T}_{2}}(1) & =\frac{1}{\sqrt{2}}\left(\sigma_{1}-\sigma_{2}\right)  \tag{6.17b}\\
\phi_{\mathrm{T}_{2}}(2) & =\frac{1}{\sqrt{2}}\left(\sigma_{3}-\sigma_{4}\right)  \tag{6.17c}\\
\phi_{\mathrm{T}_{2}}(3) & =\frac{1}{2}\left(\sigma_{1}+\sigma_{2}-\sigma_{3}-\sigma_{4}\right) \tag{6.17d}
\end{align*}
$$

These orbitals are shown in Figure 6.7.
The expression of the $\mathrm{T}_{2}$ orbitals given in $(6.17 \mathrm{~b}-\mathrm{d})$ is not unique; as we have already seen for square-planar complexes in $\oint$ 6.6.1.2, any independent linear combination of these functions is acceptable. One such, which 'favours' a vertical $C_{3}$-axis, is shown on the right-hand side of Figure 6.7 (see Exercise 6.10).

### 6.6.3. Trigonal-planar $\mathrm{ML}_{3}$ complexes

A trigonal-planar $\mathrm{ML}_{3}$ complex (6-27) belongs to the $D_{3 h}$ point group and contains the following symmetry elements: the molecular plane $\left(\sigma_{h}\right)$, three perpendicular planes $\left(\sigma_{v}\right)$, each of which contains one $\mathrm{M}-\mathrm{L}$ bond, a $C_{3}$-axis and an $S_{3}$-axis which are co-linear and perpendicular to the $\sigma_{h}$ plane, and three $C_{2}$ axes, each of which is co-linear with one of the bonds. Each type of symmetry element is illustrated in 6-28.

### 6.6.3.1. Reduction of the representation $\Gamma_{\sigma}$

Rotations around the $C_{3}$ and $S_{3}$ axes change the positions of all the ligands $(\chi=0)$, but a rotation around a $C_{2}$-axis leaves the ligand on that axis unmoved $(\chi=1)$. Reflections in the $\sigma_{h}$ and $\sigma_{v}$ planes leave all three ligands in their original positions $(\chi=3)$. The characters obtained from these considerations are given in Table $6.21\left(\Gamma_{\sigma}\right)$, together

Table 6.21. Character table for the $D_{3 h}$ point group and the characters of the reducible representation $\Gamma_{\sigma}$ for a trigonal-planar $\mathrm{ML}_{3}$ complex

| $D_{3 h}$ | $E$ | $2 C_{3}$ | $3 C_{2}$ | $\sigma_{h}$ | $2 S_{3}$ | $3 \sigma_{v}$ |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | :--- |
| $\mathrm{~A}_{1}^{\prime}$ | 1 | 1 | 1 | 1 | 1 | 1 |  |
| $\mathrm{~A}_{2}^{\prime}$ | 1 | 1 | -1 | 1 | 1 | -1 |  |
| $\mathrm{E}^{\prime}$ | 2 | -1 | 0 | 2 | -1 | 0 | $(x, y)$ |
| $\mathrm{A}_{1}^{\prime \prime}$ | 1 | 1 | 1 | -1 | -1 | -1 |  |
| $\mathrm{~A}_{2}^{\prime \prime}$ | 1 | 1 | -1 | -1 | -1 | 1 | $z$ |
| $\mathrm{E}^{\prime \prime}$ | 2 | -1 | 0 | -2 | 1 | 0 |  |
| $\Gamma_{\sigma}$ | 3 | 0 | 1 | 3 | 0 | 1 | $(x z, y z)$ |

with those of the different irreducible representations of the $D_{3 h}$ point group.

As the order of the group $(h)$ is 12 , application of the reduction formula (6.5) leads to (products where at least one factor is zero have been omitted):

$$
\begin{aligned}
& \begin{aligned}
a_{\mathrm{A}_{1}^{\prime}}=\frac{1}{12} & {[(3 \times 1 \times 1)+(1 \times 1 \times 3)+(3 \times 1 \times 1)} \\
& +(1 \times 1 \times 3)]=1
\end{aligned} \\
& \begin{aligned}
a_{\mathrm{A}_{2}^{\prime}}= & \frac{1}{12}[(3 \times 1 \times 1)-(1 \times 1 \times 3)+(3 \times 1 \times 1) \\
& \quad(1 \times 1 \times 3)]=0
\end{aligned} \\
& a_{\mathrm{E}^{\prime}}=\frac{1}{12}[(3 \times 2 \times 1)+(3 \times 2 \times 1)]=1 \\
& a_{\mathrm{A}_{1}^{\prime \prime}}=\frac{1}{12}[(3 \times 1 \times 1)+(1 \times 1 \times 3)-(3 \times 1 \times 1) \\
& \quad-(1 \times 1 \times 3)]=0
\end{aligned} \quad \begin{aligned}
& \quad+(3 \times 1 \times 3)]=0 \\
& a_{\mathrm{A}_{2}^{\prime \prime}}=\frac{1}{12}[(3 \times 1 \times 1)-(1 \times 1 \times 3)-(3 \times 1 \times 1)
\end{aligned} \quad \begin{aligned}
& a_{\mathrm{E}^{\prime \prime}}=\frac{1}{12}[(3 \times 2 \times 1)-(3 \times 2 \times 1)]=0 \\
& \text { hence }
\end{aligned}
$$

$$
\begin{equation*}
\Gamma_{\sigma}=\mathrm{A}_{1}^{\prime} \oplus \mathrm{E}^{\prime} \tag{6.18}
\end{equation*}
$$

The representation $\Gamma_{\sigma}$ is thus decomposed into a one-dimensional representation ( $A_{1}^{\prime}$, the totally symmetric representation) and a degenerate two-dimensional representation $\left(E^{\prime}\right)$.

### 6.6.3.2. Symmetry-adapted orbitals

As in the previous examples, the orbital with $\mathrm{A}_{1}^{\prime}$ symmetry (the totally symmetric representation) is a linear combination of all the $\sigma_{i}$ orbitals with coefficients that are equal in sign and magnitude. The orbitals with $\mathrm{E}^{\prime}$ symmetry are obtained by applying the projection formula to the functions $\sigma_{1}$ and $\left(\sigma_{2}-\sigma_{3}\right)$, in turn (see Exercise 6.11), as for $\mathrm{NH}_{3}$ (§ 6.4.2.2.). The results are:

$$
\begin{align*}
\phi_{\mathrm{A}_{1}^{\prime}} & =\frac{1}{\sqrt{3}}\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right)  \tag{6.19a}\\
\phi_{\mathrm{E}^{\prime}}(1) & =\frac{1}{\sqrt{6}}\left(2 \sigma_{1}-\sigma_{2}-\sigma_{3}\right)  \tag{6.19b}\\
\phi_{\mathrm{E}^{\prime}}(2) & =\frac{1}{\sqrt{2}}\left(\sigma_{2}-\sigma_{3}\right) \tag{6.19c}
\end{align*}
$$

Figure 6.8. Symmetry-adapted $\sigma$ orbitals for a trigonal-planar ML3 complex.


These three orbitals are shown in Figure 6.8.

### 6.6.4. Trigonal-bipyramidal $\mathrm{ML}_{5}$ complexes

In a trigonal-bipyramidal ML ${ }_{5}$ complex, three of the ligands are found in equatorial positions ( $\mathrm{L}_{1}, \mathrm{~L}_{2}$, and $\mathrm{L}_{3}$ ), whereas the other two ( $\mathrm{L}_{4}$ and


6-29 $\mathrm{L}_{5}$ ) are in axial sites (6-29). The apices of the bipyramid are defined by the axial ligands, its triangular base by the positions of the equatorial ligands. The angles between the equatorial bonds are $120^{\circ}$ and the axial bonds are perpendicular to the equatorial plane. Like the trigonal-planar complex $\mathrm{ML}_{3}$, a trigonal-bipyramidal molecule belongs to the $D_{3 h}$ point group (see 6-28 for the symmetry elements).

### 6.6.4.1. Reduction of the representation $\Gamma_{\sigma}$

The $C_{3}$-axis maintains the positions of the two axial ligands ( $\chi=2$ ), whereas the $S_{3}$-axis, which interchanges them, changes the positions of all the ligands $(\chi=0)$. A $C_{2}$-axis moves every ligand except the one placed on that axis $(\chi=1)$. The $\sigma_{h}$ and $\sigma_{v}$ planes do not move the three ligands found in these planes $(\chi=3)$. The characters of the $\Gamma_{\sigma}$ representation are given in Table 6.22 (second line).

The reduction formula (6.5), in combination with the characters of the irreducible representations of the $D_{3 h}$ point group found in Table 6.21, leads to:

$$
\begin{aligned}
& a_{\mathrm{A}_{1}^{\prime}}=\frac{1}{12}[(5 \times 1 \times 1)+(2 \times 1 \times 2)+(1 \times 1 \times 3)+(3 \times 1 \times 1) \\
& +(3 \times 1 \times 3)]=2 \\
& a_{\mathrm{A}_{2}^{\prime}}=\frac{1}{12}[(5 \times 1 \times 1)+(2 \times 1 \times 2)-(1 \times 1 \times 3)+(3 \times 1 \times 1) \\
& -(3 \times 1 \times 3)]=0 \\
& a_{\mathrm{E}^{\prime}}=\frac{1}{12}[(5 \times 2 \times 1)-(2 \times 1 \times 2)+(3 \times 2 \times 1)]=1 \\
& a_{\mathrm{A}_{1}^{\prime \prime}}=\frac{1}{12}[(5 \times 1 \times 1)+(2 \times 1 \times 2)+(1 \times 1 \times 3)-(3 \times 1 \times 1) \\
& -(3 \times 1 \times 3)]=0 \\
& a_{\mathrm{A}_{2}^{\prime \prime}}=\frac{1}{12}[(5 \times 1 \times 1)+(2 \times 1 \times 2)-(1 \times 1 \times 3)-(3 \times 1 \times 1) \\
& +(3 \times 1 \times 3)]=1 \\
& a_{\mathrm{E}^{\prime \prime}}=\frac{1}{12}[(5 \times 2 \times 1)-(2 \times 1 \times 2)-(3 \times 2 \times 1)]=0
\end{aligned}
$$

Table 6.22. Characters of the reducible representations $\Gamma_{\sigma}, \Gamma_{\sigma}(\mathrm{eq})$ and $\Gamma_{\sigma}(\mathrm{ax})$ of an ML ${ }_{5}$ complex with a TBP geometry

| $D_{3 h}$ | $E$ | $2 C_{3}$ | $3 C_{2}$ | $\sigma_{h}$ | $2 S_{3}$ | $3 \sigma_{v}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
| $\Gamma_{\sigma}$ | 5 | 2 | 1 | 3 | 0 | 3 |
| $\Gamma_{\sigma}(\mathrm{eq})$ | 3 | 0 | 1 | 3 | 0 | 1 |
| $\Gamma_{\sigma}(\mathrm{ax})$ | 2 | 2 | 0 | 0 | 0 | 2 |

hence

$$
\begin{equation*}
\Gamma_{\sigma}=2 \mathrm{~A}_{1}^{\prime} \oplus \mathrm{A}_{2}^{\prime \prime} \oplus \mathrm{E}^{\prime} \tag{6.20}
\end{equation*}
$$

The $\Gamma_{\sigma}$ representation is thus decomposed into three one-dimensional representations $\left(2 \mathrm{~A}_{1}^{\prime}\right.$ and $\left.\mathrm{A}_{2}^{\prime \prime}\right)$ and one doubly degenerate representation ( $\mathrm{E}^{\prime}$ ).

It is important to note that no symmetry operation exchanges an axial ligands with an equatorial one. As these two types of ligands are therefore non-equivalent, both 'chemically' and according to group theory, they can be considered separately. The characters of the representations $\Gamma_{\sigma}(\mathrm{eq})$ and $\Gamma_{\sigma}(\mathrm{ax})$ are given in Table 6.22. From the reduction formula (6.5), we find:

$$
\begin{align*}
& \Gamma_{\sigma}(\mathrm{eq})=\mathrm{A}_{1}^{\prime} \oplus \mathrm{E}^{\prime}  \tag{6.21}\\
& \Gamma_{\sigma}(\mathrm{ax})=\mathrm{A}_{1}^{\prime} \oplus \mathrm{A}_{2}^{\prime \prime} \tag{6.22}
\end{align*}
$$

It is easy to check that $\Gamma_{\sigma}=\Gamma_{\sigma}(\mathrm{eq}) \oplus \Gamma_{\sigma}(\mathrm{ax})$.

### 6.6.4.2. Symmetry-adapted orbitals

The separation of $\Gamma_{\sigma}$ into $\Gamma_{\sigma}(\mathrm{eq})$ and $\Gamma_{\sigma}(\mathrm{ax})$ leads to a considerable simplification of the determination of the symmetry-adapted orbitals. For the orbitals on the equatorial ligands, they are identical to those we have already determined for a trigonal-planar $\mathrm{ML}_{3}$ complex ( $\S$ 6.6.3.2, 6.23a,b and c and Figure 6.9).

$$
\begin{align*}
\left(\phi_{\mathrm{A}_{1}^{\prime}}\right)_{\mathrm{eq}} & =\frac{1}{\sqrt{3}}\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right)  \tag{6.23a}\\
\left(\phi_{\mathrm{E}^{\prime}}\right)_{\mathrm{eq}}(1) & =\frac{1}{\sqrt{2}}\left(\sigma_{2}-\sigma_{3}\right)  \tag{6.23b}\\
\left(\phi_{\mathrm{E}^{\prime}}\right)_{\mathrm{eq}}(2) & =\frac{1}{\sqrt{6}}\left(2 \sigma_{1}-\sigma_{2}-\sigma_{3}\right)  \tag{6.23c}\\
\left(\phi_{\mathrm{A}_{1}^{\prime}}\right)_{\mathrm{ax}} & =\frac{1}{\sqrt{2}}\left(\sigma_{4}+\sigma_{5}\right)  \tag{6.23d}\\
\left(\phi_{\mathrm{A}_{2}^{\prime \prime}}\right)_{\mathrm{ax}} & =\frac{1}{\sqrt{2}}\left(\sigma_{4}-\sigma_{5}\right) \tag{6.23e}
\end{align*}
$$

Figure 6.9. Symmetry-adapted $\sigma$ orbitals for a TBP ML 5 complex.




For the axial ligands, we only need to make normalized sums $\left(A_{1}^{\prime}\right.$ symmetry) and differences ( $\mathrm{A}_{2}^{\prime \prime}$ symmetry) (6.23d, 6.23e and Figure 6.9).

### 6.6.5. Octahedral $\mathrm{ML}_{6}$ complexes

Of the numerous symmetry elements in an octahedral complex $\mathrm{ML}_{6}$ $(6-30)$, the only ones that leave the positions of certain ligands unchanged are shown in $\mathbf{6 - 3 1}$. The three $C_{4}$ axes and their co-linear $C_{2}$ axes both conserve the positions of the two ligands situated on them $(\chi=2)$. Reflection in the three $\sigma_{h}$ planes, which are perpendicular to the $C_{4}$


6-31 axes, does not move the four ligands in these planes $(\chi=4)$. Each of the six $\sigma_{d}$ planes contains only two ligands, so reflection therein gives a character of 2 .

### 6.6.5.1. Reduction of the $\Gamma_{\sigma}$ representation

The characters of the resulting $\Gamma_{\sigma}$ representation are given in Table 6.23 $\left(\Gamma_{\sigma}\right)$, together with those of the irreducible representations of the $O_{h}$ point group.

The order of this group $(h)$ is 48 . Use of the reduction formula (6.5) shows that the only non-zero contributions are:

$$
\begin{aligned}
a_{\mathrm{A}_{1 g}}= & \frac{1}{48}[(6 \times 1 \times 1)+(2 \times 1 \times 6)+(2 \times 1 \times 3)+(2 \times 1 \times 6) \\
& \quad+(4 \times 1 \times 3)]=1 \\
a_{\mathrm{E}_{g}}= & \frac{1}{48}[(6 \times 2 \times 1)+(2 \times 2 \times 3)+(4 \times 2 \times 3)]=1 \\
a_{\mathrm{T}_{1 u}}= & \frac{1}{48}[(6 \times 3 \times 1)+(2 \times 1 \times 6)-(2 \times 1 \times 3)+(2 \times 1 \times 6) \\
& \quad+(4 \times 1 \times 3)]=1
\end{aligned}
$$

Table 6.23. Character table for the $O_{h}$ point group and the characters of the reducible representation $\Gamma_{\sigma}$ for an octahedral $\mathrm{ML}_{6}$ complex

| $O_{h}$ | $E$ | $8 C_{3}$ | $6 C_{2}^{\prime}$ | $6 C_{4}$ | $3 C_{2}$ | $i$ | $8 S_{6}$ | $6 \sigma_{d}$ | $6 S_{4}$ | $3 \sigma_{h}$ |  |
| :--- | :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | :--- |
| $\mathrm{~A}_{1 g}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |  |
| $\mathrm{~A}_{2 g}$ | 1 | 1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 | -1 | $x^{2}+y^{2}+z^{2}$ |
| $\mathrm{E}_{g}$ | 2 | -1 | 0 | 0 | 2 | 2 | -1 | 0 | 0 | 2 |  |
| $\mathrm{~T}_{1 g}$ | 3 | 0 | -1 | 1 | -1 | 3 | 0 | -1 | 1 | -1 |  |
| $\mathrm{~T}_{2 g}$ | 3 | 0 | 1 | -1 | -1 | 3 | 0 | 1 | -1 | -1 |  |
| $\mathrm{~A}_{1 u}$ | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 |  |
| $\left.\mathrm{~A}_{2 u}, x^{2}-y^{2}\right)$ |  |  |  |  |  |  |  |  |  |  |  |
| $\mathrm{E}_{u}$ | 1 | 1 | -1 | -1 | 1 | -1 | -1 | 1 | 1 | -1 |  |
| $\mathrm{~T}_{1 u}$ | 3 | -1 | 0 | 0 | 2 | -2 | 1 | 0 | 0 | -2 |  |
| $\mathrm{~T}_{2 u}$ | 3 | 0 | -1 | 1 | -1 | -3 | 0 | 1 | -1 | 1 | $(x y, x z, y z)$ |
| $\Gamma_{\sigma}$ | 6 | 0 | 0 | 2 | 2 | 0 | 0 | 2 | 0 | 4 |  |

hence

$$
\begin{equation*}
\Gamma_{\sigma}=\mathrm{A}_{1 g} \oplus \mathrm{E}_{g} \oplus \mathrm{~T}_{1 u} \tag{6.24}
\end{equation*}
$$

The $\Gamma_{\sigma}$ representation is therefore decomposed into a one-dimensional representation $\left(\mathrm{A}_{1 g}\right)$, a doubly degenerate representation $\left(\mathrm{E}_{g}\right)$, and a triply degenerate representation $\left(\mathrm{T}_{1 u}\right)$.

### 6.6.5.2. Symmetry-adapted orbitals

In view of the large number of symmetry operations, we shall limit ourselves here to giving the result of the application of the projection formula (6.10). The following orbitals are obtained; they are shown in Figure 6.10:

$$
\begin{align*}
\phi_{\mathrm{A}_{1 g}} & =\frac{1}{\sqrt{6}}\left(\sigma_{1}+\sigma_{2}+\sigma_{3}+\sigma_{4}+\sigma_{5}+\sigma_{6}\right)  \tag{6.25a}\\
\phi_{\mathrm{E}_{g}}(1) & =\frac{1}{2}\left(\sigma_{1}-\sigma_{2}+\sigma_{3}-\sigma_{4}\right)  \tag{6.25b}\\
\phi_{\mathrm{E}_{g}}(2) & =\frac{1}{\sqrt{12}}\left(-\sigma_{1}-\sigma_{2}-\sigma_{3}-\sigma_{4}+2 \sigma_{5}+2 \sigma_{6}\right)  \tag{6.25c}\\
\phi_{\mathrm{T}_{1 u}}(1) & =\frac{1}{\sqrt{2}}\left(\sigma_{1}-\sigma_{3}\right)  \tag{6.25~d}\\
\phi_{\mathrm{T}_{1 u}}(2) & =\frac{1}{\sqrt{2}}\left(\sigma_{2}-\sigma_{4}\right)  \tag{6.25e}\\
\phi_{\mathrm{T}_{1 u}}(3) & =\frac{1}{\sqrt{2}}\left(\sigma_{5}-\sigma_{6}\right) \tag{6.25f}
\end{align*}
$$

Figure 6.10. Symmetry-adapted $\sigma$ orbitals in an octahedral $\mathrm{ML}_{6}$ complex.


The degeneracy of the $T_{1 u}$ orbitals is obvious from the coefficients of the $\sigma_{i}$ orbitals (6.25d-f), but the same is not true of the $\mathrm{E}_{\mathrm{g}}$ orbitals, whose coefficients are very different ( $6.25 \mathrm{~b}, \mathrm{c}$ ).

### 6.6.6. Trigonal-planar $\mathrm{ML}_{3}$ complexes with a ' $\pi$ system' on the ligands

In this last example, we shall analyse a trigonal-planar $\mathrm{ML}_{3}$ complex in which the ligands are considered to have two $p$ orbitals perpendicular to the $\mathrm{M}-\mathrm{L}$ bond as well as the $\sigma$ orbital that points towards the metal (§ 6.6.3.). These $p$ orbitals are written $p_{\|}(6-32)$ and $p_{\perp}(\mathbf{6 - 3 3})$, depending on whether they are in the plane of the complex or perpendicular to it. By convention, each orbital $p_{\| i}$ is oriented in a clockwise sense ( $-p_{\| i}$ corresponds to the opposite orientation). This system can act as a model for a complex in which the three ligands are double-face $\pi$ donors or $\pi$ acceptors (Chapter 3).

### 6.6.6.1. Reduction of the representations $\Gamma_{p \|}$ and $\Gamma_{p \perp}$

To obtain the characters of the representations $\Gamma_{p \|}$ and $\Gamma_{p \perp}$, it is now not sufficient just to count the number of ligands whose position is unchanged by the action of the symmetry operations of the $\left(D_{3 h}\right)$ point group. Although that procedure was acceptable in the previous examples, the situation is more complicated now, since an orbital can be transformed into itself ( $\chi=1$ ) or into its opposite $(\chi=-1)$.

We consider first the $\Gamma_{p \|}$ representation. Action of the $C_{3}$ and $S_{3}$ axes changes all three orbitals $(\chi=0)$, a $C_{2}$-axis transforms the orbital of the atom on that axis into its opposite and also interchanges the two others while changing their sign (6-34 for the axis that passes through
the ligand $L_{1}$ ). The associated character is therefore -1 . The molecular plane $\sigma_{h}$ maintains the three orbitals $(\chi=3)$ in their original positions, whereas reflection in a $\sigma_{v}$ plane has the same consequences as rotation around a $C_{2}$-axis ( $6-35, \chi=-1$ ).


6-34



6-35

We now turn to the $\Gamma_{p \perp}$ representation. The character is still zero for the $C_{3}$ and $S_{3}$ axes, and still -1 for a $C_{2}$-axis (6-36). Reflection in the $\sigma_{h}$ plane changes each orbital into its opposite ( $6-37, \chi=-3$ ), whereas a $\sigma_{v}$ plane maintains one orbital and interchanges the two others (6-38, $\chi=1$ ).




6-36

The characters associated with the $\Gamma_{p \|}$ and $\Gamma_{p \perp}$ representations are listed in Table 6.24, together with the characters of the irreducible representations of the $D_{3 h}$ point group.



Table 6.24. Character table for the $D_{3 h}$ point group and the characters of the reducible representations $\Gamma_{p \|}$ and $\Gamma_{p \perp}$ for a trigonal-planar $\mathrm{ML}_{3}$ complex
$\left.\left.\begin{array}{lrrrrrrl}\hline D_{3 h} & E & 2 C_{3} & 3 C_{2} & \sigma_{h} & 2 S_{3} & 3 \sigma_{v} & \\ \hline \mathrm{~A}_{1}^{\prime} & 1 & 1 & 1 & 1 & 1 & 1 & \\ \mathrm{~A}_{2}^{\prime} & 1 & 1 & -1 & 1 & 1 & -1 & \\ \mathrm{E}^{\prime} & 2 & -1 & 0 & 2 & -1 & 0 & (x, y) \\ \mathrm{A}_{1}^{\prime \prime} & 1 & 1 & 1 & -1 & -1 & -1 & \\ \mathrm{~A}_{2}^{\prime \prime} & 1 & 1 & -1 & -1 & -1 & 1 & z\end{array}\right] z^{2}-y^{2}, x y\right)$

For the $\Gamma_{p \|}$ representation, application of the reduction formula (6.5) gives:

$$
\begin{aligned}
& a_{\mathrm{A}_{1}^{\prime}}=\frac{1}{12}[(3 \times 1 \times 1)-(1 \times 1 \times 3)+(3 \times 1 \times 1)-(1 \times 1 \times 3)]=0 \\
& a_{\mathrm{A}_{2}^{\prime}}=\frac{1}{12}[(3 \times 1 \times 1)+(1 \times 1 \times 3)+(3 \times 1 \times 1)+(1 \times 1 \times 3)]=1 \\
& a_{\mathrm{E}^{\prime}}=\frac{1}{12}[(3 \times 2 \times 1)+(3 \times 2 \times 1)]=1 \\
& a_{\mathrm{A}_{1}^{\prime \prime}}=\frac{1}{12}[(3 \times 1 \times 1)-(1 \times 1 \times 3)-(3 \times 1 \times 1)+(1 \times 1 \times 3)]=0 \\
& a_{\mathrm{A}_{2}^{\prime \prime}}=\frac{1}{12}[(3 \times 1 \times 1)+(1 \times 1 \times 3)-(3 \times 1 \times 1)-(1 \times 1 \times 3)]=0 \\
& a_{\mathrm{E}^{\prime \prime}}=\frac{1}{12}[(3 \times 2 \times 1)-(3 \times 2 \times 1)]=0
\end{aligned}
$$

hence

$$
\begin{equation*}
\Gamma_{p \|}=\mathrm{A}_{2}^{\prime} \oplus \mathrm{E}^{\prime} \tag{6.26}
\end{equation*}
$$

For the $\Gamma_{p \perp}$ representation, we obtain:
$a_{\mathrm{A}_{1}^{\prime}}=\frac{1}{12}[(3 \times 1 \times 1)-(1 \times 1 \times 3)-(3 \times 1 \times 1)+(1 \times 1 \times 3)]=0$
$a_{\mathrm{A}_{2}^{\prime}}=\frac{1}{12}[(3 \times 1 \times 1)+(1 \times 1 \times 3)-(3 \times 1 \times 1)-(1 \times 1 \times 3)]=0$
$a_{\mathrm{E}^{\prime}}=\frac{1}{12}[(3 \times 2 \times 1)-(3 \times 2 \times 1)]=0$
$a_{\mathrm{A}_{1}^{\prime \prime}}=\frac{1}{12}[(3 \times 1 \times 1)-(1 \times 1 \times 3)+(3 \times 1 \times 1)-(1 \times 1 \times 3)]=0$
$a_{\mathrm{A}_{2}^{\prime \prime}}=\frac{1}{12}[(3 \times 1 \times 1)+(1 \times 1 \times 3)+(3 \times 1 \times 1)+(1 \times 1 \times 3)]=1$
$a_{\mathrm{E}^{\prime \prime}}=\frac{1}{12}[(3 \times 2 \times 1)+(3 \times 2 \times 1)]=1$
hence

$$
\begin{equation*}
\Gamma_{p \perp}=\mathrm{A}_{2}^{\prime \prime} \oplus \mathrm{E}^{\prime \prime} \tag{6.27}
\end{equation*}
$$

### 6.6.6.2. Symmetry-adapted orbitals

We shall limit ourselves here to the determination of the symmetryadapted $\mathrm{A}_{2}^{\prime \prime}$ and $\mathrm{E}^{\prime \prime}$ orbitals in the $\Gamma_{p \perp}$ representation. Table 6.25 shows the action of all the symmetry operations of the $D_{3 h}$ point group on the generating functions $p_{\perp 1}$ and $\left(p_{\perp 2}-p_{\perp 3}\right)$, as well as the characters of the $A_{2}^{\prime \prime}$ and $\mathrm{E}^{\prime \prime}$ irreducible representations.

The symmetry-adapted orbitals are obtained from formula (6.10), by using the $p_{\perp_{1}}$ generating function for the $\phi_{\mathrm{A}_{2}^{\prime \prime}}$ orbital and for the first degenerate orbital $\phi_{\mathrm{E}^{\prime \prime}}(1)$, and the function ( $p_{\perp 2}-p_{\perp_{3}}$ ) for the second degenerate orbital $\phi_{\mathrm{E}^{\prime \prime}}(2)$. We obtain:

$$
\begin{aligned}
\phi_{\mathrm{A}_{2}^{\prime \prime}} & =4 \times\left(p_{\perp 1}+p_{\perp 2}+p_{\perp 3}\right) \\
\phi_{\mathrm{E}^{\prime \prime}}(1) & =2 \times\left(2 p_{\perp 1}-p_{\perp 2}-p_{\perp 3}\right) \\
\phi_{\mathrm{E}^{\prime \prime}}(2) & =6 \times\left(p_{\perp 2}-p_{\perp 3}\right)
\end{aligned}
$$

so the normalized orbitals are

$$
\begin{align*}
\phi_{A_{2}^{\prime \prime}} & =\frac{1}{\sqrt{3}}\left(p_{\perp 1}+p_{\perp 2}+p_{\perp 3}\right)  \tag{6.28a}\\
\phi_{\mathrm{E}^{\prime \prime}}(1) & =\frac{1}{\sqrt{6}}\left(2 p_{\perp 1}-p_{\perp 2}-p_{\perp 3}\right)  \tag{6.28b}\\
\phi_{\mathrm{E}^{\prime \prime}}(2) & =\frac{1}{\sqrt{2}}\left(p_{\perp 2}-p_{\perp 3}\right) \tag{6.28c}
\end{align*}
$$

Table 6.25. Action of the symmetry operations of the $D_{3 h}$ point group on the $p_{\perp}$ orbitals (see 6-33) and the characters of the $\mathrm{A}_{2}^{\prime \prime}$ and $\mathrm{E}^{\prime \prime}$ irreducible representations. [See Scheme 6-28 for the definition of the symmetry elements. A $C_{2}(i)$-axis is co-linear with the $\mathrm{M}-\mathrm{L}_{i}$ bond, and a $\sigma_{v}(i)$ plane contains the $\mathrm{M}-\mathrm{L}_{i}$ bond.]

| $D_{3 h}$ | E | $\mathrm{C}_{3}^{1}$ | $\mathrm{C}_{3}^{2}$ | $\mathrm{C}_{2}(1)$ | $\mathrm{C}_{2}(2)$ | $\mathrm{C}_{2}(3)$ | $\sigma_{h}$ | $\mathrm{~S}_{3}^{1}$ | $\mathrm{~S}_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $R_{k}\left(p_{\perp 1}\right)$ | $p_{\perp 1}$ | $p_{\perp 2}$ | $p_{\perp 3}$ | $-p_{\perp 1}$ | $-p_{\perp 3}$ | $-p_{\perp 2}$ | $-p_{\perp 1}$ | $-p_{\perp 2}$ | $-p_{\perp 3}(3)$ | $p_{\perp 1}$ | $p_{\perp 3}$ |
| $R_{k}\left(p_{\perp 2}-p_{\perp 3}\right)$ | $p_{\perp 2}-p_{\perp 3}$ | $p_{\perp 1}-p_{\perp 2}$ | $p_{\perp 3}-p_{\perp 1}$ | $p_{\perp 2}-p_{\perp 3}$ | $p_{\perp 1}-p_{\perp 2}$ | $p_{\perp 3}-p_{\perp 1}$ | $p_{\perp 3}-p_{\perp 2}$ | $p_{\perp 2}-p_{\perp 1}$ | $p_{\perp 1}-p_{\perp 3}$ | $p_{\perp 3}-p_{\perp 2}$ | $p_{\perp 2}-p_{\perp 1}$ |
| $\mathrm{~A}_{2}^{\prime \prime}$ | 1 | 1 | 1 | -1 | -1 | -1 | -1 | -1 | -1 | 1 | 1 |
| $\mathrm{E}^{\prime \prime}$ | 2 | -1 | -1 | 0 | 0 | 0 | -2 | 1 | 1 | 0 | 0 |

Figure 6.11. $\pi_{\|}$and $\pi_{\perp}$ symmetry-adapted orbitals in a trigonal-planar $\mathrm{ML}_{3}$ complex.

$\pi_{/ /}\left(\mathrm{A}_{2}^{\prime}\right)$

$\pi_{\perp}\left(\mathrm{A}_{2}^{\prime \prime}\right)$

$\pi_{/ /}\left(\mathrm{E}^{\prime}\right)$



For the $\mathrm{A}_{2}^{\prime}$ and $\mathrm{E}^{\prime \prime}$ symmetry-adapted orbitals $\left(\Gamma_{p \|}\right)$, we obtain (see Exercise 6.12):

$$
\begin{align*}
\phi_{\mathrm{A}_{2}^{\prime}} & =\frac{1}{\sqrt{3}}\left(p_{\| 1}+p_{\| 2}+p_{\| 3}\right)  \tag{6.29a}\\
\phi_{\mathrm{E}^{\prime}}(1) & =\frac{1}{\sqrt{6}}\left(2 p_{\| 1}-p_{\| 2}-p_{\| 3}\right)  \tag{6.29b}\\
\phi_{\mathrm{E}^{\prime}}(2) & =\frac{1}{\sqrt{2}}\left(p_{\| 2}-p_{\| 3}\right) \tag{6.29c}
\end{align*}
$$

The different symmetry-adapted orbitals are shown in Figure 6.11.
Although the $\sigma$-type orbitals point along the $\mathrm{M}-\mathrm{L}$ bonds, the orbitals above are constructed from ligand orbitals that are perpendicular to these bonds. This is why they are often written $\pi_{\|}$and $\pi_{\perp}$, even though this notation is not strictly correct according to group theory.

## Exercises

## Symmetry elements and symmetry operations

## 6.1

How many reflection planes are there in the following molecules? (1) water $\left(\mathrm{H}_{2} \mathrm{O}\right)$; (2) ammonia $\left(\mathrm{NH}_{3}\right)$; (3) ethylene $\left(\mathrm{C}_{2} \mathrm{H}_{4}\right)$; (4) (Z)-1,2-difluoroethylene; (5) (E)-1,2-difluoroethylene; (6) aluminium trichloride $\left(\mathrm{AlCl}_{3}\right.$, a trigonal-planar molecule).

## 6.2

Which of the following molecules contain an inversion centre? (1) carbon dioxide ( $\mathrm{CO}_{2}$, linear); (2) hydrogen cyanide (HCN, linear); (3) dimethylether (shown in 6-1); (4) benzene $\left(\mathrm{C}_{6} \mathrm{H}_{6}\right.$, hexagonal);
(5) (Z)-1,2-difluoroethylene; (6) (E)-1,2-difluoroethylene; (7) aluminium trichloride $\left(\mathrm{AlCl}_{3}\right.$, a trigonal-planar molecule); (8) methane $\left(\mathrm{CH}_{4}\right.$, tetrahedral); (9) $\left[\mathrm{PtCl}_{4}\right]^{2-}$ (a square-planar complex, Scheme 6-10); (10) ethane $\left(\mathrm{C}_{2} \mathrm{H}_{6}\right)$ in the staggered conformation; (11) ethane in the eclipsed conformation.

## 6.3

1. Are there any rotation axes in the ammonia molecule (6-7) in addition to the $C_{3}$-axis?
2. Locate the three $C_{3}$ axes and the four $C_{2}$ axes in the methane molecule.
3. How many $C_{2}$ axes are there in (i) ethylene; (ii) ( $Z$ )-1,2difluoroethylene; (iii) (E)-1,2-difluoroethylene ?
4. Which rotation axes are present in $\mathrm{AlCl}_{3}$ ?

## 6.4

1. Consider a $C_{2}$-axis and the co-linear $S_{4}$-axis in methane (see Scheme 6-11). Show that $S_{4}^{2}=C_{2}$ and that $S_{4}^{4}=E$.
2. Carry out the $S_{3}^{2}$ and $S_{3}^{4}$ operations in the $\mathrm{PF}_{5}$ molecule (see Scheme 6-12). Show that these operations are equivalent to operations associated with the $C_{3}$-axis or the $\sigma_{h}$ plane.
3. (i) Show that ethane $\left(\mathrm{C}_{2} \mathrm{H}_{6}\right)$ in the eclipsed conformation possesses an improper axis that is co-linear with the $\mathrm{C}-\mathrm{C}$ bond. What is the order of this axis? (ii) Repeat the question for ethane in the staggered conformation.
6.5

The symmetry elements present in the ammonia molecule $\left(\mathrm{NH}_{3}\right)$ are a $C_{3}$-axis (6-7) and three planes of symmetry $\left(\sigma_{v}(1), \sigma_{v}(2)\right.$, and $\left.\sigma_{v}(3)\right)($ Exercise 6.1, question 2).

1. Perform the $C_{3}$ operation followed by $\sigma_{v}(1)$ (written $\sigma_{v}(1) C_{3}$ ); carry out $\sigma_{v}(1)$ followed by $C_{3}$ (written $\left.C_{3} \sigma_{v}(1)\right)$. What conclusion(s) can you draw?
2. Perform $\left(C_{3} \sigma_{v}(1)\right) \sigma_{v}(2)$ then $C_{3}\left(\sigma_{v}(1) \sigma_{v}(2)\right)$. What conclusion(s) can you draw?
3. Which operations are the inverses of $E, C_{3}, C_{3}^{2}, \sigma_{v}(1)$ ?
6.6

Determine the point-group symmetries of the following molecules:
(1) $\mathrm{O}_{2}$; (2) HCl ; (3) ethylene $\left(\mathrm{C}_{2} \mathrm{H}_{4}\right)$; (4) (Z)-1,2-difluoroethylene;
(5) (E)-1,2-difluoroethylene; (6) $\mathrm{AlCl}_{3}$, a trigonal-planar molecule;
(7) tetrachloromethane $\left(\mathrm{CCl}_{4}\right)$; (8) dichloromethane $\left(\mathrm{CH}_{2} \mathrm{Cl}_{2}\right)$; (9)
$\mathrm{NH}_{3}(6-7)$; (10) $\left.\left[\mathrm{PtCl}_{4}\right]^{2-}(\mathbf{6 - 1 0}) ;(11) \mathrm{PF}_{5}(\mathbf{6 - 1 2}) ; 12\right)$ trifluoroethylene ( $\mathrm{F}_{2} \mathrm{C}=\mathrm{CHF}$, planar).

## Symmetry-adapted orbitals

## 6.7

Show that in the $\mathrm{NH}_{3}$ molecule, the functions $(1 / \sqrt{3})\left(1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}+\right.$ $\left.1 s_{\mathrm{H}_{3}}\right)$, a basis for the $\mathrm{A}_{1}$ representation, and $(1 / \sqrt{6})\left(2 \times\left(1 s_{\mathrm{H}_{1}}\right)-\right.$ $1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}$ ), one of the basis functions for the E representation, are orthogonal. The $1 s_{\mathrm{H}}$ orbitals are assumed to be normalized, and the overlap between two $1 s_{\mathrm{H}}$ orbitals on different atoms is written $S$.

## 6.8

1. Reduce the $\Gamma_{\mathrm{H}}$ basis that is constituted by the $1 s_{\mathrm{H}}$ orbitals on the hydrogen atoms in the ethylene molecule (the characters of this representation are given in Table 6.8, and the character table for the $D_{2 h}$ point group is given below).
2. Find the linear combinations of these orbitals that are bases for irreducible representations.

| $D_{2 h}$ | $E$ | $\mathrm{C}_{2}(z)$ | $C_{2}(y)$ | $C_{2}(x)$ | $i$ | $\sigma(x y)$ | $\sigma(x z)$ | $\sigma(y z)$ |  |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | :--- | :--- |
| $\mathrm{A}_{g}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | $x^{2}, y^{2}, z^{2}$ |
| $\mathrm{~B}_{1 g}$ | 1 | 1 | -1 | -1 | 1 | 1 | -1 | -1 | $x y$ |
| $\mathrm{~B}_{2 g}$ | 1 | -1 | 1 | -1 | 1 | -1 | 1 | -1 | $x z$ |
| $\mathrm{~B}_{3 g}$ | 1 | -1 | -1 | 1 | 1 | -1 | -1 | 1 | $y z$ |
| $\mathrm{~A}_{u}$ | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 |  |
| $\mathrm{~B}_{1 u}$ | 1 | 1 | -1 | -1 | -1 | -1 | 1 | 1 | $z$ |
| $\mathrm{~B}_{2 u}$ | 1 | -1 | 1 | -1 | -1 | 1 | -1 | 1 | $y$ |
| $\mathrm{~B}_{3 u}$ | 1 | -1 | -1 | 1 | -1 | 1 | 1 | -1 | $x$ |

## 6.9

Construct the molecular orbitals of $\mathrm{NH}_{3}$ from the symmetryadapted orbitals that are given in Figure 6.3. To help you construct the interaction diagram, note that the energies of the orbitals on $\mathrm{H}_{3}$ fragment and of the nitrogen $2 p_{\mathrm{N}}$ orbitals are close to -13.5 eV , and that the energy of the $2 s_{\mathrm{N}}$ orbital is -26.0 eV .
6.10

Calculate the coefficients of $\sigma_{1}, \sigma_{2}, \sigma_{3}$, and $\sigma_{4}$ (see below for the numbering of the atoms) in the symmetry-adapted $\mathrm{T}_{2}$ orbitals given on the right-hand side of Figure 6.7. You should use the normalization relationships and the fact that these orbitals are mutually orthogonal and orthogonal to the symmetry-adapted $\mathrm{A}_{1}$ orbital.

### 6.11

In a trigonal-planar $\mathrm{ML}_{3}$ complex, the $\Gamma_{\sigma}$ representation is reduced thus: $\Gamma_{\sigma}=\mathrm{A}_{1}^{\prime} \oplus \mathrm{E}^{\prime}$ (§ 6.6.3.1, formula (6.18)).

1. Characterize the $\mathrm{A}_{1}^{\prime}$ symmetry-adapted orbital, using $\sigma_{1}$ as the generating function.
2. Characterize the $\mathrm{E}^{\prime}$ symmetry-adapted orbitals, using $\sigma_{1}$ followed by $\left(\sigma_{2}-\sigma_{3}\right)$ as the generating functions. The projection formula (6.10) and the character table for the $D_{3 h}$ point group (Table 6.21) should be helpful.

### 6.12

In a trigonal-planar $\mathrm{ML}_{3}$ complex, the $\Gamma_{p \|}$ representation is reduced thus:
$\Gamma_{p \|}=\mathrm{A}_{2}^{\prime} \oplus \mathrm{E}^{\prime}$ (§ 6.6.6.1, formula (6.26)).

1. Characterize the $\mathrm{A}_{2}^{\prime}$ symmetry-adapted orbital, using $p_{\| 1}$ as the generating function.
2. Characterize the $\mathrm{E}^{\prime}$ symmetry-adapted orbitals, using $p_{\| 1}$ followed by $\left(p_{\| 2}-p_{\| 3}\right)$ as the generating functions.

The projection formula (6.10) and the character table for the $D_{3 h}$ point group (Table 6.24) should be helpful. The complete set of symmetry operations is given in Table 6.25.

### 6.13

Consider a square-planar $\mathrm{ML}_{4}$ complex in which each ligand possesses a $\pi$ system made up of two $p$ orbitals that are perpendicular to the $\mathrm{M}-\mathrm{L}$ bonds. These orbitals are $p_{\perp}$ and $p_{\|}$, depending on whether they are perpendicular to the plane of the complex or in that plane. The orientation of these orbitals is shown below.


Tables 6.18 and 6.19 , which give the character table for the $D_{4 h}$ point group and the full set of the symmetry operations of this group, should be helpful, as should Scheme 6-24 in which the different symmetry elements are shown.

1. Determine the characters of the $\Gamma_{p \perp}$ representation.
2. Decompose this representation into a sum of irreducible representations of the $D_{4 h}$ point group.
3. Characterize the symmetry-adapted orbitals in each of these representations. The generating function $p_{\perp_{1}}$ should be used for the one-dimensional representations, and the generating functions $p_{\perp 1}$ and $p_{\perp 2}$ in turn for the two-dimensional representation.
4. Sketch the resulting orbitals.
5. From an analysis of the interactions between the ligand orbitals, establish the energetic ordering of the symmetry-adapted MO.
6. Repeat questions $1-5$ for the $\Gamma_{p \|}$ representation.

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# Answers to the exercises (in skeletal form) 

## Chapter 1: Setting the scene

1.1. The coordination modes are $\eta^{1}$ (X-type ligand) and $\eta^{3}$ (LX-type ligand).

$\eta^{1}$-allyl

$\eta^{3}$-allyl
1.2.

|  | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 | 11 | 12 | 13 | 14 | 15 | 16 | 17 | 18 | 19 | 20 | 21 | 22 | 23 | 24 |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $l$ | 6 | 5 | 5 | 0 | 3 | 5 | 4 | 5 | 0 | 2 | 4 | 0 | 0 | 2 | 1 | 0 | 4 | 6 | 4 | 4 | 4 | 0 | 5 | 2 |
| $x$ | 0 | 0 | 1 | 4 | 1 | 1 | 2 | 0 | 9 | 7 | 0 | 3 | 5 | 4 | 4 | 8 | 4 | 0 | 4 | 5 | 2 | 6 | 0 | 2 |
| $q$ | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | -2 | 0 | 0 | -2 | -3 | -1 | -1 | -2 | +1 | +1 | 0 | 0 | 0 | 0 | 0 | 0 |
| no | 0 | 0 | 1 | 4 | 1 | 1 | 2 | 0 | 7 | 7 | 0 | 1 | 2 | 3 | 3 | 6 | 5 | 1 | 4 | 5 | 2 | 6 | 0 | 2 |
| $d^{n}$ | $d^{6}$ | $d^{6}$ | $d^{6}$ | $d^{0}$ | $d^{8}$ | $d^{6}$ | $d^{6}$ | $d^{8}$ | $d^{0}$ | $d^{0}$ | $d^{10}$ | $d^{10}$ | $d^{8}$ | $d^{6}$ | $d^{6}$ | $d^{0}$ | $d^{2}$ | $d^{6}$ | $d^{0}$ | $d^{0}$ | $d^{6}$ | $d^{0}$ | $d^{8}$ | $d^{10}$ |
| $N_{t}$ | 18 | 16 | 18 | 8 | 16 | 18 | 18 | 18 | 18 | 18 | 18 | 16 | 18 | 18 | 18 | 16 | 18 | 18 | 16 | 18 | 18 | 12 | 18 | 18 |

1.3. $n o=1, d^{8}, N_{t}=16 ; n o=3, d^{6}, N_{t}=18 ; n o=3, d^{6}, N_{t}=16$; no $=3, d^{6}, N_{t}=18 ; n o=3, d^{6}, N_{t}=16 ; n o=3, d^{6}, N_{t}=18 ;$ $n o=1, d^{8}, N_{t}=16$.
1.4. (1) $\mathrm{CO}, \mathrm{Cl}^{-}, \mathrm{Et}^{-}, \mathrm{PR}_{3}, \mathrm{H}^{-}, \mathrm{H}_{2}, \mathrm{SiR}_{3}^{-}, \mathrm{SR}^{-}, \mathrm{CN}^{-}, \mathrm{I}^{-}, \mathrm{Me}^{-}$, $\mathrm{COMe}^{-}, \mathrm{F}^{-}, \mathrm{O}^{2-}, \mathrm{NR}_{3}, \mathrm{C}_{2} \mathrm{H}_{4}, \mathrm{C}_{6} \mathrm{H}_{6}, \mathrm{C}_{5} \mathrm{H}_{5}^{-}$.
(2)

| Complex | Ionic model | no |
| :--- | :--- | :--- |
| $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$ | $\left[\mathrm{Fe}(\mathrm{CO})_{5}\right]$ | 0 |
| $\left[\mathrm{Ir}(\mathrm{CO})(\mathrm{Cl})\left(\mathrm{PPh}_{3}\right)_{2}\right]$ | $\left[\mathrm{Ir}^{+}(\mathrm{CO})\left(\mathrm{Cl}^{-}\right)\left(\mathrm{PPh}_{3}\right)_{2}\right]$ | 1 |
| $\left[\mathrm{Mn}(\mathrm{CO})_{6}\right]^{+}$ | $\left[\mathrm{Mn}^{+}(\mathrm{CO})_{6}\right]$ | 1 |
| $\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-}$ | $\left.\left[\mathrm{Ni}^{2+}(\mathrm{CN})_{5}\right)_{5}\right]$ | 2 |
| $\left[\mathrm{Zn}(\mathrm{Cl})_{4}\right]^{2-}$ | $\left[\mathrm{Zn}^{2+}\left(\mathrm{Cl}^{-}\right)_{4}\right]$ | 2 |
| $\left[\mathrm{~V}(\mathrm{Cl})_{4}\right]$ | $\left[\mathrm{V}^{4+}\left(\mathrm{Cl}^{-}\right)_{4}\right]$ | 4 |
| $\left[\mathrm{Cr}(\mathrm{CO})_{3}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\right]$ | $\left[\mathrm{Cr}(\mathrm{CO})_{3}\left(\eta^{6}-\mathrm{C}_{6} \mathrm{H}_{6}\right)\right]$ | 0 |
| $\left[\mathrm{Fe}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\right]$ | $\left[\mathrm{Fe}^{2+}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}^{-}\right)_{2}\right]$ | 2 |
| $\left[\mathrm{Cu}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)\left(\mathrm{PR}_{3}\right)\right]$ | $\left[\mathrm{Cu}^{+}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}^{-}\right)\left(\mathrm{PR}_{3}\right)\right]$ | 1 |
| $\left[\mathrm{Zr}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}\right)_{2}\left(\mathrm{CH}_{3}\right)\right]^{+}$ | $\left[\mathrm{Zr}^{4+}\left(\eta^{5}-\mathrm{C}_{5} \mathrm{H}_{5}^{-}\right)_{2}\left(\mathrm{CH}_{3}^{-}\right)\right]$ | 4 |
| $\left[\mathrm{Ti}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})_{3}\left(\mathrm{CH}_{3}\right)\right]$ | $\left[\mathrm{Ti}^{4+}\left(\mathrm{PR}_{3}\right)_{2}\left(\mathrm{Cl}^{-}\right)_{3}\left(\mathrm{CH}_{3}^{-}\right)\right]$ | 4 |
| $\left[{\left.\mathrm{~W}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{CO})_{3}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]}^{\left[\mathrm{W}^{2}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{CO})_{3}\left(\eta^{2}-\mathrm{H}_{2}\right)\right]}\right.$ | 0 |  |
| $\left[\operatorname{Ir}\left(\mathrm{PR}_{3}\right)_{2}(\mathrm{Cl})(\mathrm{H})_{2}\right]$ | $\left[\mathrm{Ir}^{3+}\left(\mathrm{PR}_{3}\right)_{2}\left(\mathrm{Cl}^{-}\right)\left(\mathrm{H}^{-}\right)_{2}\right]$ | 3 |
| $\left[\mathrm{Ni}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ | $\left[\mathrm{Ni}^{2+}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]$ | 2 |
|  |  |  |

1.5.

1. $\left[\mathrm{CrL}_{6}\right], n o=0, d^{6}, N_{t}=18 ;\left[\mathrm{RuL}_{5}, n o=0, d^{8}, N_{t}=18\right.$.
2. The two compounds are 18 -electron complexes. An $\eta^{6}$ coordination of the two ligands in the ruthenium complex would lead to a 20 -electron organometallic (and thus strong-field) complex.
3. (i) $\eta^{5}$; (ii) $\eta^{3}$; (iii) $\eta^{1}$.
1.6. The borohydride ligand is of L -type in $\eta^{1}$-coordination, $\mathrm{L}_{2}$-type in $\eta^{2}$-coordination, and $\mathrm{L}_{3}$-type in $\eta^{3}$-coordination. In each case, the coordination mode of this ligand allows an 18 -electron complex to be formed.
1.7. (1) no $=0$; (2) no $=3$; (3) no $=2$; (4) no $=2$ (each bridging Cl acts as an X-type ligand towards one metallic centre and as an L-type ligand towards the other); (5) no $=2$ (SR = ligand of XL type (or LX), like Cl).
1.8.
4. The $s$ orbitals have spherical symmetry, so their overlap is nonzero for any relative positions of the metal and the ligand. The bonding orbital is mainly concentrated on the ligand and the antibonding orbital on the metal.


5. (i) On the x-axis; (ii) in the nodal plane $y z: S=0$.
6. $x y$ : the overlap is largest when the ligand is placed on the bisectors of the $x$ - and $y$-axes, and smallest if it is in the $x z$ or $y z$ nodal planes $(S=0)$; $z^{2}$ : largest overlap along the $z$-axis and smallest ( $S=0$ ) when the ligand is on the nodal cone.


Chapter 2: Principal ligand fields: $\sigma$ interactions
2.1. $1-3$

| M | Sc | Ti | V | Cr | Mn | Fe | Co | Ni | Cu | Zn |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $d^{n}$ | $d^{1}$ | $d^{2}$ | $d^{3}$ | $d^{4}$ | $d^{5}$ | $d^{6}$ | $d^{7}$ | $d^{8}$ | $d^{9}$ | $d^{10}$ |
| Electronic configuration | $\mathrm{t}_{2 g}^{1} \mathrm{e}_{g}^{0}$ | $\mathrm{t}_{2 g}^{2} \mathrm{e}_{g}^{0}$ | $\mathrm{t}_{2 g}^{3} \mathrm{e}_{g}^{0}$ | $\mathrm{t}_{2 g}^{3} \mathrm{e}_{g}^{1}$ | $\mathrm{t}_{2 g}^{3} \mathrm{e}_{g}^{2}$ | $\mathrm{t}_{2 g}^{4} \mathrm{e}_{g}^{2}$ | $\mathrm{t}_{2 g}^{5} \mathrm{e}_{g}^{2}$ | $\mathrm{t}_{2 g}^{6} \mathrm{e}_{g}^{2}$ | $\mathrm{t}_{2 g}^{6} \mathrm{e}_{g}^{3}$ | $\mathrm{t}_{2 g}^{6} \mathrm{e}_{g}^{4}$ |
| Unpaired electrons | 1 | 2 | 3 | 4 | 5 | 4 | 3 | 2 | 1 | 0 |

2.2.
(1) $d^{6}$;
(2) Strong-field complex (organometallic): $\mathrm{t}_{2 g}^{6}$;
(3) The anion would be a 19-electron complex $\left(d^{7}\right)$ with a very high-energy $d \mathrm{MO}$ (strong field) containing an electron. Dissociation gives a 17-electron $\mathrm{ML}_{5}$ complex $\left(d^{7}\right)$.
2.3.

 $=\mathrm{e}_{g}$



$\mathrm{t}_{1 u}$

2.4. The $x y$ and $x^{2}-y^{2}$ orbitals are interchanged compared to the description given in $\S$ 2.2.1.

2.5. (1) A fourteen-electron $d^{8}$ complex; (2) and (3) Two empty nonbonding MOs:
2.6. (1) and (2) For the orbitals of the square-planar complex, see $\oint 2.2$. The $z^{2}$ orbital becomes strictly nonbonding (ligands on the nodal cone).

2.7.

1. See $\S$ 2.3.1.1, Figure 2.7.
2. (a) A low-spin $d^{7}$ complex with electronic configuration $(x y)^{2}(x z)^{2}(y z)^{2}\left(z^{2}\right)^{1}$. Compared to $d^{6}$ low-spin complexes, there is an additional electron in the $z^{2}$ orbital.

This orbital is particularly antibonding towards the apical ligand $\Rightarrow$ substantial lengthening of this bond.
(b) Two electrons in $z^{2}$ (low-spin $d^{8}$ ): even greater lengthening of the apical bond.
3. A high-spin $d^{4}$ complex with electronic configuration $(x y)^{1}(x z)^{1}(y z)^{1}\left(z^{2}\right)^{1}$. The $\mathrm{M}-\mathrm{Cl}_{a}$ bond will be longer than $\mathrm{M}-\mathrm{Cl}_{b}$, due to the electron in $z^{2}\left(\exp : \mathrm{M}-\mathrm{Cl}_{a}=2.58 \AA\right.$, $\mathrm{M}-\mathrm{Cl}_{b}=2.30 \AA$ ).
4. $(x y)^{2}(x z)^{2}(y z)^{2}\left(z^{2}\right)^{2} \rightarrow(x y)^{2}(x z)^{2}(y z)^{2}\left(z^{2}\right)^{1}\left(x^{2}-y^{2}\right)^{1}$. Loss of an electron from $z^{2}\left(\mathrm{M}-\mathrm{L}_{a}\right.$ antibonding) shortens the $\mathrm{M}-\mathrm{L}_{a}$ bond. Addition of an electron to $x^{2}-y^{2}\left(\mathrm{M}-\mathrm{L}_{b}\right.$ antibonding) lengthens the $\mathrm{M}-\mathrm{L}_{b}$ bonds. Compared to the case of a low-spin complex $\left(\mathrm{M}-\mathrm{L}_{a}>\mathrm{M}-\mathrm{L}_{b}\right)$, the bond lengths tend to become equal in a high-spin complex.
2.8.

1. See $\S$ 2.2.1, Figure 2.6 and $\S$ 2.3.1, Figure 2.9.
2. Low-spin $d^{8}\left(\mathrm{~b}_{2 g}\right)^{2}\left(\mathrm{e}_{g}\right)^{4}\left(\mathrm{a}_{1 g}\right)^{2}$ for square-planar; $d^{8}$ high-spin $(e)^{4}\left(\mathrm{t}_{2}\right)^{4}$ for the tetrahedron, with two unpaired electrons in the $t_{2}$ block.
3. Square-planar: six electrons in nonbonding orbitals, and two in very weakly antibonding orbitals, in the d block; tetrahedron: four electrons in nonbonding orbitals and four in weakly antibonding orbitals. The bonds are weaker in the tetrahedron.
2.9.
4. 



2. Compare with Figure 2.15 (right-hand side): the $x^{2}-y^{2}$ orbital $\left(1 \mathrm{a}_{1}\right)$ is slightly antibonding, since the two ligands are not exactly on the nodal planes, as they are when $L_{1}-M-L_{2}=$ $90^{\circ}$. However, the $x y$ orbital $\left(\mathrm{b}_{2}\right)$ is less antibonding, as in this case the three ligands are no longer pointing exactly towards the regions of greatest amplitude for $x y$. The gap between $x^{2}-y^{2}$ and $x y$ therefore becomes smaller.
2.10. For the $\mathrm{ML}_{4}$ complex, use the orientation given in Exercise 2.4. The result obtained is exactly the same as that established in $\oint$ 2.8.4 (Figure 2.16), since the bond angles are $90^{\circ}$ in both cases.
2.11. (1) and (2) The symmetry-adapted orbitals on the ligands have $a_{1}$ and $t_{2}$ symmetries (Chapter 6, $\S 6.6 .2$ ). Due to the (weakly) bonding overlaps, the $a_{1}$ orbital is lower in energy than the $t_{2}$ orbitals. The character table for the $T_{d}$ point group (Chapter $6, \S 6.6 .2$, Table 6.20) shows the symmetries of the metal orbitals: $\mathrm{a}_{1}(s), \mathrm{t}_{2}$ $(p), \mathrm{e} \oplus \mathrm{t}_{2}(d)$. The following interaction diagram is obtained, characterized by a two-orbital interaction ( $\mathrm{a}_{1}$ ), interactions involving two sets of three orbitals $\left(\mathrm{t}_{2}\right)$, and nonbonding e orbitals (the order of the antibonding MOs $2 \mathrm{a}_{1}$ and $3 \mathrm{t}_{2}$ is not obvious, and may depend on the particular system considered):

2.12. Trigonal-planar $\mathrm{ML}_{3}$ complex

2.13.

1. The $s$ orbital, which is totally symmetric like $z^{2}$ ( $\mathrm{a}_{1}$ symmetry).
2. The interaction between the $s$ orbital on the metal and the ligand orbitals is bonding. As a result, the amplitude of $z^{2}$ decreases in the plane of the complex (xy) but increases along the $z$-axis. The metal-ligand interactions are therefore less antibonding after the polarization.

## Chapter 3: $\boldsymbol{\pi}$-type interactions

3.1.

1. See Figure 2.6 (right-hand side)
2. (a) $z^{2}$ is unchanged (overlap zero by symmetry with the lone pairs $\left(p_{\mathrm{Cl}}\right)$ on the chloride ligands); one orbital $(x y)$ is destabilized by four antibonding interactions with $p_{\mathrm{Cl}}$; two orbitals ( $x z$ and $y z$ ) are destabilized less strongly by
two antibonding interactions with $p_{\mathrm{Cl}}$. Without detailed calculations, the position of $z^{2}$ with respect to the three other orbitals is not clear.

(b) $x y$ is now stabilized by four bonding interactions with the $\pi_{\mathrm{CN}}^{*}$ orbitals (CN is a double-face $\pi$-acceptor, like CO), $x z$ and $y z$ by two bonding interactions.

(c) $z^{2}$ and one of the three nonbonding MO are unchanged, but the two others are destabilized by an antibonding interaction with one $p_{\mathrm{Cl}}$.

3.2. 1. See Figure 2.10.
3. 1a: Large overlaps between $y z$ (nonbonding) and $\pi_{y}^{*}$, and between $x z$ (nonbonding) and $\pi_{x}^{*}$; very small overlaps between $x^{2}-y^{2}$ and $\pi_{y}^{*}$ (one bonding and one antibonding interaction, which do not completely cancel due to the
polarization of $x^{2}-y^{2}($ see question 1$)$ ); the same situation for $x y$ and $\pi_{x}^{*}$.
1b: By symmetry (the $x y$ and $y z$ planes), there are interactions between $y z$ (nonbonding) and $\pi_{z}^{*}$, and between $x y$ (antibonding) and $\pi_{x}^{*}$.


There are two stabilizing interactions in each case. The strongest interaction involves $x y$ (antibonding) and $\pi_{x}^{*}$ (these orbitals are closest in energy, and the overlap is larger due to the polarization of $x y$ ).
3. The most favourable site for substitution so far as the $\pi$ interactions are concerned is axial for a $d^{4}$ complex and equatorial for a $d^{8}$ complex.
4. The same type of interactions, but antibonding (a double-face $\pi$-donor ligand).
5. $d^{4}$ : equatorial substitution (one destabilizing interaction instead of two for axial substitution); $d^{8}$ : two destabilizing interactions in both isomers, but axial substitution is favoured since one of the overlaps is larger for equatorial substitution.
3.3.

1. See Chapter 1, § 1.1.1.2
2. (i) a nonbonding $\sigma$ orbital, concentrated on C ; (ii) see Chapter 1, Figure 1.7; (iii) a $\pi$ acceptor; (iv) through carbon.
3. The MO are similar to those in $\left[\mathrm{M}(\mathrm{CO})_{6}\right]$ (see Scheme 3-35).
4. The bonds are shorter in $\left[\mathrm{Fe}(\mathrm{CN})_{6}\right]^{4-}$ (an $\mathrm{Fe}(\mathrm{II})$ complex, with six electrons in $\pi$-bonding MO in the d block) than in $\left[\mathrm{Fe}(\mathrm{CN})_{6}\right]^{3-}$ (an $\mathrm{Fe}(\mathrm{III})$ complex with only five electrons in $\pi$-bonding MO).
5. No, since the $\mathrm{Fe}^{3+}$ cation is smaller than the $\mathrm{Fe}^{2+}$ cation.
6. (i) $d^{3}$ for each of them; (ii) CN is a $\pi$ acceptor, F a $\pi$ donor (see Figure 3.8).
3.4.
7. $x z\left(P_{1}\right)$ and $x y\left(P_{2}\right)$.
8. $x^{2}-y^{2}: \mathbf{S S} ; x z: \mathbf{S A} ; y z: \mathbf{A A}$
9. 


SA

AA

SS

AS
4. Three.
5. and 6 .

7. (i) 12 ; (ii) 18 .
3.5.

1. Two, $x z$ and $y z$.
2. (a) Two bonding interactions with the $\pi_{\mathrm{CO}}^{*}$ orbitals for axial substitution, only one for equatorial substitution;
(b) axial.
3.6. Use the orbitals established in Exercise 6.13 (questions 4 and 6), but for $\left[\operatorname{Pt}(\mathrm{CN})_{4}\right]^{2-}$, the $p$ orbitals on the ligands should be replaced by the $\pi_{\mathrm{CN}}^{*}$ orbitals. The interaction diagram may be obtained by considering the symmetry properties of these orbitals and of the
metal d orbitals (see Table 6.18). There are three interactions: two involve $\mathrm{e}_{g}$ orbitals, and one concerns $\mathrm{b}_{2 g}$ orbitals.
3.7.
3. See Chapter 2, Figure 2.11.
4. Use Figure 6.11.
5. Two interactions between $e^{\prime}$ orbitals and two between $e^{\prime \prime}$ orbitals.
6. It is not possible to establish easily the two-fold degeneracy of the orbitals.

## Chapter 4: Applications

4.1.

1. Structure 2 (the same analysis as in $\S$ 4.1.1).
2. and 3.: See T. A. Albright, R. Hoffmann, J. C. Thibeault, D. Thorn J. Am. Chem. Soc. 101, 3801 (1979).
4.2.
3. See Chapter 2, § 2.8.2.
4. See Chapter 2, § 2.6.1 and Chapter 4, § 4.2.2.
5. A 'T-shaped' structure for $\left[\mathrm{Rh}\left(\mathrm{PPh}_{3}\right)_{3}\right]^{+}\left(d^{8}\right)$ and trigonal for $\left[\mathrm{Pt}\left(\mathrm{PPh}_{3}\right)_{3}\right]\left(d^{10}\right)$.
4.3.
6. See Chapter 2, § 2.1.2.4 (trans) and 2.1.2.5 (cis).
7. The same types of diagram as in the dichloro octahedral complexes (trans: see Chapter 3, § 3.3.1; cis: see Chapter 3, Exercise 3.4).
8. $d^{0}:$ cis; $d^{2}:$ trans; see D. M. P. Mingos J. Organomet. Chem. 179, C29 (1979).
9. $\left[\mathrm{OsO}_{2} \mathrm{~F}_{4}\right]^{2-}$ : trans; $\left[\mathrm{MoO}_{2} \mathrm{Cl}_{4}\right]^{2-}$ : cis.
4.4.
10. $d^{6}$
11. See the analysis in Chapter 4, § 4.1.3 and A. Jarid, A. Lledos, D. Lauvergnat, Y. Jean New J. Chem. 21, 953 (1997).
12. $(0,0),(0,90)$, and $(90,90)$.
13. Coplanar or perpendicular, the first being more stable.
4.5.
14. See the analysis for the bis-ethylene complexes (Chapter 4, § 4.1.3).
15. $d^{2}: 1 ; d^{4}: 2 ; d^{6}: 2$.
16. The back-donation interactions towards the carbenes decrease, so the rotation barrier also decreases.
17. For the $d$ orbitals in the octahedron, use the sketches given in Figure 2.4. For those in the square-planar structure, use the representation in Exercise 2.4 (rotated by $90^{\circ}$ ). For the MO that characterize the broken M-H bonds, see Figure 4.12.
18. Yes.
19. An orbital in the reactant that characterizes a bond becomes a d-block orbital in the product.

## Chapter 5: The isolobal analogy

5.1.

1. $\mathrm{CH}_{3}-\mathrm{CH}_{3} ;\left[(\mathrm{CO})_{5} \mathrm{Mn}-\mathrm{Mn}(\mathrm{CO})_{5}\right]$
2. 




$\mathrm{CH}_{3}$
$\left[\mathrm{Mn}(\mathrm{CO})_{5} \mathrm{CH}_{3}\right]$
5.2.
1.

2. $[\mathrm{CpRh}(\mu-\mathrm{CO})]_{2}$, with two bridging CO .
5.3.

1. and 2 .

5.4.
2. 

$\left[\mathrm{Fe}(\mathrm{CO})_{4}\right]: d^{8}-\mathrm{ML}_{4} \longleftrightarrow \mathrm{CH}_{2}$

$[\mathrm{Pt}(\mathrm{CO}) \mathrm{L}]: d^{10}-\mathrm{ML}_{2} \longleftrightarrow 0 \mathrm{CH}_{2}$

2.


Planar coordination around the platinum atoms, octahedral around the iron.

5.5.

1. $\mathrm{CH}_{2}^{+}$.
2. $\mathrm{BH}_{2}$.
3. Diborane $\mathrm{B}_{2} \mathrm{H}_{6}$.
5.6.
4. (a).
5. (b) (analogues of cyclopropane).
5.7.
6. 

$\left[\operatorname{Ir}(\mathrm{CO})_{3}\right]: d^{9}-\mathrm{ML}_{3} \longleftrightarrow 0 \mathrm{CH} \quad$ (scheme 5-11, last line)
2. These are also isolobal analogues of tetrahedrane.
3. Tetrahedral representations, like structures 1 and 2 in the question.
5.8. Yes for the organic molecule ('orthogonal' ethylene), but not for the complex. There is a $\pi$ interaction, weaker than that in conformation (a), which involves a nonbonding $d$ orbital on the metal (see Chapter 4, § 4.1.1.).

## Chapter 6: Elements of group theory and applications

6.1. (1) Two; (2) three, each of which contains an $\mathrm{N}-\mathrm{H}$ bond and bisects the opposite $\mathrm{H}-\mathrm{N}-\mathrm{H}$ angle (planes written $\sigma_{\mathrm{v}}(1), \sigma_{\mathrm{v}}(2)$, and $\left.\sigma_{\mathrm{v}}(3)\right)$; (3) three: the molecular plane and two planes perpendicular to it, one perpendicular to the $\mathrm{C}-\mathrm{C}$ bond, the other bisecting the $\mathrm{H}-\mathrm{C}-\mathrm{H}$ angles; (4) two, since the third plane described above for ethylene is not a symmetry element here ( $\mathrm{H}-\mathrm{C}-\mathrm{F}$ angles); (5) one, the molecular plane; (6) four: the molecular plane and the three planes perpendicular to it which each contain an $\mathrm{Al}-\mathrm{Cl}$ bond.
6.2. (1) Yes; (2) no; (3) no; (4) yes; (5) no; (6) yes; (7) no; (8) no; (9) yes; (10) yes; (11) no.
6.3.

1. No.
2. The $C_{3}$ axes are co-linear with one of the $\mathrm{C}-\mathrm{H}$ bonds; the $C_{2}$ axes bisect two opposite $\mathrm{H}-\mathrm{C}-\mathrm{H}$ angles.
3. (a) three, at the intersections of the three planes found in Exercise 6.1 (question 3);
(b) one, the intersection of the two reflection planes (Exercise 6.1, question 4);
(c) one, perpendicular to the molecular plane.
4. a $C_{3}$ axis perpendicular to the molecular plane and three $C_{2}$ axes, each of which is co-linear with one of the $\mathrm{Al}-\mathrm{Cl}$ bonds.
6.4.
5. $S_{3}^{2}=C_{3}^{2}$ and $S_{3}^{4}=C_{3}$.
6. (i) $S_{3}$; (ii) $S_{6}$.
6.5.
7. $\sigma_{\mathrm{v}}(1) C_{3}=\sigma_{\mathrm{v}}(2) ; C_{3} \sigma_{\mathrm{v}}(1)=\sigma_{\mathrm{v}}(3)$ (non-commutative).
8. $\left(C_{3} \sigma_{\mathrm{v}}(1)\right) \sigma_{\mathrm{v}}(2)=C_{3}\left(\sigma_{\mathrm{v}}(1) \sigma_{\mathrm{v}}(2)\right)=C_{3}^{2}$ (associative).
9. $E, C_{3}^{2}, C_{3}, \sigma_{\mathrm{v}}(1)$, respectively.
6.6. Consult § 6.2.2. (1) point 1: $D_{\infty h}$; (2) point 1: $C_{\infty v}$; (3) point 4: $D_{2 h}$; (4) point 3: $C_{2 v}$; (5) point 3: $C_{2 h}$; (6) point 4: $D_{3 h}$; (7) point 1: $T_{d} ;(8)$ point 3: $C_{2 v} ;(9)$ point 3: $C_{3 v} ;(10)$ point 4: $D_{4 h} ;(11)$ point 4: $D_{3 h}$; (12) point 2: $C_{s}$.
6.7.

$$
\begin{aligned}
& \int_{\text {space }}\left[2\left(1 s_{\mathrm{H}_{1}}\right)-\left(1 s_{\mathrm{H}_{2}}\right)-\left(1 s_{\mathrm{H}_{3}}\right)\right] \times\left[1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}+1 s_{\mathrm{H}_{3}}\right] \mathrm{d} \tau \\
& \quad=(2-1-1)+S(2+2-1-1-1-1)=0
\end{aligned}
$$

6.8.

1. $\Gamma_{\mathrm{H}}=\mathrm{A}_{g} \oplus \mathrm{~B}_{1 g} \oplus \mathrm{~B}_{2 u} \oplus \mathrm{~B}_{3 u}$.
2. 

$$
\begin{aligned}
& \mathrm{A}_{g}: \frac{1}{2}\left(1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}+1 s_{\mathrm{H}_{3}}+1 s_{\mathrm{H}_{4}}\right) \\
& \mathrm{B}_{1 g}: \frac{1}{2}\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}+1 s_{\mathrm{H}_{3}}-1 s_{\mathrm{H}_{4}}\right) \\
& \mathrm{B}_{2 u}: \frac{1}{2}\left(1 s_{\mathrm{H}_{1}}-1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}+1 s_{\mathrm{H}_{4}}\right) \\
& \mathrm{B}_{3 u}: \frac{1}{2}\left(1 s_{\mathrm{H}_{1}}+1 s_{\mathrm{H}_{2}}-1 s_{\mathrm{H}_{3}}-1 s_{\mathrm{H}_{4}}\right)
\end{aligned}
$$

6.9.

6.10.

$$
\begin{aligned}
\phi_{\mathrm{A}_{1}} & =\frac{1}{2}\left(\sigma_{1}+\sigma_{2}+\sigma_{3}+\sigma_{4}\right) \\
\phi_{\mathrm{T}_{2}}(2) & =\frac{1}{\sqrt{2}}\left(\sigma_{3}-\sigma_{4}\right) \\
\phi_{\mathrm{T}_{2}}(1) & =\lambda \sigma_{1}-\mu\left(\sigma_{2}+\sigma_{3}+\sigma_{4}\right)
\end{aligned}
$$

where $\lambda^{2}+3 \mu^{2}=1$ (normalization) $\phi_{\mathrm{T}_{2}}(1)$ and $\phi_{\mathrm{A}_{1}}$ orthogonal $\Rightarrow \lambda=3 \mu$

$$
\begin{aligned}
\Rightarrow \phi_{\mathrm{T}_{2}}(1) & =\frac{1}{\sqrt{12}}\left(3 \sigma_{1}-\sigma_{2}-\sigma_{3}-\sigma_{4}\right) \\
\phi_{\mathrm{T}_{2}}(3) & =\lambda \sigma_{2}-\mu\left(\sigma_{3}+\sigma_{4}\right)
\end{aligned}
$$

where $\lambda^{2}+2 \mu^{2}=1$ (normalization)
$\phi_{\mathrm{T}_{2}}(3)$ and $\phi_{\mathrm{A}_{1}}$ orthogonal $\Rightarrow \lambda=2 \mu$

$$
\Rightarrow \phi_{\mathrm{T}_{2}}(3)=\frac{1}{\sqrt{6}}\left(2 \sigma_{2}-\sigma_{3}-\sigma_{4}\right)
$$

6.11. 1. and 2 .

| $D_{3 h}$ | E | $C_{3}^{1}$ | $C_{3}^{2}$ | $C_{2}(1)$ | $C_{2}(2)$ | $C_{2}(3)$ | $\sigma_{h}$ | $S_{3}^{1}$ | $S_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $R_{k}\left(\sigma_{1}\right)$ | $\sigma_{1}$ | $\sigma_{2}$ | $\sigma_{3}$ | $\sigma_{1}$ | $\sigma_{3}$ | $\sigma_{2}$ | $\sigma_{1}$ | $\sigma_{2}$ | $\sigma_{3}$ | $\sigma_{1}$ | $\sigma_{3}$ |
| $R_{k}\left(\sigma_{2}-\sigma_{3}\right)$ | $\sigma_{2}-\sigma_{3}$ | $\sigma_{3}-\sigma_{1}$ | $\sigma_{1}-\sigma_{2}$ | $\sigma_{3}-\sigma_{2}$ | $\sigma_{2}-\sigma_{1}$ | $\sigma_{1}-\sigma_{3}$ | $\sigma_{2}-\sigma_{3}$ | $\sigma_{3}-\sigma_{1}$ | $\sigma_{1}-\sigma_{2}$ | $\sigma_{3}-\sigma_{2}$ | $\sigma_{2}-\sigma_{1}$ |
| $\mathrm{~A}_{1}^{\prime}$ | 1 | 1 | $\sigma_{1}-\sigma_{3}$ |  |  |  |  |  |  |  |  |
| $\mathrm{E}^{\prime}$ | 2 | -1 | -1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |

$\phi_{\mathrm{A}_{1^{\prime}}}=4 \times\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right)$, so the normalized function is $(1 / \sqrt{3})\left(\sigma_{1}+\sigma_{2}+\sigma_{3}\right)$.
$\phi_{E^{\prime}}(1)=2 \sigma_{1}-\sigma_{2}-\sigma_{3}+2 \sigma_{1}-\sigma_{2}-\sigma_{3}=4 \sigma_{1}-2 \sigma_{2}-2 \sigma_{3}$, so the normalized function is $(1 / \sqrt{6})\left(2 \sigma_{1}-\sigma_{2}-\sigma_{3}\right)$.

$$
\phi_{\mathrm{E}^{\prime}}(2)=2 \times\left(\sigma_{2}-\sigma_{3}\right)-\left(\sigma_{3}-\sigma_{1}\right)-\left(\sigma_{1}-\sigma_{2}\right)+2 \times\left(\sigma_{2}-\right.
$$ $\left.\sigma_{3}\right)-\left(\sigma_{3}-\sigma_{1}\right)-\left(\sigma_{1}-\sigma_{2}\right)=6 \sigma_{2}-6 \sigma_{3}$ so the normalized function is $(1 / \sqrt{2})\left(\sigma_{2}-\sigma_{3}\right)$.

6.12. 1. and 2 .

| $D_{3 h}$ | $E$ | $C_{3}^{1}$ | $C_{3}^{2}$ | $C_{2}(1)$ | $C_{2}(2)$ | $C_{2}(3)$ | $\sigma_{h}$ | $S_{3}^{1}$ | $S_{3}^{2}$ | $\sigma_{v}(1)$ | $\sigma_{v}(2)$ | $\sigma_{v}(3)$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $R_{k}\left(p_{\\| 1}\right)$ | $p_{\\| 1}$ | $p_{\\| 2}$ | $p_{\\| 3}$ | $-p_{\\| 1}$ | $-p_{\\| 3}$ | $-p_{\\| 2}$ | $p_{\\| 1}$ | $p_{\\| 2}$ | $p_{\\| 3}$ | $-p_{\\| 1}$ | $-p_{\\| 3}$ | $-p_{\\| 2}$ |
| $R_{k}\left(p_{\\| 2}-p_{\\| 3}\right)$ | $p_{\\| 2}-p_{\\| 3}$ | $p_{\\| 3}-p_{\\| 1}$ | $p_{\\| 1}-p_{\\| 2}$ | $p_{\\| 2}-p_{\\| 3}$ | $p_{\\| 1}-p_{\\| 2}$ | $p_{\\| 3}-p_{\\| 1}$ | $p_{\\| 2}-p_{\\| 3}$ | $p_{\\| 3}-p_{\\| 1}$ | $p_{\\| 1}-p_{\\| 2}$ | $p_{\\| 2}-p_{\\| 3}$ | $p_{\\| 1}-p_{\\| 2}$ | $p_{\\| 3}-p_{\\| 1}$ |
| $\mathrm{~A}_{2}^{\prime}$ | 1 | 1 | 1 | -1 | -1 | -1 | 1 | 1 | 1 | -1 | -1 | -1 |
| $\mathrm{E}^{\prime}$ | 2 | -1 | -1 | 0 | 0 | 0 | 2 | -1 | -1 | 0 | 0 | 0 |

$\phi_{\mathrm{A}_{2}^{\prime}}=4 \times\left(p_{\| 1}+p_{\| 2}+p_{\| 3}\right)$, so the normalized function is $(1 / \sqrt{3})\left(p_{\| 1}+p_{\| 2}+p_{\| 3}\right)$.

$$
\phi_{\mathrm{E}^{\prime}}(1)=2 p_{\| 1}-p_{\| 2}-p_{\| 3}+2 p_{\| 1}-p_{\| 2}-p_{\| 3}=2 \times
$$ $\left(2 p_{\| 1}-p_{\| 2}-p_{\| \mid 3}\right)$, so the normalized function is $(1 / \sqrt{6})\left(2 p_{\| 1}\right.$ $\left.-p_{\| 2}-p_{\| 3}\right)$.

$\phi_{\mathrm{E}^{\prime}}(2)=2 \times\left(p_{\| 2}-p_{\| 3}\right)-\left(p_{\| 3}-p_{\| 1}\right)-\left(p_{\| 1}-p_{\| 2}\right)+2 \times$ $\left(p_{\| 2}-p_{\| 3}\right)-\left(p_{\| 3}-p_{\| 1}\right)-\left(p_{\| 1}-p_{\| 2}\right)=6 \times\left(p_{\| 2}-p_{\| 3}\right)$, so the normalized function is $(1 / \sqrt{2})\left(p_{\| 2}-p_{\| 3}\right)$.
6.13. 1. and 2 .

| $D_{4 h}$ | $E$ | $2 C_{4}$ | $C_{2}$ | $2 C_{2}^{\prime}$ | $2 C_{2}^{\prime \prime}$ | $i$ | $2 S_{4}$ | $\sigma_{h}$ | $2 \sigma_{v}$ | $2 \sigma_{d}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\Gamma_{p \perp}$ | 4 | 0 | 0 | -2 | 0 | 0 | 0 | -4 | 2 | 0 |

$$
=>\Gamma_{p \perp}=\mathrm{A}_{2 u} \oplus \mathrm{~B}_{2 u} \oplus \mathrm{E}_{g}
$$

3. 

| $D_{4 h}$ | $E$ | $C_{4}^{1}$ | $C_{4}^{3}$ | $C_{2}$ | $C_{2 a}^{\prime}$ | $C_{2 b}^{\prime}$ | $C_{2 a}^{\prime \prime}$ | $C_{2 b}^{\prime \prime}$ |
| :--- | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| $R_{k}\left(p_{\perp 1}\right)$ | $p_{\perp 1}$ | $p_{\perp 4}$ | $p_{\perp 2}$ | $p_{\perp 3}$ | $-p_{\perp 1}$ | $-p_{\perp 3}$ | $-p_{\perp 4}$ | $-p_{\perp 2}$ |
| $R_{k}\left(p_{\perp 2}\right)$ | $p_{\perp 2}$ | $p_{\perp 1}$ | $p_{\perp 3}$ | $p_{\perp 4}$ | $-p_{\perp 4}$ | $-p_{\perp 2}$ | $-p_{\perp 3}$ | $-p_{\perp 1}$ |
| $\mathrm{~A}_{2 u}$ | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 |
| $\mathrm{~B}_{2 u}$ | 1 | -1 | -1 | 1 | -1 | -1 | 1 | 1 |
| $\mathrm{E}_{g}$ | 2 | 0 | 0 | -2 | 0 | 0 | 0 | 0 |
| $D_{4 h}$ | $i$ | $S_{4}^{1}$ | $S_{4}^{3}$ | $\sigma_{h}$ | $\sigma_{v a}$ | $\sigma_{v b}$ | $\sigma_{d a}$ | $\sigma_{d b}$ |
| $R_{k}\left(p_{\perp 1}\right)$ | $-p_{\perp 3}$ | $-p_{\perp 4}$ | $-p_{\perp 2}$ | $-p_{\perp 1}$ | $p_{\perp 1}$ | $p_{\perp 3}$ | $p_{\perp 4}$ | $p_{\perp 2}$ |
| $R_{k}\left(p_{\perp 2}\right)$ | $-p_{\perp 4}$ | $-p_{\perp 1}$ | $-p_{\perp 3}$ | $-p_{\perp 2}$ | $p_{\perp 4}$ | $p_{\perp 2}$ | $p_{\perp 3}$ | $p_{\perp 1}$ |
| $\mathrm{~A}_{2 u}$ | -1 | -1 | -1 | -1 | 1 | 1 | 1 | 1 |
| $\mathrm{~B}_{2 u}$ | -1 | 1 | 1 | -1 | 1 | 1 | -1 | -1 |
| $\mathrm{E}_{g}$ | 2 | 0 | 0 | -2 | 0 | 0 | 0 | 0 |

$\phi_{\mathrm{A}_{2 u}}=4 \times\left(p_{\perp 1}+p_{\perp 2}+p_{\perp 3}+p_{\perp 4}\right)$, so the normalized function is

$$
\begin{aligned}
\phi_{\mathrm{A}_{2 u}} & =\frac{1}{2} \times\left(p_{\perp 1}+p_{\perp 2}+p_{\perp 3}+p_{\perp 4}\right) \\
\phi_{\mathrm{B}_{2 u}} & =\frac{1}{2} \times\left(p_{\perp 1}-p_{\perp 2}+p_{\perp 3}-p_{\perp 4}\right) \quad(\text { normalized }) \\
\phi_{\mathrm{E}_{g}}(1) & =\frac{1}{\sqrt{2}}\left(p_{\perp 1}-p_{\perp 3}\right) \quad(\text { normalized }) \\
\phi_{\mathrm{E}_{g}}(2) & =\frac{1}{\sqrt{2}}\left(p_{\perp 2}-p_{\perp 4}\right) \quad(\text { normalized })
\end{aligned}
$$

4. 




5. $\varepsilon_{\mathrm{A}_{2 u}}<\varepsilon_{\mathrm{Eg}}<\varepsilon_{\mathrm{B}_{2 u}}$
6.

| $D_{4 h}$ | $E$ | $2 C_{4}$ | $C_{2}$ | $2 C_{2}^{\prime}$ | $2 C_{2}^{\prime \prime}$ | $i$ | $2 S_{4}$ | $\sigma_{h}$ | $2 \sigma_{v}$ | $2 \sigma_{d}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\Gamma_{p \\|}$ | 4 | 0 | 0 | -2 | 0 | 0 | 0 | 4 | -2 | 0 |

$\Rightarrow \Gamma_{p \|}=\mathrm{A}_{2 g} \oplus \mathrm{~B}_{2 g} \oplus \mathrm{E}_{u}$
Normalized functions:

$$
\begin{aligned}
\phi_{A_{2 g}} & =\frac{1}{2} \times\left(p_{\| 1}+p_{\| 2}+p_{\| 3}+p_{\| 4}\right) \\
\phi_{\mathrm{B}_{2 g}} & =\frac{1}{2} \times\left(p_{\| 1}-p_{\| 2}+p_{\| 3}-p_{\| 4}\right) \\
\phi_{\mathrm{E}_{u}}(1) & =\frac{1}{\sqrt{2}}\left(p_{\| 1}-p_{\| 3}\right) \\
\phi_{\mathrm{E}_{u}}(2) & =\frac{1}{\sqrt{2}}\left(p_{\| 2}-p_{\| 4}\right)
\end{aligned}
$$





where $\varepsilon_{\mathrm{B} 2 g}<\varepsilon_{\mathrm{E}_{u}}<\varepsilon_{\mathrm{A}_{2 g}}$

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$\left[\mathrm{Ni}(\mathrm{CN})_{4}\right]^{4-} 66$
$\left[\mathrm{Ni}(\mathrm{CN})_{5}\right]^{3-} 9,11,12,15,60,62,72,87$
$\mathrm{Ni}(\mathrm{CO})_{2} \quad 195$
$\mathrm{Ni}(\mathrm{CO})_{3} 74,138$
$\mathrm{Ni}(\mathrm{CO})_{4} 66$
$\left[\mathrm{Ni}\left(\mathrm{H}_{2} \mathrm{O}\right)_{6}\right]^{2+}$ 9, 11, 12
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[^0]:    ${ }^{3}$ The dimension of the representation is given by the value of the character associated with the identity operation $E$ ( 3 in the case of $\mathrm{T}_{1 \mathrm{u}}$ ).

[^1]:    ${ }^{6}$ The tetrahedral geometry is also favoured by the bonding MO, lower in energy than the d block, as shown by the experimental structures of $d^{0}$ complexes (§ 2.4.2.2) which have tetrahedral geometries.

[^2]:    ${ }^{3}$ These expressions are not normalized.

[^3]:    ${ }^{8}$ J. Chatt, L. A. Duncanson J. Chem. Soc. 2939 (1953); M.J.S. Dewar Bull. Soc. Chim. Fr. 18, C79 (1951).

[^4]:    ${ }^{1}$ If the symmetry labels of the $C_{2 v}$ group are used, the symmetries $\mathbf{S S}, \mathbf{S A}, \mathbf{A S}$, and AA correspond to $a_{1}, b_{1}, b_{2}$, and $a_{2}$, respectively.

[^5]:    ${ }^{8}$ If the opposite holds, the ground-state electronic configuration of the reactant correlates with a doubly excited configuration of the product. This would be a symmetry-forbidden reaction, according to Woodward and Hoffmann.

[^6]:    ${ }^{3}$ L. Pauling The Nature of Chemical Bond, 3rd edn., Cornell University Press, Ithaca, NY (1960).

